

PLASMA PROCESSING OF CELLULOSE SURFACES AND THEIR INTERACTIONS WITH FLUIDS

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To my lovely daughter, Athira Balu

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LIST OF SYMBOLS AND ABBREVIATIONS

θ^*	Contact angle of on rough surface
θ	Contact angle of on flat surface
r	Roughness factor
f	Fraction of solid water interface
θ_A	Advancing contact angle
θ_R	Receding contact angle
CA_{adv}	Advancing contact angle
CA_{rec}	Receding contact angle
γ_{SV}	Interfacial energies between solid-vapor
γ_{SL}	Interfacial energies between solid-liquid
γ_{LV}	Interfacial energies between liquid-vapor
μL	Microliter
F	Force exerted on the drop
W_{drop}	Width of the drop perpendicular to the direction of sliding
W_{dot}	Width of the ink dot
W_{line}	Width of the ink line
ρ	Density of the liquid
V	Volume of the drop
g	Acceleration due to gravity
α	Critical sliding angle
F_E	Experimental force for a drop to slide
F_P	Predicted force for a drop to slide
F_{PO}	Predicted force for a drop to slide in “outside” configuration
F_{PP}	Predicted force for a drop to slide in “perimeter” configuration

F_I	Predicted force for a drop to slide on ink film
F_{SH}	Predicted force for a drop to slide on SH paper
θ_{RI}	Receding contact angle on ink
θ_{RSH}	Receding contact angle on SH paper
θ_{AI}	Advancing contact angle on ink
θ_{ASH}	Advancing contact angle on SH paper
W_{adh}	Work of adhesion
L_{char}	Characteristic length scale
$V_{predicted}$	Predicted volume
β	Proportionality constant
W	Width of ink pattern
σ	Area fraction
MEMS	Micro-electromechanical systems
LOC	Lab on a Chip
CTMP	Chemothermomechanical pulp
FTMS	Fluorotrimethylsilane
PFPDMS	Pentafluorophenyldimethylsilane
FDA	Food and drug administration
PTFE	Polytetrafluoroethylene
HDMS	Hexamethyldisilazane
SH	Superhydrophobic
CA	Contact angle
ATRP	Atom transfer radical polymerization
AKD	Alkyl ketene dimer
POC	Point of care

LOP	Lab on a paper
SEM	Scanning electron microscope
ATR	Attenuated total reflectance
FTIR	Fourier transform infrared
XPS	X-ray photoelectron spectroscopy
PECVD	Plasma-enhanced chemical vapor deposition
PFE	Pentafluoroethane
RH	Relative humidity
RF	Radio frequency
FIB	Focused ion beam
ASTM	American society for testing and materials
HS-OD	Handsheet - overnight drying (50% hardwood and 50% softwood)
HS	Handsheet (50% hardwood and 50% softwood)
H	Handsheet (100% hardwood)
S	Handsheet (100% softwood)
sccm	Standard cubic centimeter
TAPPI	Technical association of pulp and paper
CP	Copy paper
PT	Paper towel
PDMS	Polydimethylsiloxane
PMMA	Polymethylmethacrylate
OEW	Optoelectrowetting
2D	Two-dimensional
3D	Three-dimensional
HELP	Hysteresis enabled lab-on-paper

I

Inkfilm

ACNT

Aligned carbon nanotubes

SUMMARY

Cellulose is a biodegradable, renewable, flexible, inexpensive, biopolymer which is abundantly present in nature. In spite of these inherent advantages, cellulose fibers cannot be used directly in a number of potential industrial applications because of their hydrophilic nature; a surface modification is often required to alter the surface properties of cellulose. This thesis work reports a fabrication method that results in superhydrophobic properties (contact angle (CA) $> 150^\circ$) on cellulose (paper) surfaces. Superhydrophobicity was obtained by domain-selective etching of amorphous portions of the cellulose fiber in an oxygen plasma, and by subsequently coating the etched surface with a thin fluorocarbon film deposited via plasma enhanced chemical vapor deposition from a pentafluoroethane precursor. Two forms of superhydrophobicity with vastly different degrees of adhesion were obtained by varying the plasma treatment conditions, in particular the duration of oxygen etching: “roll-off” (contact angle (CA): $166.7^\circ \pm 0.9^\circ$ and CA hysteresis: $3.4^\circ \pm 0.1^\circ$) and “sticky” (CA: $153.4^\circ \pm 4.7^\circ$ and CA hysteresis: $149.8^\circ \pm 5.8^\circ$) superhydrophobicity. The CA hysteresis could be tuned between the two extremes by adjusting the oxygen etching time to control the formation of nano-scale features on the cellulose fibers. The effects of fiber types (soft vs. hard wood) and paper making parameters on fabricating superhydrophobic paper were also investigated. There were no significant differences in the formation of the nano-scale features created via oxygen etching on paper substrates obtained from different fiber types and paper making parameters. Because “roll-off” superhydrophobicity is primarily determined by the nano-scale roughness, this property is therefore not significantly affected by the fiber types or

paper making parameters. While the fiber type does not affect “roll-off” or “sticky” superhydrophobicity, paper making process parameters affect the structure of the paper web on the micro-scale and thus lead to variations in “sticky” superhydrophobicity.

Superhydrophobic paper substrates were patterned with high surface energy ink deposited using a commercial desktop printer. The patterns could be used to manipulate the drag and extensional adhesion of water drops on the substrates. Classic ‘drag’ and ‘extensional’ adhesion expressions were used to model the behavior of water drops on basic dot and line patterns of variable dimensions. A fundamental understanding of the adhesive forces of water drops as a function of pattern shape and size was thus obtained. Based on this knowledge, patterned paper substrates were then designed and fabricated to perform simple unit operations, such as storage, transfer, mixing and merging of water drops. These basic functionalities were combined in the design of a simple two-dimensional lab-on-paper (LOP) device. Further studies of more complicated pattern shapes led to the generation of patterns that allowed directional mobility and tunable adhesion of water drops. These developments are critical for designing novel components for two-dimensional LOP devices such as flow paths, gates/diodes, junctions and drop size filters.

CHAPTER 1

INTRODUCTION

1.1 Background

1.1.1 Paper

Paper is one of the most significant inventions in the history of mankind. Applications range from toilet tissues to currency bills. Paper is defined as a felted sheet formed on a fine mesh from an aqueous suspension of cellulose fibers. The word paper derives its name from the Greek term papyrus, the stem of a reedy plant used as a writing material in ancient Egypt [1]. The first authenticated papermaking originated in China in 105 CE [2]. Paper making was considered an art for several centuries after that until it was extended to other parts of the world. In the 1700s, the introduction of science and engineering to the paper making process resulted in great advances with respect to both production and usage. In 2002, the world-wide production of paper exceeded 300 million tons and it is expected to reach 500 million tons by 2013 [1, 3]. For the past two centuries, a very large fraction of paper and paperboard production was used for printing, writing and packaging applications. Currently, paper is more than just a substrate for writing, printing and packaging; recent scientific research has established its potential as an inexpensive, biodegradable, renewable, flexible polymer substrate. Innovative concepts of paper-based devices include transistors [4, 5], batteries [6, 7], super-capacitors [7], MEMS devices [8], sensors [9] and lab-on-chip (LOC) microfluidic devices [10-16].

The fibers used to make paper come from a variety of sources such as wood, straws, grasses, canes, reeds and bamboos. In spite of the diversity of sources, all fibers contain the following three components: cellulose, hemicellulose, and lignin. Of the above, cellulose is the single most important component for papermaking [1]. As shown in Figure 1.1a, cellulose molecules $(C_6H_{10}O_5)_n$ are essentially polymer chains of β -d-glucose residues covalently coupled via glucosidal linkages [17, 18]. The 'n' value, also known as the degree of polymerization (DP), varies from 600 to 1500 depending on the fiber source and treatment techniques applied [1]. Cellulose molecules form long segments in several ordered levels of orientation developing crystalline and amorphous domains as shown in Figure 1.1b [19-22]. The crystalline segments constitute the microfibrils which are bundled together to form macrofibrils as shown in Figure 1.2 [20, 21, 23, 24]. A single fiber is made up of several aggregated macrofibrils. Hemi-celluloses are shorter-chain polysaccharide molecules with DP between 15 and 90. They constitute five different sugars to constitute the polymer chains: glucose, mannose, galactose, xylose and arabinose [1]. Lignins are complex amorphous polymers made from phenylpropane units [25] and give plants their brown color. Lignin's principal role in the wood is to cement the fibers together; it is thus responsible for the rigidity of the plants [26]. In addition to the three primary components, wood contains miscellaneous constituents such as resin acids, fats, terpenes and alcohols which can be extracted with water or organic solvents such as ethanol, acetone or dichloromethane and are thus termed 'extractives'.

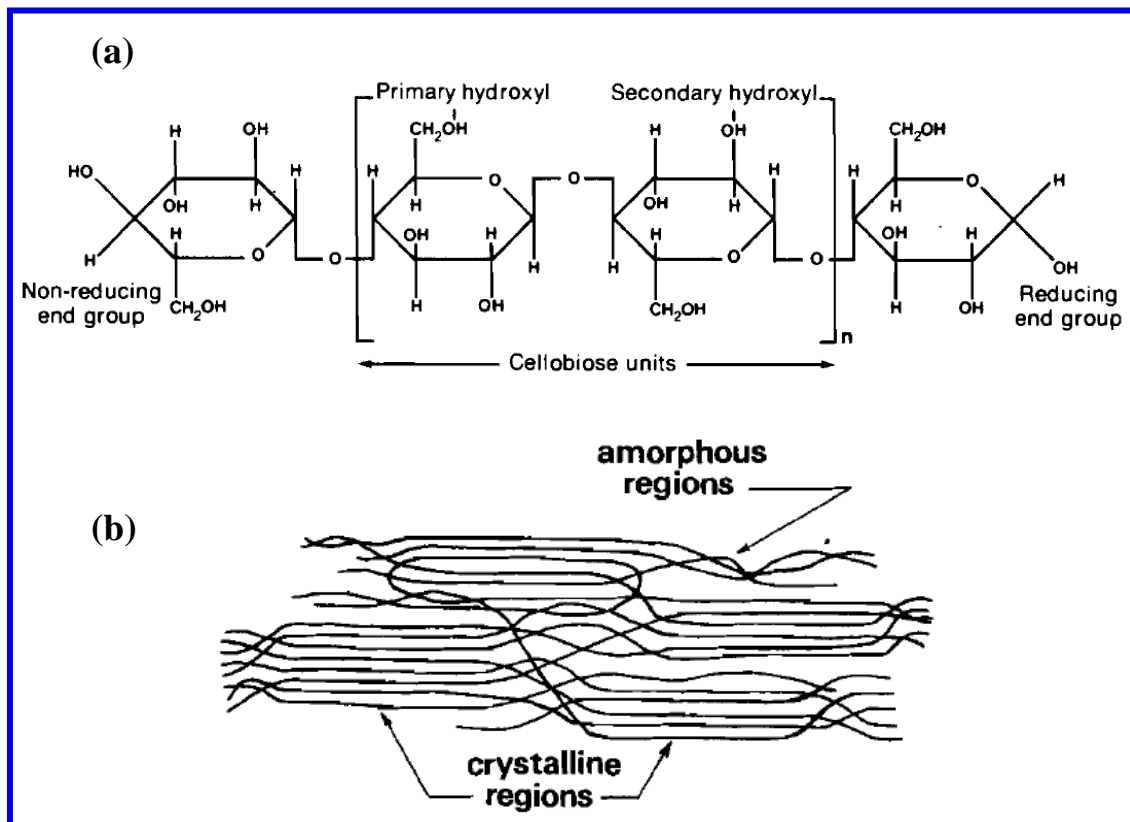


Figure 1.1 Cellulose structure (a) and schematic of molecular organization within a cellulose microfibril (b) [1]

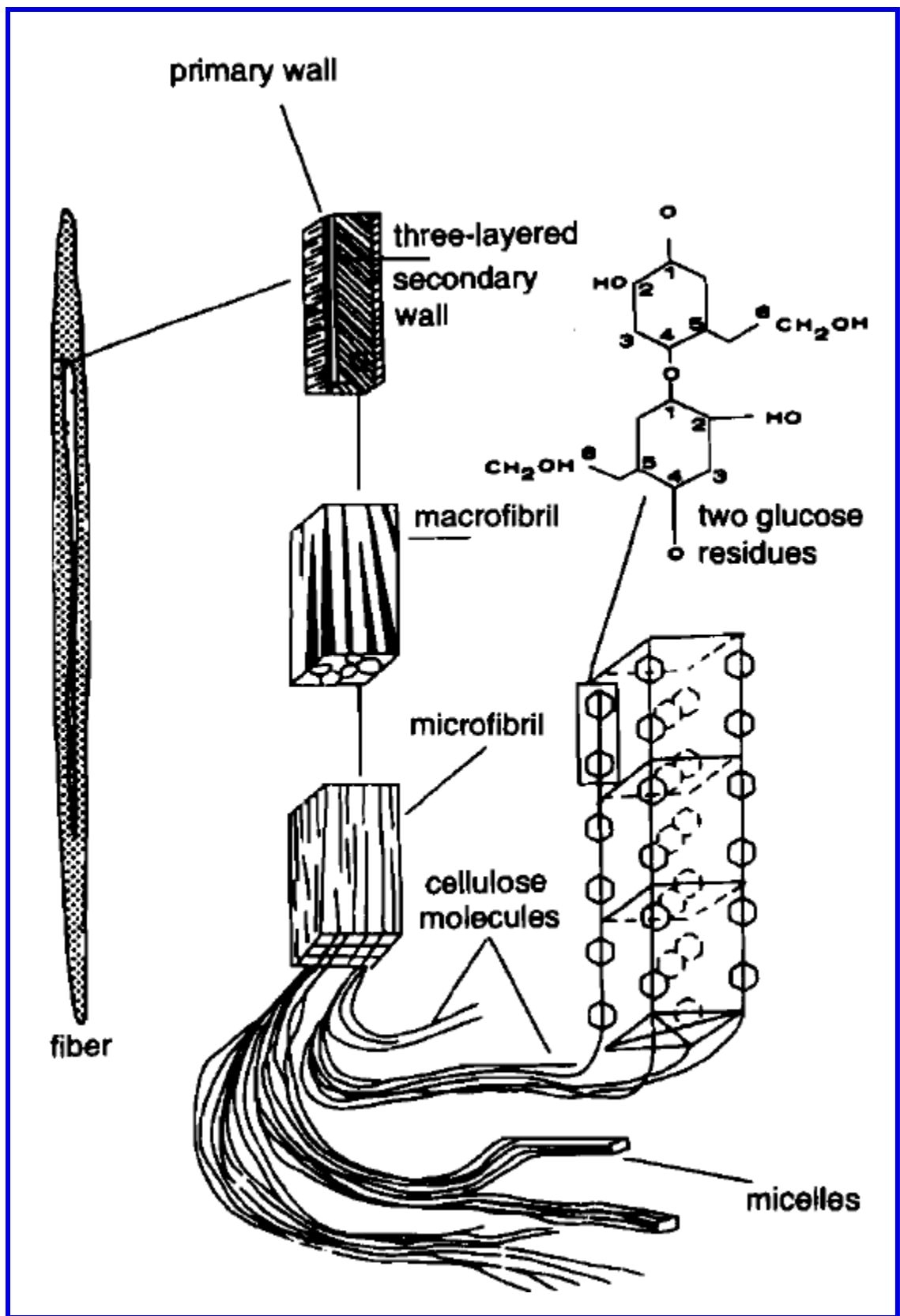


Figure 1.2 Microscopic and submicroscopic structure of cellulose [1]

1.1.2 Surface modification of paper

The hydrophilic nature of cellulose is critical in the paper making process [27] , since it is this property that bonds the fibers together, thereby allowing the production of paper. However, the surface chemistry and morphology of the cellulose fibers also restrict its use in a number of important industrial applications. Hence, the ability to modify cellulose fibers to tailor their physical and chemical properties without changing bulk paper properties is of considerable importance. Such a process whereby the surface morphology and chemistry of the solid surfaces are changed to control the interfacial interactions between solid-solid, solid-liquid and solid-gas, without modifying the bulk properties of the solid is known as surface modification.

Two important properties essential to all applications and that therefore drive research on the surface modification of paper are wettability and adhesion. Wettability is determined by the contact angle (θ) displayed by a liquid drop on a surface. If the liquid is water, a surface can be classified as hydrophilic ($\theta < 90^\circ$), hydrophobic ($\theta > 90^\circ$), or superhydrophobic ($\theta > 150^\circ$) based on the contact angle values. The wettability of paper surfaces is critical for applications such as printing and packaging. On the other hand, adhesion between fibers or to additives in the paper is critical in determining the internal bonding strength of paper and in making fiber reinforced composites with various polymer matrices.

Sizing is one of the traditional surface treatments applied to paper in order to control the wettability. Sizing methods can be classified into two categories [1]: internal sizing and

surface sizing. Internal sizing is accomplished by adding natural (rosin based) or synthetic (alkylketene dimer and alkenyl succinic anhydride) sizing agents to pulp slurries before forming the paper [28]. Surface sizing is performed in later stages of the process by filling the pores of the paper with starch and waxes, thereby retarding the rate of liquid penetration. The adhesion strength between fibers can be improved by adding starches, gums and resins to the pulp [1]. Cellulose fibers can also be modified via mechanical treatments (roughening and polishing), wet chemical treatments, and vapor phase treatments to obtain a variety of surface properties.

The traditional surface modification techniques are quite effective when considering traditional paper applications such as printing, writing and packaging. However, as mentioned in the previous section, applications of paper have expanded in the past decade, and the paper industry must implement innovations in their manufacturing processes to allow competition with plastics and biomaterials. Such considerations suggest that so called ‘out of the box thinking’ is needed to develop and implement novel surface treatment techniques that can ensure paper substrate compatibility with a variety of liquids and solid materials.

Vapor phase treatment, in particular cold plasma processing, has gained popularity in the past few decades because of its versatility to impart cellulose surfaces with various desirable surface properties by choosing the appropriate chemical reactants and experimental parameters. Moreover, since plasmas or glow discharges represent solvent free processes, they have been recently included in the list of “green” technologies [29].

The following sections will discuss the reaction mechanisms involved in plasma processing and will describe applications of plasmas to paper surface modifications.

1.1.3 Plasma processing

A plasma is a partially ionized quasineutral gas containing electrons (0-10 eV), ions (10-30 eV), photons (200-400 nm) and a variety of neutral species [30]. Irving Langmuir first observed this phenomenon in his lab in 1927 when he tried to strike an electrical discharge in a gas. In naming this interesting phenomenon, several suggestions were made by his research group: ‘uniform discharge’, ‘homogenous discharge’, ‘equilibrium discharge’, ‘auras’ and ‘haloes’. Finally, Langmuir suggested the term ‘plasma’, as the way the electric discharge carries the electrons, ions, neutrals and gas impurities reminded him of the way the ‘blood plasma’ carries red corpuscles, white corpuscles and germs [31, 32].

Typically, in plasma processing, organic gases or vapors are partially ionized by an electric discharge resulting in the creation of highly reactive species. Thus, by using electrical energy, even at room temperatures, highly reactive species can be generated that otherwise can only be obtained at higher temperatures via thermal excitation.

Generation of reactive species at low temperatures is a significant advantage when treating paper fibers which will decompose at high temperatures. For instance, the degree of polymerization (DP) of Kraft paper decreases by a factor of 7 around 150° C [33]. In plasmas the highly reactive species at ambient temperatures provide the capability to perform surface modification of paper without the loss of mechanical properties [30].

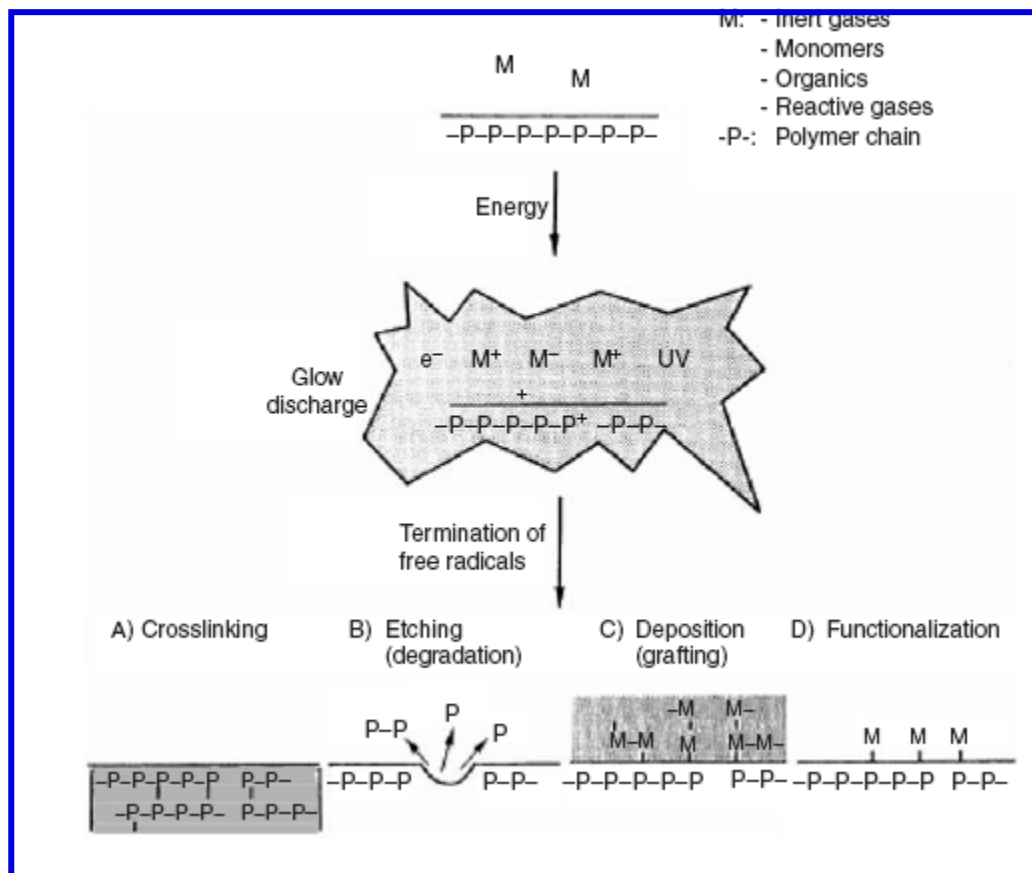


Figure 1.3 Reaction mechanisms for plasma surface modification [30]

The energetic species formed in the plasma may undergo various reactions with the surface as shown in Figure 1.3: a) crosslinking, b) etching (degradation), c) deposition (grafting) and d) functionalization. Apart from these reactions at the surfaces, there can also be crosslinking of gas species which results in the formation of particles in the vapor phase. All the above mentioned reactions can be controlled by selecting the appropriate type of gas or vapor, reactor design, and operating conditions, such as gas flow rate, pressure, power, and temperature of the substrate.

1.1.4 Plasma processing of paper surfaces

The level of interest in the use of plasma-based technologies to modify cellulose can be assessed by considering the number of publications reported in ISI Web of Science that contain the terms “plasma” and “cellulose”. The results from this search were scrutinized to identify only papers that are relevant to this study. As shown in Figure 1.4, the number of publications related to plasma modification of cellulose has increased exponentially every 10 years. This trend clearly shows the increasing popularity of using plasma based technologies for cellulose surface modification.

The first application of plasma processing to modify cellulose surfaces dates back to the 1970s, when a microwave plasma was used to modify the bonding properties of cellulose [34]. Thereafter, plasma processing of cellulose surfaces has been studied by a number of researchers for a variety of precursor gases: oxygen, argon, nitrogen, hydrogen, ammonia, sulfur dioxide, hydrocarbons, fluorocarbons, halogens, organo-silanes and halo-silane derivatives. In this section we will discuss how these different precursors along with the other operational parameters have been used to impart various surface properties to cellulose fibers.

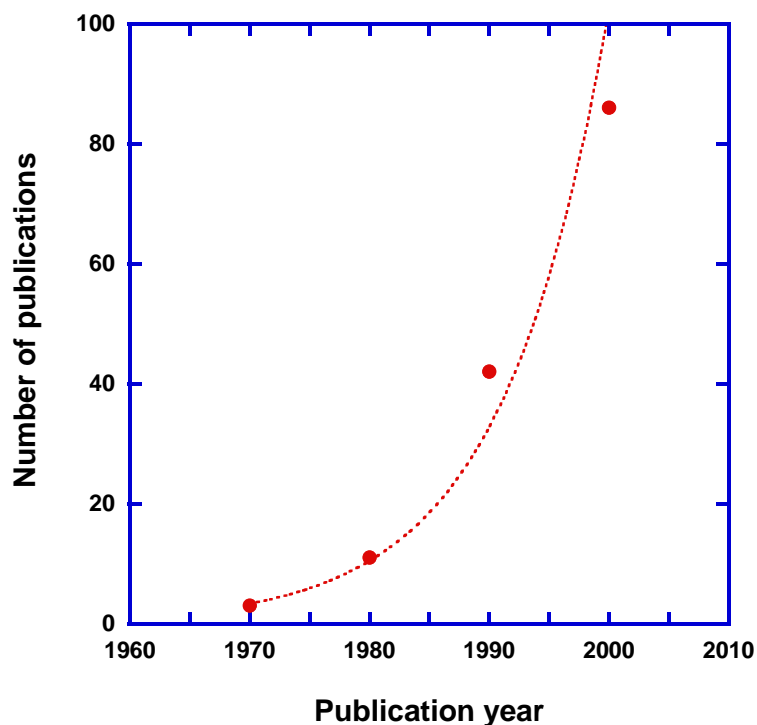


Figure 1.4 Plot of number of publications versus publication year (accumulated for every 10 years) obtained from ISI Web of Science search for terms ‘plasma’ and ‘cellulose’

1.1.4.1 Hydrophilicity

Although hydrophilicity and hydrophobicity are defined as displaying a water contact angle less than and greater than 90° respectively, these terms should be considered only in a relative sense. Therefore, in the following sections, we will refer to modification of paper to impart hydrophilic or hydrophobic properties will indicate simply that the surface energy of the paper has been increased or decreased, respectively.

As mentioned previously, pure cellulose is hydrophilic in nature with a water contact angle of $\sim 20^\circ$. But, depending upon the pulping techniques used, the pulp will contain different proportions of cellulose lignin, hemicellulose and extractives which may affect the wettability of the final paper. Controlled hydrophilicity of the paper is critical because

this property determines dye uptake and printability efficiency. Plasma treatments of paper have been effective in altering the wettability of cellulose fibers. Carlsson et al. [35, 36] studied the effects of oxygen plasma treatment on three kraft pulps with different proportions of lignins and extractives. They reported that the oxygen plasma oxidizes the lignin and extractives present on the fiber surface, resulting in increased hydrophilicity and water sorption characteristics. They also studied the effects of a hydrogen plasma treatment on a resin-rich paper. For a specific reaction time, hydrogen plasma sputters the light molecular weight resins from the fiber surface leading to an increase in the O/C ratio, which resulted in enhanced wettability. In another study, the same researchers[37] reported that an oxygen plasma treatment increased the hydrophilicity of Chemical Thermo Mechanical Pulp (CTMP) paper that had high lignin content. [37]

In some cases internal sizing is added during one of the steps in the paper making process, which creates the need to make the paper hydrophilic during post-treatment for specific applications. Vesel et al. [38] reported that ink jet paper, which is moderately hydrophobic due to the addition of alkyl ketene dimer (AKD) during the paper making process, can be rendered hydrophilic using an oxygen plasma treatment. Denes et al. [39] studied the effects of both SiCl_4 and oxygen plasmas on sized security paper. Both SiCl_4 and oxygen plasmas increased the oxygen content on the surface and thus decreased the contact angles of water on these substrates.

1.1.4.2 Hydrophobicity

Plasma processing techniques have been successfully implemented to make paper surfaces hydrophobic by depositing hydrophobic polymer films. The advantage of plasma assisted deposition is that with sufficiently thin polymer film deposition, the coating covers only the individual fibers and thus does not block the intrafiber pores.

In the 1980s, many researchers attempted to fluorinate various polymer surfaces using fluorocarbon plasmas. Along this line, Sapiuha [27, 40] first attempted to fluorinate paper surfaces using tetrafluoromethane plasmas to yield hydrophobic properties. In this work, paper surfaces were treated with different ratios of CF_4 and O_2 feed gases and hydrophobicity was measured in terms of absorption time of water on the treated paper surfaces. The water absorption time increased with increasing CF_4 fraction in the feed gas. The increase in hydrophobicity was attributed to the incorporation of CF_2 groups on the fiber surfaces. Adding to this contribution, Sahin et al. [41] reported that the atomic fluorine concentration on the paper surface can be controlled by the external plasma parameters such as power, pressure and reaction time.

Surface properties of a paper substrate differ on its two sides depending on whether it is on the felt or wire side during the paper formation process. This leads to a difference in the surface morphology on the micro-scale between the two sides of the paper, often referred to as “two-sidedness”. Sahin [42] studied the effect of this “two-sidedness” on the fluorination of papers by CF_4 plasmas. He reported that the felt side which has coarse

fibrous materials usually is more prone to fluorination than is the wire side containing fiber fines.

Mukhopadhyay et al. [43] carried out an extensive study on the treatment of paper fibers with a variety of fluorocarbon precursors: penta-deca-fluoro-octyl-acrylate, per-fluoro-dodecene, tri-deca-fluoro-octene and per-fluoro-methyl-cyclo-hexane. They concluded that all these precursors resulted in the incorporation of the CF_2 and CF_3 groups which are key to obtaining hydrophobic properties. They also concluded that a minimum of 15% surface fluorine concentration is sufficient to impart hydrophobicity on the paper surface, irrespective of the monomers used. The chemical composition of the paper fibers, which is determined by the paper making parameters, does not have a significant impact on plasma fluorination. Soon after Sahin et al. [41] reported the hydrophobization of bleached kraft paper, Navarro et al. [44] reported the hydrophobization of chemothermomechanical pulp (CTMP) by treating it in a fluorotrimethylsilane (FTMS) plasma environment. The hydrophobicity of the treated surfaces was attributed to the formation of C-Si-F and C-O-Si-F groups. The same researchers later used a pentafluorophenyldimethylsilane (PFPDMS) precursor and were successful in replicating the hydrophobic properties on the CTMP substrate [45].

Vaswani et al. [46] reported the deposition of fluoropolymer thin films from pentafluoroethane (CF_3CHF_2) and octafluorocyclobutane (C_4F_8) precursors onto paper surfaces. The treatment made the substrates water repellent (hydrophobic), but allowed water vapor to diffuse through the film and substrate readily. The researchers reported

that for cellulose surfaces, a minimum of a 70 nm thick fluorocarbon film was required to obtain a uniform fiber coating that completely retarded the absorption of water into the paper. They also observed that treated paper showed good resistance against the permeation of oleic acid and grease after the plasma treatments (lipophobicity).

An interesting aspect of the fluorocarbon precursors is that they simultaneously provide polymerization and etching environments when excited in the plasma environment.

While the CF_x species are responsible for polymerization, fluorine free radicals can etch the paper surfaces. The rate of etching and deposition can be controlled by adjusting the ratio of fluorine to carbon in the precursor and by tuning various reactor parameters such as power, pressure, substrate temperature and residence time. It is worth mentioning that if the critical experimental parameters were not properly selected, the plasma treatments mentioned above would not have resulted in a fluorocarbon film on the paper surfaces; rather these plasmas may have etched the paper surfaces.

There is a general concern within the scientific community regarding the incorporation of fluorine containing compounds into the hydrophobic coatings; these coatings may make the product less environmental friendly [47]. Food packing is one of the most sensitive applications of plasma fluorinated paper coatings. The Food and Drug Administration (FDA) regulations indicate that fluorine can be used in a water or oil repellant material with a basis weight of 0.22-2.44 g/m² depending upon the chemical form in which it is present. The plasma deposited fluorocarbon films are often referred to as Teflon-like films because of their chemical and physical properties being similar to those of PTFE.

Hence, it is reasonable to assume a density of ~ 2 g/cc (typical for PTFE film) for the plasma deposited fluorocarbon films. Studies have suggested that a ~ 100 nm fluorocarbon film is sufficient to inhibit liquid water absorption. For this film thickness (~ 100 nm) the basis weight corresponds to ~ 0.2 g/ m² which falls well within the specified limits. In spite of this compliance with FDA regulations, there has been an increased interest in the use of non-fluorinated precursors for rendering paper surfaces hydrophobic. Tu et al. [48] first reported the use of cyclohexane, a hydrocarbon precursor, to yield hydrophobicity on cellulose fibers. However, these coatings were not conformal on the fibers and formed bridges between the fibers that altered the porosity of the treated paper [47]. Organic silicone compounds were also investigated as precursors to obtain conformal fluorine-less hydrophobic films. Organic silicone compounds have stronger adhesion to cellulose surfaces than do hydrocarbon moieties because of the formation Si-O bonds at the interface. Tan et al. [47] and Mahlberg et al. [49] reported the hydrophobization of paper surfaces using hexamethyldisilazane (HDMS) precursors which yielded conformal film coatings. Tan et al. [47] also proposed a dual film deposition method in which a double layer film was obtained by first depositing HDMS, followed by a deposition of n-hexane. The double layer was reported to have good mechanical stability because of the presence of HDMS film at the interface (between n-hexane and cellulose) and also cost effective because of using n-hexane. With the increasing interest in green technologies, Gaiolas et al. [29] very recently reported the hydrophobization of paper using renewable, natural precursors such as limonene and myrcene.

Apart from using polymerizable precursors to decrease the surface free energy of paper, non-polymer forming gases have also been used as precursors to functionalize paper surfaces. Carlsson et al.[35] reported that hydrogen plasma treatment of paper surfaces reduced the hydroxyl functional groups of the cellulose and created low molecular weight, low polarity functional groups, which resulted in an increase in the time needed for water absorption. Deslandes et al. [50] studied the effects of nitrogen plasma treatments on paper surfaces and found that the paper became more hydrophobic, as indicated by increased water absorption times. However, in this case the change in the surface energy was attributed to plasma induced degradation of the cellulose molecules and not to the incorporation of nitrogen containing functional groups on the cellulose surfaces.

1.1.4.3 Adhesion

The adhesion between cellulose and other polymers is critical for the use of cellulose fibers as reinforcing elements in macromolecular composites. As discussed previously, cellulose fibers have a non-polar surface due to the presence of OH groups. However, most polymer matrices for making reinforced composites are hydrophobic, which results in a low interfacial compatibility with cellulose fibers. Furthermore, the loss of mechanical properties of cellulose fibers when exposed to moisture is also not desirable in these applications [51]. Hence, cellulose fibers need to be surface modified to retard moisture absorption and increase their interfacial adhesion with the polymer matrices. Plasma processing of fibers can increase the adhesion in the following ways [52]: 1) by introducing polar groups and radical sites to enhance bonding with the matrix, 2) by

forming a new polymer layer which can form covalent bonds with the matrix and 3) by creating roughness on the fiber surface to increase mechanical interlocks between the fiber and the matrix.

Wertheim et al. [34] first reported improvement of adhesion between cellulose and polyethylene using a microwave discharge. Zanini et al. [53] studied the effects of argon plasma treatments on lignocellulosic fibers in an attempt to create radical sites on the fiber surfaces for improved fiber adhesion. Results from electron paramagnetic resonance (EPR) spectroscopy confirmed the presence of phenoxy radicals on the fiber surface after the plasma treatment. Morales et al. [54] extended this idea by creating radical sites on a thin film of polymer deposited on the fiber surface. They synthesized radical rich polystyrene-like films on the cellulose surface which provided a strong interface between the fiber and the polystyrene matrix, resulting in good fiber-matrix adhesion. Ahlblad et al. [55] reported the grafting of cellulose fibers with butadiene or divinylbenzene by plasma treatment which increased the adhesion between the cellulose and polymer.

Young et al. [56] provided new insight regarding the improvement of adhesion between cellulose fibers and polymers. A cyclohexane plasma was chosen in an attempt to make the cellulose fiber surface rich in non-polar functional groups to enhance adhesion with a non-polar polypropylene surface. Counter-intuitively, T-peel strength (hot pressing two substrates together and peeling in a direction orthogonal to the substrate's surface; the structure looks like a "T" while peeling the substrates) results showed that the cyclohexane plasma treatment did not significantly enhance the adhesion. On the other

hand, when the polypropylene was treated with an oxygen and argon plasma, the T-peel strength increased significantly. From these results, Young et al. concluded that making the cellulose surface non-polar is not sufficient to increase the adhesion strength; however, enhancing the surface polarity can increase the adhesion strength. Possible reasons for such observations were not discussed. Carlsson and Strom [36, 59] reported that reduction of a cellulose surface in a hydrogen plasma environment, which decreases the number of polar surface groups, does not have an effect on the adhesion of cellulose to polyethylene. Additionally, they reported a decrease in adhesion that was attributed to a weak boundary layer formed between the low molecular weight degradation products of cellulose and polymer. They also reported that an oxygen plasma treatment of cellulose enhances the cellulose-polymer adhesion because the oxygen plasma removed low molecular weight resin-compounds on fibers which enhanced adhesion. Finally, some researchers have reported an enhancement of adhesion with non-polar compounds by adding polar groups to the cellulose surfaces; these observations are in opposition to the concepts described in the beginning of this section. Interestingly, this counter-intuitive conclusion has received little opposition to date.

1.1.4.4 Bioactivity

Recently, there has been an increased interest in the fabrication of bioactive paper substrates [16]. Anti-bacterial paper finds potential applications in cost-effective lab accessories and medical textiles. Jampala et al. [57] reported the synthesis of quaternary ammonium groups on cellulose surfaces by using plasma processing techniques to obtain bactericidal properties. They first treated the paper surfaces with an ethylenediamine

plasma to create secondary and tertiary amines on the surface, which were then treated with hexyl bromide and methyl iodide to yield immobilized quaternary ammonium groups. The treated paper surface inactivated 98.7% and 96.8% of *S. aureus* and *K. pneumoniae*, respectively. Vasile et al. [58] reported the surface modification of cellulose fibers using a variety of carboxylic acids to obtain anti-bacterial and anti-allergic properties via the formation of cellulose esters.

Plasma modified cellulose surfaces also find application as a biomaterial (e.g., hemodialysis membrane). Poncin-Epaillard [59] proposed that treatment of cellulose surfaces with CF₄ plasmas reduced the bio-activation for use in hemodialysis membranes by decreasing surface OH groups on the cellulose surfaces.

Plasma treatments have also been used for the removal of microbial contaminations from paper substrates that contained printed cultural relics, in an attempt to protect the cultural heritage available in books and paintings. In 2001, Vohrer et al. [60] reported the removal of *Bacillus subtilis* and *Aspergillus niger* contamination on naturally aged ground paper from the 1950s using oxygen and hydrogen precursors.

1.2 Motivation and objectives

Many living organisms exhibit unusual wettability behavior that causes water drops to be repelled and roll off the surface; examples are plants [61, 62] (lotus leaves, cabbage, Indian cress, etc) and insects wings [61-63] (butterflies, cicada, etc.). There has been a significant interest in fabricating such water repellant surfaces by artificial means since the late 1930s. A surface is generally classified as superhydrophobic (SH) if it has a static water contact angle (CA) greater than 150° [61, 63, 64].

When the free energy of a homogeneous, flat surface is lowered through modifications of its chemical composition, the water CA increases [63, 64]. However, the CA values of chemically modified flat surfaces are bound by certain limits: the highest water CA recorded on a flat surface was 120° , with a free energy of 6.7 mJ/m^2 [63-65]. The water CA can be further enhanced only by creating some type of roughness or patterns on the hydrophobic surface.

Wenzel [66, 67] proposed a model (Figure 1.5a) describing the increased CA due to surface roughness as follows:

$$\cos \theta^* = r \cos \theta \quad (1.1)$$

where θ^* and θ are the macroscopic contact angles of the rough substrate and of a flat substrate of the same chemical composition, respectively; r is the roughness factor which is defined as the ratio of the actual area of the rough surface to the geometrically

projected surface area. The apparent interfacial energy on a rough surface differs from the true interfacial energy, because the area of the liquid-solid interface is increased (proportional to r) on the rough surface. As a result, the presence of roughness amplifies the actual contact angle: a hydrophobic surface ($CA > 90^\circ$) becomes more hydrophobic, and a hydrophilic surface ($CA < 90^\circ$) becomes more hydrophilic. In some cases, the liquid does not completely wet the roughened surface; this scenario cannot be addressed by Wenzel's model, but is explained by Cassie's model (Figure 1.5b) [68]:

$$\cos \theta^* = f \cos \theta + (1 - f) \cos 180^\circ = f \cos \theta - (1 - f) \quad (1.2)$$

where f and $(1-f)$ are the fractions of solid/water and air/water interface between the drop and the substrate, respectively.

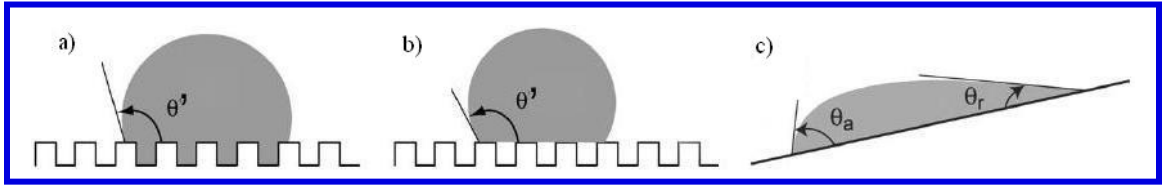


Figure 1.5 a) Wenzel state CA, b) Cassie state CA, c) Definition of CA hysteresis, $\Delta\theta = \theta_A - \theta_R$ [62]

Although the sliding of a water drop is typically visualized as the movement of the water molecules at the solid-liquid interface, only the interfacial (solid-liquid) water molecules at the three-phase contact line move. All other interfacial (solid-liquid) water molecules remain stationary, thus obeying the no-slip boundary condition of fluid mechanics [69, 70]. When a drop slides on a surface, the contact line on the front edge advances over the surface and the contact line on the trailing edge recedes over the surface. Any kinetic energy barrier to the movement of the receding contact lines results in a difference in contact line dynamics for the two drop edges, which is referred to as hysteresis [71]. Such kinetic energy barriers can be caused by physical heterogeneity (roughness) and/or chemical heterogeneity. Figure 1.5c describes the definition of CA hysteresis. This value is a measure of the stickiness of the water drop to the surface. A lower value of CA hysteresis implies that the water drop can easily slide off the surface. Johnson and Dettre [72-74] pointed out that the hysteresis is a balance between energy barrier heights of between metastable states and the vibrational energy of the drop. They formulated a mathematical model in which the height of the energy barriers was changed to obtain different hysteresis values. Nevertheless, they could not use this model to quantitatively predict the hysteresis on their real surfaces and gave only qualitative explanations. Recently, Chen et al.[75] pointed out that the energy barrier can be attributed to the line tension of the contact line at the three-phase interface. To date, kinetic energy barriers for movement of the contact line (causing the hysteresis) and/or the line tension has not been modeled mathematically; thus, prediction of the hysteresis for a drop sliding on a surface has not been possible.

Superhydrophobic paper opens a new window of applications for paper substrates as a substitute for polymers in textile, chemical and biomedical industries. As described in the previous sections, two hundred years of surface modification of paper (including plasma processing) have resulted mainly in turning hydrophilic paper hydrophobic; at the start of this project in 2006, fabrication of superhydrophobic paper had not been reported. The main reason is that research efforts had concentrated on changing the surface energy of paper by changing the surface chemistry; the surface roughness aspect which is critical in achieving surface superhydrophobicity, had been neglected.

Contemporary to this thesis work a couple of other research groups also started focusing on making paper superhydrophobic: in one study, in-situ atom transfer radical polymerization (ATRP) was used to graft and grow a polymer layer with nanoscale roughness on a fiber surface, which led to superhydrophobicity (water contact angles $>160^\circ$) [76]. In another study, polymethylsilsesquioxane coatings were covalently bonded onto the cellulose fiber surfaces via a polycondensation reaction between the hydroxyl groups of silanol and cellulose. Deposition of this low surface energy film also resulted in the formation of nano-scale roughness resulting in a superhydrophobic paper surface with a contact angle of 157° [77]. However, because the nanoscale roughness of the surfaces in the above two methods is associated with a deformable polymer film, robustness and durability are limited. Superhydrophobicity was also imparted onto paper surfaces by depositing nano-scale silica particles on the fibers using sol-gel approaches and fluorinating the resulting surface using perfluorooctyltriethoxysilane to obtain low surface energy [78]. In another study, fine micron-sized particles of alkyl ketene dimer

(AKD) were formed by rapid expansion of AKD solutions in supercritical carbon dioxide. The deposited AKD layer had the appropriate roughness scale for superhydrophobicity and yielded a contact angle of 173° [79].

In all the above referenced work, surface roughness was obtained by depositing an additional material onto paper. This raises concerns about the adhesion of those materials to the cellulose fiber surfaces and the mechanical robustness. It would be advantageous to change the surface topography of cellulose fibers inherently, instead of through deposition of a foreign material. This thesis project aimed to address this issue by uncovering the inherent roughness scales present on fiber surfaces by plasma etching and by depositing a thin fluorocarbon film via plasma polymerization. Thus, there would be no need to deposit foreign materials to obtain roughness. Also, the above mentioned methods (reported by other researchers) used a solvent based approach in at least one of the fabrication steps; this approach is not ideal because of safety and environmental considerations [80]. The current project uses a more environmentally friendly plasma processing (vapor-based) approach which has been recently included in the list of “green” technologies [29].

In principle, different levels of adhesion (‘stickiness’) can be achieved on a superhydrophobic surface depending on the sizes and shapes of the roughness scales. This tunability has not been reported on any superhydrophobic surface to date, including superhydrophobic paper. The current thesis project investigated approaches to obtaining

tunable ‘stickiness’ on superhydrophobic paper surfaces by manipulating plasma processing parameters.

Out of approximately 6.1 billion people in the world, one-third of them do not have access to electricity,[81] nearly one-fourth live below the poverty line of \$1.25/day[82] and one-sixth are not educated enough to read or write their names.[83] In spite of this, the focus of biomedical research is primarily on industrialized nations which is notoriously referred as the ‘10/90 gap’.[84] Considering these facts, lab-on-chip (LOC) biomedical diagnostic devices targeted for developing countries must be made as simple and affordable as possible in order for them to be handled by less skilled technicians in resource limited point-of-care (POC)[81] environments. Recently, paper-based lab-on-Chip (LOC) devices (also referred to as lab-on-paper (LOP)[16]) have emerged as a promising alternative technology for the commercially available LOC microfluidic devices. In a recent report on the top ten biotechnologies for improving health in developing countries “modified molecular technologies for affordable, simple diagnosis of infectious diseases” was ranked as the number one priority [85, 86]. Another report on the grand challenges for global health ranked the development of technologies to “measure disease and health status accurately and economically in poor countries” first among the top 14 priorities [86, 87]. Due to their affordability and potentially simple fabrication technology, LOP devices may offer improved global availability of medical technology. This thesis project aimed to develop a two-dimensional LOP device capable of storage, transfer, mixing and sampling of liquid drops by decorating superhydrophobic paper substrates with high surface energy ink patterns.

With the brief introduction of the goals described above, the specific objectives of the thesis project were:

- a) To fabricate superhydrophobic paper surfaces *via* plasma processing
- b) To obtain tunability of adhesion of water drops on superhydrophobic paper surfaces by controlling the etching and deposition times
- c) To study the effect of fiber types and paper making parameters in rendering paper superhydrophobic
- d) To pattern superhydrophobic paper with phaser printer inks to control the mobility and adhesion of micro-liter drops and thus obtain a two-dimensional droplet manipulation device
- e) To optimize the patterning geometries to control the directional mobility and adhesion of water drops on a two-dimensional droplet manipulation device

1.3 Thesis organization

Chapter 2 (Objective a) describes the appropriate plasma processing conditions to impart superhydrophobic property to paper surfaces. The characterization of superhydrophobic paper surfaces via SEM, ATR-FTIR, XPS and contact angle goniometry is presented.

Chapter 3 (Objective b) describes methods to obtain tunable adhesion of water drops on superhydrophobic paper surfaces. The fundamental wettability mechanisms that control the tunability of the stickiness are discussed in detail. *Chapter 4* (Objective c) discusses the effect of fiber source and other paper making parameters on the ability to make paper superhydrophobic using plasma processing and this work provides additional insight that is useful for designing superhydrophobic paper. *Chapter 5* (Objective d) describes the

methods to pattern superhydrophobic paper surfaces with high surface energy islands to obtain tunability in adhesion forces for water drops. The adhesion forces with respect to the island sizes and geometries are extensively studied and the mobility of water drops on these patterned substrates is modeled using classical adhesion force expressions. The fabrication of a two-dimensional lab-on-paper device using the patterned superhydrophobic papers is described. *Chapter 6* (Objective e) discusses the optimization of patterning geometries on superhydrophobic paper surfaces in order to maximize control over the directional mobility and adhesion of drops. *Chapter 7* presents the overall summary and conclusions of the project, as well as recommendations for future research.

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CHAPTER 2

FABRICATION OF “ROLL-OFF” AND “STICKY” SUPERHYDROPHOBIC CELLULOSE SURFACES VIA PLASMA PROCESSING*

2.1 Introduction

As mentioned in the previous chapter, it has long been recognized that superhydrophobic surfaces [1] (water contact angle $>150^\circ$) require a unique combination of two fundamental properties: 1) surface roughness and 2) low surface energy. To date, superhydrophobic surfaces have been fabricated with a variety of length scales, *e.g.* single length scale (nanometer range [2-4] or micrometer range [5-7]) or hierarchical combination of length scales (micrometer-micrometer [8] or micrometer-nanometer [9-11]) with various topographies. Since the late 1930s, significant interest has existed in designing water repellent surfaces by artificially generating these two properties on a variety of materials. From a processing viewpoint, the desired combination can be obtained via several routes: 1) add roughness to an inherently low surface energy material; 2) add roughness to a hydrophilic surface and then modify it with a hydrophobic surface treatment; or 3) modify a surface with a low surface energy material which adds

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1) Balu, B., Breedveld, V., and Hess, D.W., *Fabrication of "roll-off" and "sticky" superhydrophobic cellulose surfaces via plasma processing*. Langmuir, 2008. **24**(9): p. 4785-4790.
2) Balu, B., Kim, J.S., Breedveld, V., and Hess, D.W., *Tunability of the Adhesion of Water Drops on a Superhydrophobic Paper Surface Via Selective Plasma Etching*. Journal of Adhesion Science & Technology, 2009. **23**: p. 361-380.

inherent roughness, for example the deposition of hydrophobic nanoparticles. Surface roughness at nano- and micro-meter scales has been obtained through a variety of methods: controlled crystallization [12], plasma etching [13], laser etching [14], plasma enhanced chemical vapor deposition [15], lithography [16], electro-spinning and spraying [17], sol-gel processing [18], stretching [19], nanocasting [20], and graft-on-graft polymerization [21]. Similarly, there are numerous ways to modify the surface chemistry, such as sol-gel processing [18], graft polymerization [21], electrochemical deposition [22], plasma enhanced chemical vapor deposition [23], chemical vapor deposition [13], and atomic layer deposition [24]. The selection of the appropriate method to create roughness and/or low surface energy depends on the mechanical and physiochemical properties of the substrate.

Artificial superhydrophobic surfaces have been fabricated on a variety of organic and inorganic substrates, *e.g.* polymers [8, 13, 25-27], Si wafers [5, 9, 11], glass slides [10, 28], metals [29, 30]. With the increased environmental interest in the use of renewable resources and biodegradable materials, the inorganic substrates listed above are less than ideal. Moreover, they are not mechanically flexible, which limits their processability and therefore the range of potential applications. Organic polymer substrates, on the other hand, are flexible, but they tend to be fairly expensive and often lack biodegradability and renewability. Hence, the search for alternative substrates for superhydrophobic surfaces is important. Cellulose, a biodegradable, renewable, inexpensive, biopolymer, which is abundantly present in nature, has been targeted as a candidate. If cellulose-based paper substrates can be rendered superhydrophobic by simple processing schemes, they will

offer a promising alternative to conventional superhydrophobic substrates. Due to relatively low cost and mechanical flexibility, these surface-modified materials could have applications in a vast array of products, including fast food and microwavable food packages, beverage containers, self-cleaning cartons, labels, paper boards, textiles, heat transfer surfaces (to remove condensed water quickly), microfluidic devices and membranes with low degrees of surface fouling.

Paper surfaces have micrometer roughness due to the network of overlapping fibers, but lack the smaller nanoscale roughness that is usually associated with superhydrophobicity. As mentioned in section 1.2 (Chapter 1), it would be advantageous to change the surface topography of cellulose fibers inherently, instead of achieving roughness via a deposited surface layer. To proceed in this direction, it is essential to understand the internal structure of cellulose fibers; considerable information has been published on this topic [31-37]. As shown in Figure 1.2 (Chapter 1) each cellulose fiber consists of several microfibril-bundles, which in turn are composed of microfibrils that have diameters ranging from 3 to 30 nm [31, 32, 34, 36]. In combination with the already mentioned micrometer length scale of the fiber network, the microfibrils could therefore provide the necessary additional length scale for superhydrophobicity. Because the surface of native cellulose fibers is relatively smooth, the challenge is to find a method to roughen the surface by exposing the microfibrils. One key characteristic of the microfibrils is that they contain mostly crystalline cellulose moieties, while the matrix surrounding the microfibrils is predominantly amorphous in nature [31-33, 38]. Taking advantage of this inherent difference in material properties may offer a viable approach to generate

superhydrophobicity on cellulose surfaces. Unfortunately, cellulose fibers represent a relatively complex substrate and their properties, in particular their water absorbency, stretchability, compressibility, porosity, non-uniform surface chemistry and thermal degradability, inhibit the use of most of the commonly used techniques for creating smaller characteristic length scales (see above discussion).

Domain-selective etching has been used to create surface roughness on polymer substrates with crystalline and amorphous domains [13, 27, 39, 40]. In these studies, etching was performed by a vapor phase plasma process at low temperature, under solvent-free conditions. Selectivity is based on the premise that amorphous domains of polymer substrates will be more susceptible to reactive plasma etching than crystalline domains. Etching thus preferentially erodes the amorphous domains, leaving behind the crystalline domains. This process ultimately results in a roughened polymer surface with a characteristic length scale that is determined by the size of and distance between crystalline domains. Because of the difference in crystallinity between microfibrils and surrounding matrix, the hypothesis is that selective etching should also be possible on cellulose fibers, provided that the appropriate etching conditions are generated. Previous studies applied plasma-assisted etching of paper surfaces to investigate interfiber bonding, internal structures of fibers, and coating distribution within the paper [41]. The focus of these studies was to establish plasma processing conditions that were aggressive (using CF_4 and O_2 gases), so that the etching process served as a fine microtome to generate a z-directional cross-section of the paper for analysis of various properties with respect to depth. This Chapter focuses on surface modification (altering only the top

layers of fibers) while maintaining bulk paper properties. As a result, this investigation invoked considerably less aggressive plasma etching conditions to achieve selective etching of the surface layers of cellulose fibers in order to generate appropriate roughness scales for superhydrophobicity. However, as discussed previously, roughness is not sufficient to establish superhydrophobicity. The inherently hydrophilic cellulose fibers must also be hydrophobized. The hydrophobization was imparted by plasma-enhanced chemical vapor deposition (PECVD) of a fluorocarbon coating.

In this chapter we demonstrate that superhydrophobic properties can be created on paper surfaces via a combination of selective etching by an oxygen plasma and deposition of a fluorocarbon film (~100 nm on Si wafer) from pentafluoroethane (PFE) monomer via plasma-enhanced chemical vapor deposition (PECVD). Because both steps alter primarily the surface properties of cellulose fibers, it is expected that the bulk material properties remain unchanged. Also, because the approach relies on uncovering roughness present inherently on cellulose fibers, the roughness obtained by this approach should be more mechanically robust than roughness generated from foreign materials that have been used previously to impart superhydrophobicity to paper.

2.2 Experimental

2.2.1 Handsheet formation

Commercial paper contains filler particles and other additives to enhance its physical and optical properties[42]. In order to initially avoid the interaction of these particles with the surface modification techniques employed, control samples were used, generally referred

to as handsheets, which were prepared using southern hardwood kraft (Alabama River Pulp Co.) and southern softwood kraft (North Carolina International Paper). The handsheets were fabricated to ensure that the most relevant properties of the handsheets were in the same range as those of commercial copy paper: basis weight ($66.7 \pm 2.7 \text{ g/m}^2$), thickness ($105.6 \pm 4.9 \text{ }\mu\text{m}$) and roughness (amplitude $R_a = 6.17 \pm 0.24 \text{ }\mu\text{m}$). Handsheets were prepared using TAPPI standardized method T205 sp-02: dry sheets of soft and hard wood fibers (1:1 mass ratio) were soaked overnight, beaten to a pulp in a valley beater and diluted to appropriate consistency; the pulp was poured into a mesh-bottom mold and the water allowed to drain under gravity to form a handsheet, which was then pressed between blotter paper sheets and further dried in a hot press. The handsheets were then placed in a humidity controlled environment (50% RH) for over- night drying.

2.2.2 Plasma processing

A 6-inch parallel plate plasma reactor was used for the plasma processing (Figure 2.1). The stainless steel bottom electrode was grounded and heated to $110 \text{ }^\circ\text{C}$ using Omegalux CIR 2015 cartridge heaters (Omega Engineering Inc., Stamford, CT). The temperature at the bottom electrode was monitored using a type K thermocouple controlled by a Syskon RKC temperature controller (RKC Instrument Inc., Southbend, IN). The stainless steel top electrode of the reactor was connected to a HF-300 13.56 MHz, RF power supply (ENI Power Systems, Rochester, NY). To minimize reflected power in the plasma reactor a matching network (Heathkit SA-2060A, Heath Company, Benton Harbor, MI) was placed between the top electrode and the power supply. The reactor pressure was

monitored and maintained using a pressure gauge (Varian Inc., Lexington, MA) and an Alcatel 2063 C rotary vacuum pump (Alcatel, Annecy, France).

Figure 2.1 Diagram of 6 inch parallel plate plasma reactor setup for plasma processing of paper substrates (monomer was oxygen and pentafluoroethane (PFE) for etching and deposition respectively)

with cellulose (P) to form water vapor, CO, and CO₂, thereby removing material from the surface according to the following reactions [41]:



Step-2, Deposition. A thin fluorocarbon film was then deposited onto the substrates. The deposition gas mixture consisted of a precursor gas (pentafluoroethane [PFE] flowing at 20 sccm), and a carrier gas (argon flowing at 75 sccm). After the reactor reached a stable (steady state) pressure (1 Torr), an RF power of 120 Watts was applied to the top electrode for 2 minutes. Electron impact collisions with the precursor form various C_xF_yH_z moieties, which react primarily at the substrate surface to form an adherent cross-linked fluorocarbon film [43-45]. At the end of the deposition, the plasma power was turned off and the reactor evacuated to base pressure. Finally, the reactor was repressurized to atmospheric pressure by backfilling with N₂ gas and the sample was removed from the reactor for surface characterization. Under these conditions, the thickness of the deposited fluorocarbon film on a silicon wafer, which can readily be measured via ellipsometry, was ~100 nm. The full process, including step-1 (etching) and step-2 (deposition) will be referred to as *superhydrophobic (SH)-treatment* in this report. In addition to this treatment, some samples were exposed only to step-2 (fluorocarbon deposition), thereby eliminating step-1 entirely; this treatment is termed *PFE-treatment*.

Pentafluoroethane (PFE) monomer gas (N4 grade, 99.99%) was donated by Dupont (Wilmington, DL). Argon carrier gas (Ultra High Purity, 99.99%) was purchased from Air Products and Chemicals Inc. (Allentown, PA). Nitrogen (Ultra High Purity, 99.999%) and oxygen (Ultra Pure Carrier, 99.996%) were purchased from Airgas Inc. (Radnor, PA).

2.2.3 SEM

SEM measurements were obtained with a LEO scanning electron microscope (model 1530, Carl Zeiss SMT Inc., Peabody, MA) operated at a pressure of $\sim 1.0 \times 10^{-7}$ Torr at room temperature. The operating voltage was adjusted between 5 and 10 kV depending on the magnifications used to avoid damaging the paper samples. Since both the paper and PFE film are insulators, the substrates were sputtercoated (EMS 350; Electron Microscopy Sciences, Hatfield, PA) with a thin gold film (~ 15 nm) prior to SEM imaging.

2.2.4 X-Ray photoelectron spectroscopy (XPS)

Spectra were collected using a PHI model 1600 spectrometer (Physical Instruments, Inc., Chanhassen, MN) with Al K α X-rays operated at a power of 350 W. The operational pressure was typically held below 5×10^{-9} Torr. Further details of this equipment can be found elsewhere [45].

2.2.5 ATR-FTIR

FTIR spectra were obtained with a Bruker Equinox 55 FT-IR spectrometer (Bruker Optics Inc., Billerica, MA), equipped with a nitrogen cooled MCT detector. Further details of this equipment can be found elsewhere [46]. A thin strip of the substrate (0.76 cm by 6.35 cm) was cut and pressed against a ZnSe ATR crystal using a flat metal strip. The pressure on the sample was controlled by a set screw and kept approximately constant for all the samples.

2.2.6 FIB

The focused ion beam (FIB) milling steps and SEM images were obtained using Quanta 200 3D and Nova 200 NanoLab Dual Beam FIB/SEM systems (FEI Company, Hillsboro, OR).

2.2.7 Contact angle goniometer

Water contact angle measurements were performed with a Rame-Hart contact angle goniometer (model 100, Netcong, NJ). The substrate was placed on a translation stage, a 4 μ L drop suspended from a needle was brought into contact, and the stage was moved in the x -direction (left to right) [47]. The CA on the advancing and receding sides of the drop was measured. The advantage of this dynamic method is that it scans a large substrate area and thus yields better statistically averaged values of CA, especially for heterogeneous substrates like paper. One of the disadvantages of this method is that receding CA values of less than 10° cannot be measured, because the drop can break apart while being dragged on the surface [47]. For the paper substrates which had

undergone PFE-treatment, the receding CA was in this range and the breakup of drops was observed. In these cases, the standard volume increment/decrement method for measuring advancing and receding CAs was used. The droplet volume was increased from 4 μL to 12 μL in increments of 4 μL to measure the advancing CA. Thereafter, small decrements of 0.17 μL were used to measure the receding CA.

2.3 Results and discussion

Before measuring the wetting properties of the modified substrates, the substrates were subjected to basic physicochemical tests to determine their physical and chemical properties after oxygen etching and PFE deposition in the plasma reactor.

2.3.1 XPS

Cellulose handsheet surfaces were characterized after etching and deposition of PFE layers by XPS analysis. As described in Chapter 1, cellulose molecules $(\text{C}_6\text{H}_{10}\text{O}_5)_n$ are essentially polymer chains of β -d-glucose residues covalently coupled via glucosidal linkages [48, 49]. Consistent with this molecular structure, the XPS survey spectrum of the cellulose handsheet (Figure 2.2a) displays two intense peaks assigned to O1s (oxygen) and C1s (carbon). Moreover, the theoretical O/C ratio of cellulose (0.83) correlates well with the ratio of 0.84 ± 0.01 determined by XPS, which indicates that the fabricated handsheets do not contain fillers or other impurities found in commercial paper, and therefore that the handsheets can be used as model substrates for cellulose fibers.

The XPS survey spectra of the etched handsheets (data not shown) did not detect the presence of additional elements, indicating that impurities, at least to the detection limit of XPS (~0.1 atomic %) were not added to the handsheet surface during plasma etching. Table 2.1 shows the %carbon, %oxygen and O/C ratio of the handsheets etched in the oxygen plasma for different etching times. The O/C ratio increases as a function of etching time, consistent with continued oxidation, surface conversion and etching of cellulose by oxygen species.

Table 2.1 Atomic percentages of carbon (C) and oxygen (O), and O/C ratio with respect to etching time

Etching time, min	%C	%O	O/C Ratio
0	54.38 ± 0.20	45.62 ± .20	0.84± 0.01
10	49.28 ± 2.27	50.72 ± 2.27	1.03± 0.10
15	48.05 ± 0.81	51.95 ± 0.81	1.08± 0.04
30	44.12 ± 1.20	55.88 ± 1.20	1.27± 0.06

An XPS survey spectrum after deposition of the fluorocarbon film on Si wafer from the PFE precursor is shown in Figure 2.2b. Due to the two intense fluorine peaks (F1s and F2s) and the reduction in the intensity of the oxygen peak, these surfaces are expected to be hydrophobic, consistent with contact angle measurements described below. High resolution C1s spectra of the PFE film revealed the highly cross-linked nature of the film (F/C ratio ~1.1) and the presence of various hydrophobic fluorinated moieties (CF_x) [44, 45].

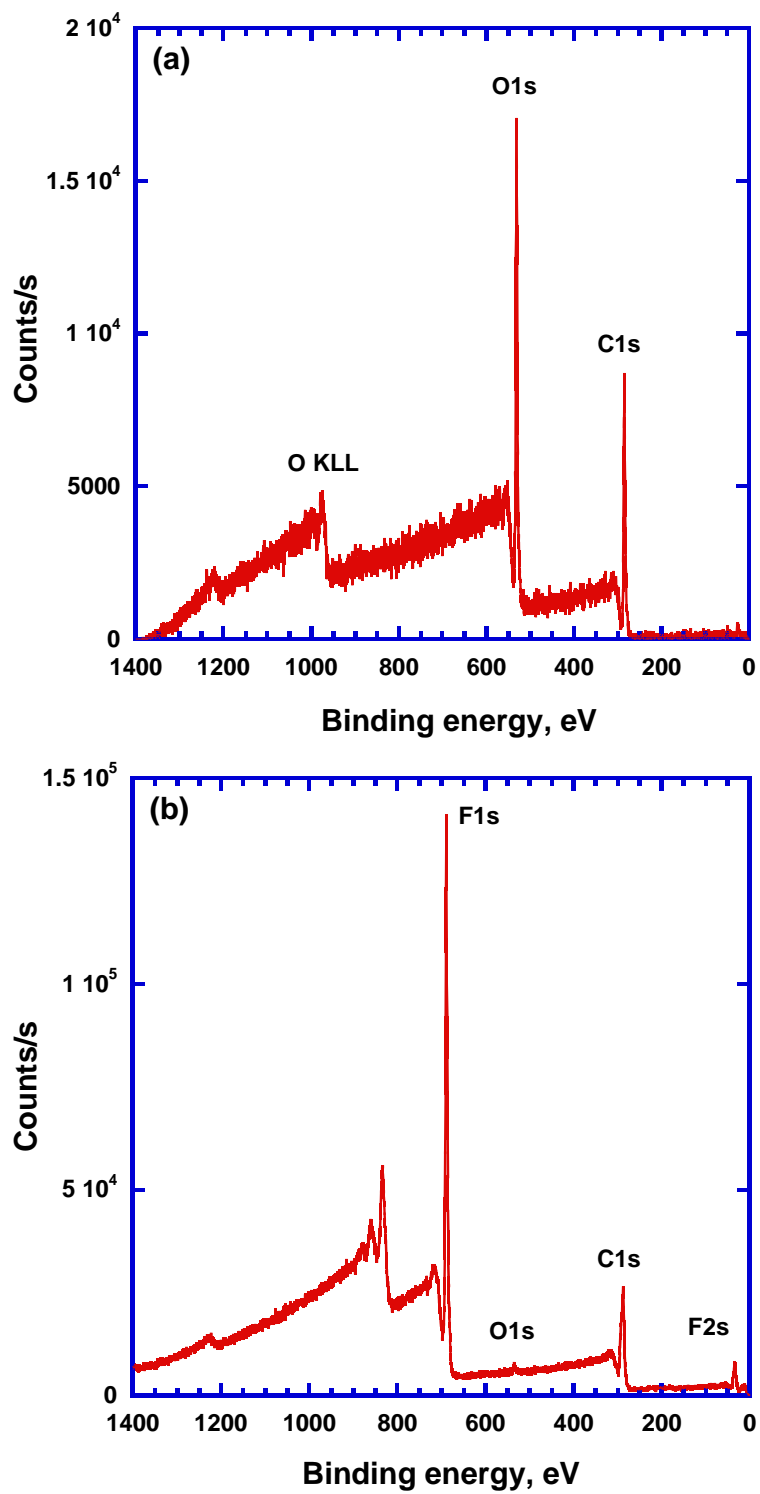


Figure 2.2 XPS survey spectra of untreated handsheet (a) and PFE film on Si wafer (b)

2.3.2 ATR-FTIR

XPS is a surface analysis technique with a probing depth of electrons ranging from 8 to 10 nm. To study the chemical nature of the bulk PFE-cellulose structure in more detail, FTIR spectra of untreated and PFE deposited handsheets were obtained with a penetration depth of 1.2 μm at 900 cm^{-1} [46]. Figures 2.3a and b show the FTIR spectra of the untreated and PFE coated handsheets (film thickness $\sim 400\text{ nm}$). The spectrum of the untreated handsheet correlates well with the FTIR spectrum of cellulose surfaces reported previously [50, 51]. Deposition of the PFE film on the handsheet (Figure 2.3b) results in suppression of cellulose absorption bands: 3348 cm^{-1} (O-H stretch), 1336 cm^{-1} (O-H in-plane deformation), 2902 cm^{-1} (C-H stretch), 1430 cm^{-1} (C-H deformation, asymmetric), and 1059 cm^{-1} (C-O stretch), which indicates a reduced level of oxygen and hydrogen present on the surface. The presence of new absorption bands at 1200 cm^{-1} (CF_x stretch) and 1700 cm^{-1} (unsaturated fluorocarbon bonds) confirms the presence of a fluorocarbon film. Thus, the spectrum of the handsheet after PFE deposition exhibits features of both an untreated handsheet (Figure 2.3a) and a PFE film [45]. This supports the hypothesis that plasma deposition does not significantly affect the chemical nature of the cellulose fibers, although the PFE is chemically bonded to the cellulose surface.

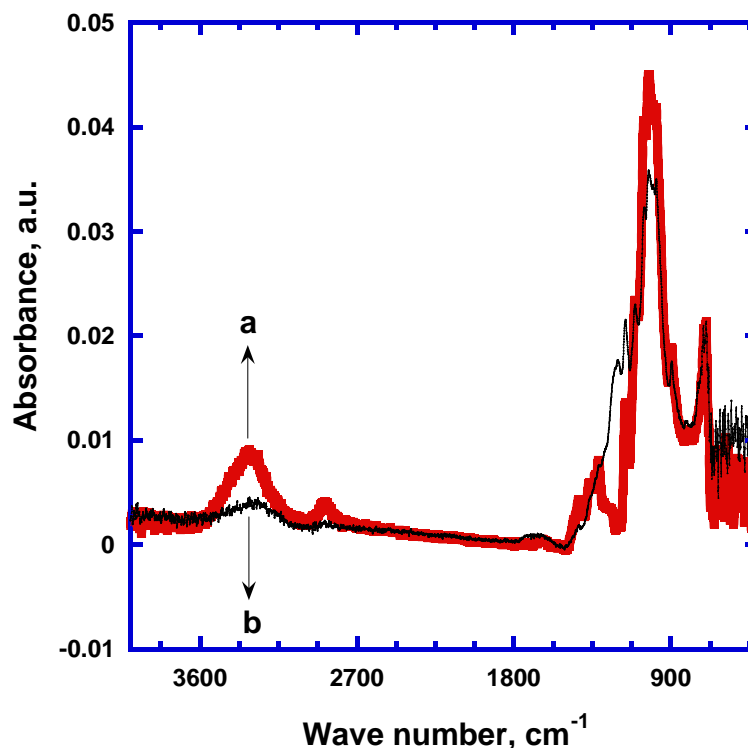


Figure 2.3 ATR-FTIR survey spectra of untreated (a) and PFE coated (~400 nm film) handsheets (b)

2.3.3 FIB

Paper is a complex porous substrate containing tortuous pores with a wide pore size distribution. Although plasma processing of paper substrates is widely referred to as a surface modification process, one might expect that there will be some penetration and reaction of plasma species into the paper bulk [43]. To study this, the untreated and PFE treated handsheets were analyzed for the penetration depth of the PFE film by Focused Ion Beam (FIB) methods. The idea was to cross-section the paper by bombarding it with gallium ions, followed by imaging the cross-sectional area via SEM.

Figure 2.4 shows cross-sectional images of untreated and PFE treated handsheets. The plasma-deposited PFE film (~ 400 nm on Si wafer under these conditions) is visible on the paper surface. However, no evidence existed for PFE deposition in subsequent fiber layers (figures not shown) within the detection limit of the SEM. Hence, it can be concluded that although there may be some pore wall fluorination deeper into the paper caused by fluorine free radicals, these do not form a film on the pore walls of detectable thickness under our experimental conditions.

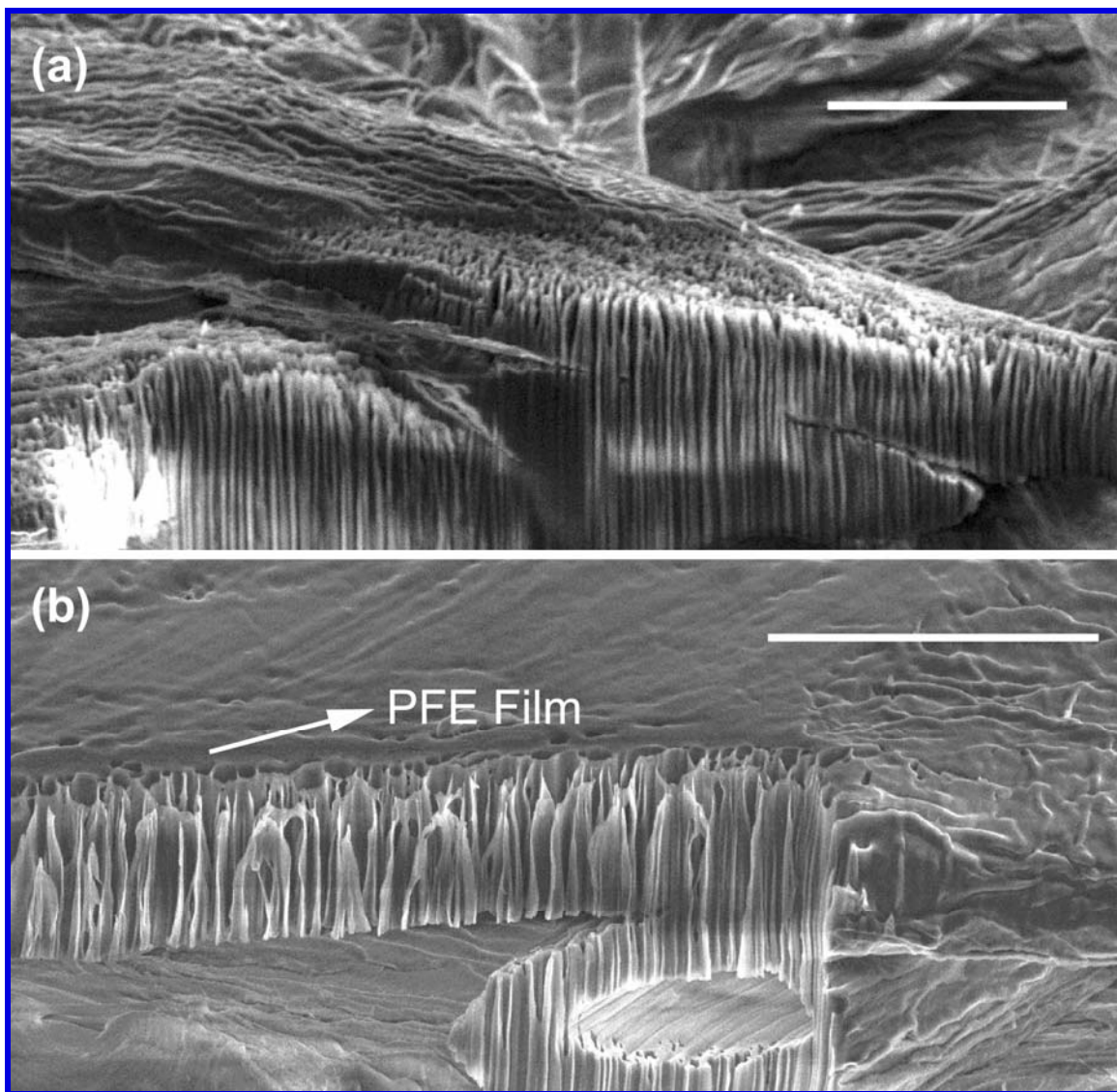


Figure 2.4 SEM images of the untreated handsheet (a) and PFE treated handsheet (b) cross-sectioned with FIB. Scale bars correspond to 10 µm

2.3.4 “Roll-off” and “sticky” superhydrophobicity

Before presenting the wettability results, it is appropriate to briefly discuss the definition of superhydrophobicity as described in previous studies. The most common definition for superhydrophobicity is the existence of a static water contact angle (CA) greater than 150° [52, 53]. Although the contact angle is a good descriptor of the interaction between

water and solid surfaces, the threshold of 150° is not sufficient to guarantee the water repellent behavior associated with lotus leaves, *i.e.* droplet roll-off and self-cleaning properties. In order to predict the mobility of water droplets, it is also necessary to determine the contact angle *hysteresis*, *i.e.* the difference between advancing and receding contact angles at the leading and trailing edge of a moving droplet. It has been found that for a CA hysteresis less than 10° , water drops roll off surfaces, while for a hysteresis greater than 10° drops tend to stick to the surface, even if such a surface has a CA greater than 150° . In previous publications, surfaces have been considered superhydrophobic purely based on the criterion of contact angles $>150^\circ$, with hysteresis values greater than 10° [5, 9, 16, 54-56], less than 10° [8, 10, 11, 15, 16, 26], or even unreported [12, 19, 40]. A debate exists on whether or not to include CA hysteresis in the definition of superhydrophobicity [15]. In this thesis, both the water CA and CA hysteresis will be included to categorize the interactions of hydrophobic substrates with liquids.

The charts in Figure 2.5 show a comparison of the water CA and CA hysteresis of copy paper, handsheet, and silicon wafer subjected to two tests: 1) SH-treatment, and 2) PFE-treatment. The copy paper used in this study was a standard copy-grade paper purchased from Office Depot ("Premium white copy paper"; brightness: 104, weight 76 g/m^2). The results for substrates without PFE deposition are not presented in the chart. CA measurements on untreated and oxygen plasma treated (only step-1 of SH-treatment) copy paper showed a CA of $81.9^\circ \pm 3.6^\circ$, while for untreated and oxygen plasma treated handsheets the water drop was absorbed into the paper in less than one second, so that the

CA could not be measured. Even with copy paper, slow water absorption takes place and as a result, CA hysteresis could not be measured for either substrate without PFE deposition. In spite of these challenges, CA measurements clearly revealed the hydrophilic nature of these samples in the absence of a fluorocarbon film.

The SH-treatment resulted in a high water CA ($>150^\circ$) and low CA hysteresis ($<10^\circ$) for both copy paper and handsheets. These substrates were superhydrophobic according to the classical definition. In contrast, the copy paper and handsheets exposed to PFE-treatment resulted in high CA ($>150^\circ$) along with large CA hysteresis ($\sim 150^\circ$). Selected images of the advancing and receding CA measurements on a PFE-treated handsheet are shown in Figure 2.6. It is evident from the image sequence in Figure 2.6 that during the receding CA measurement, the apparent solid-liquid contact area did not decrease; this underlines the extreme stickiness of the water drop to the surface.

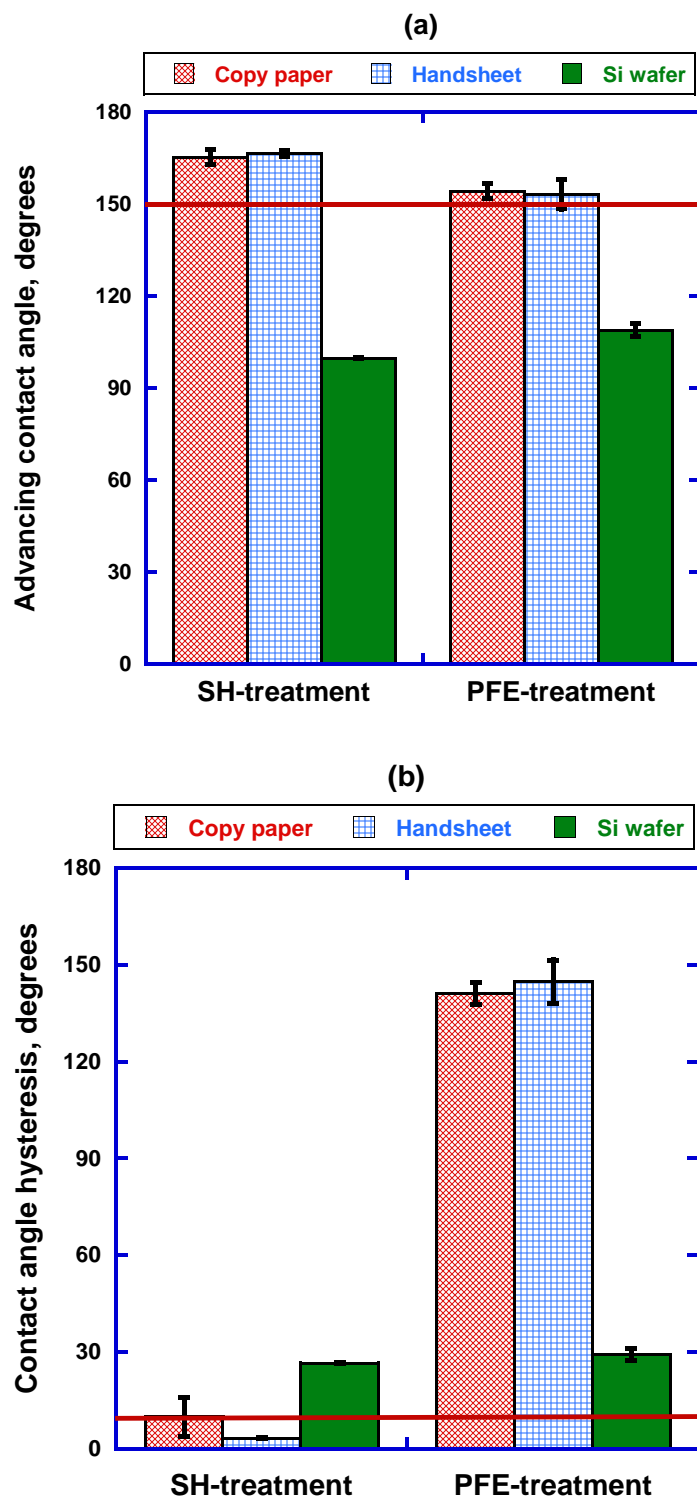


Figure 2.5 Plots of contact angle (*a*) and contact angle hysteresis (*b*) measurements for copy paper, handsheets, and Si wafers for the two plasma treatments. Red lines in (a) and (b) indicate the cut-off value for “roll-off” superhydrophobicity. Error bars represent 95% confidence intervals

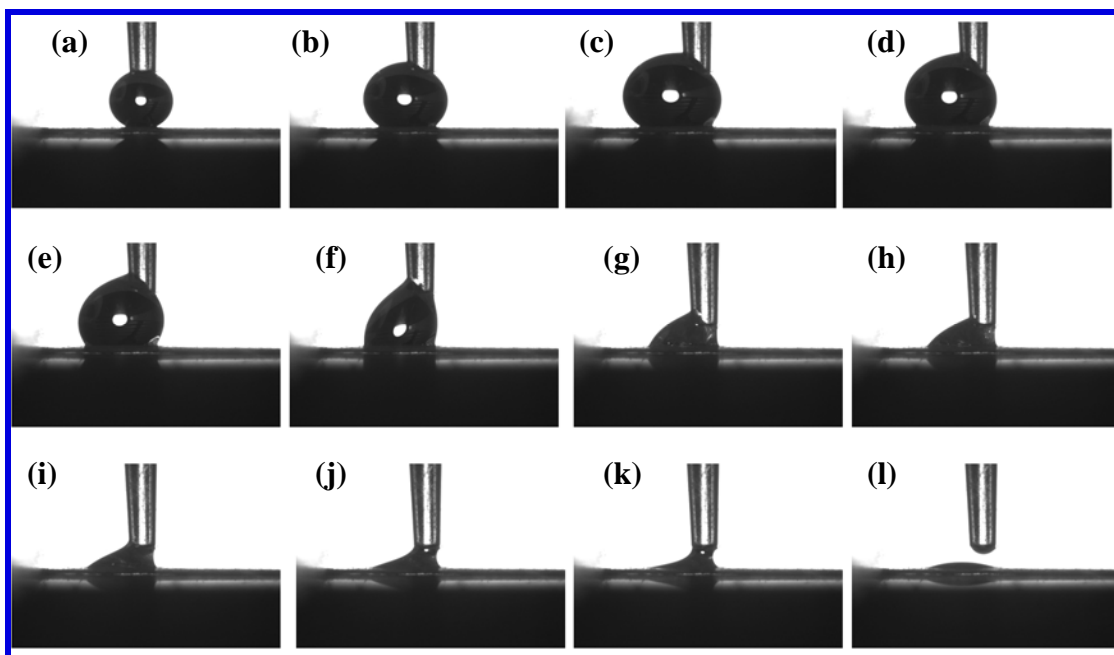






Figure 2.6 Selected images of advancing (a-c) and receding contact angle (d-l) measurements on a PFE-treated handsheet

To further illustrate the difference between the results of SH-treatments and PFE treatments, images of water drops on the treated handsheet samples are shown in Table 2.2. The static contact angle does not show a significant difference between the two substrates, since both substrates exhibit a very high contact angle ($>150^\circ$). But when the water drop is dragged across the substrate, the substrate with PFE-treatment reveals its ‘sticky’ nature as a result of a very high CA hysteresis $\sim 150^\circ$. Whereas typical “*roll-off superhydrophobicity*” is observed after SH-treatment, the properties of the PFE treated sample can best be described as “*sticky superhydrophobicity*”. In the remainder of this thesis, we will use the terms roll-off superhydrophobicity (contact angle $>150^\circ$, hysteresis $<10^\circ$) and sticky superhydrophobicity (contact angle $>150^\circ$, hysteresis $>10^\circ$) to categorize our substrates.

Table 2.2 Images of water drop on handsheets for SH-treatment and PFE-treatment

Treatment	Static	When dragged	Result
SH-treatment			“Roll-off” Superhydrophobic
PFE-treatment			“Sticky” Superhydrophobic

2.3.5 Oxygen plasma etching of amorphous cellulose domains.

The data in Figure 2.5 and Table 2.2 strongly suggest that the SH-treatment generates the desired roughness topography for “roll-off” superhydrophobicity due to etching in the oxygen plasma. However, macroscopic contact angle measurements do not provide undisputable proof for this hypothesis. To verify the effect of plasma etching (step-1) and deposition (step-2) in more detail, high-resolution SEM images were obtained for three samples: 1) untreated handsheets, 2) oxygen etched handsheets (step-1 of the SH-treatment) and 3) oxygen etched and PFE deposited handsheets (SH-treatment). Figure 2.7 shows the direct comparison at two SEM magnifications: approximately 5000X and 20,000X, respectively. It should be noted that the images display single fibers. The fluffy, “cotton-like” surface of the untreated sample can be attributed to the soft amorphous primary layer of the fibers. After oxygen etching, the fibers display a roughened surface with nanometer-scale features (Figure 2.7c, 2.7d) that are not observed on the untreated sample. The features on the roughened surface can be attributed to the

crystalline portions of the fiber, which protrude after selective etching of the amorphous portions of the fibers by the oxygen plasma treatment.

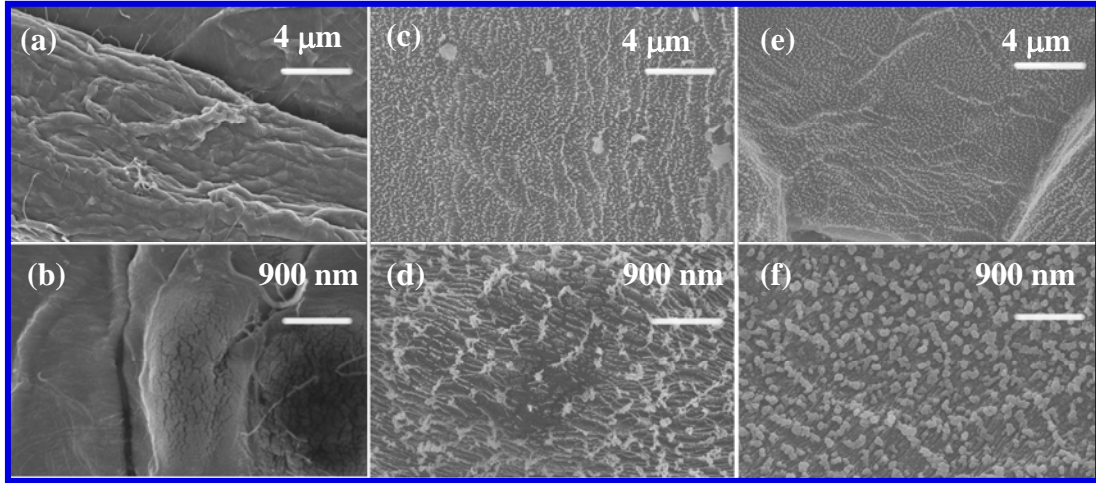


Figure 2.7 High-magnification ($\sim 5000\times$ and $\sim 20,000\times$) SEM images of: (a,b) “untreated” handsheet fiber, (c,d) oxygen-etched handsheet fiber (step-1 of SH-treatment), (e,f) oxygen-etched and PFE coated handsheet fiber (SH-treatment)

The feature sizes present on the etched sample shown in Figure 2.7c and d are consistent with the dimensions of crystalline microfibrils [31-34], which supports the basic hypothesis presented in this chapter (section 2.1). PFE deposition (~ 100 nm) on the etched surface (Figure 2.7e, 2.7f) accentuates the features and partially covers ridges created by oxygen etching. It can be concluded that the surface roughness created by the SH-treatment is sufficient to generate “roll-off” superhydrophobicity.

2.3.6 Significance of the natural topography of the cellulose fibers for the two extreme behaviors

The fundamental difference between the SH treatment and the PFE treatment is that in the latter substrates are not subjected to oxygen etching, which creates nanometer scale roughness in addition to the natural micrometer scale roughness of paper surfaces. Since paper is a porous substrate, both samples are considered to be physically heterogeneous (air pockets at the interface) irrespective of their different roughness scales. Thus, the superhydrophobic behavior observed in the “roll-off” and “sticky” superhydrophobic paper surfaces is expected to be modeled by Cassie’s model [57], which assumes that a liquid does not completely wet rough hydrophobic surfaces and which attributes the increased CA to the presence of air pockets (composite surface) at the liquid-solid interface. However, recent studies [58] have reported different contact angles for the same wetted surface fraction merely by changing surface topography. The variation in contact angles was attributed to differences in the contact line topology and tension. Thus, both the three-phase contact line topology and the wetted surface fraction are involved in establishing the contact angle; such considerations were not addressed by Cassie[58, 59]. It can therefore be concluded that the higher contact angles obtained for the “roll-off” and “sticky” superhydrophobic samples are likely due to the decreased wetted surface fraction and the difference in surface topography which changes the contact line topology.

It has been reported that the contact angle hysteresis of a superhydrophobic surface primarily depends upon two properties [8, 27], 1) metastable state energy, 2) barrier

energy for the drop to move from one metastable state to another metastable state. These two energies depend on the chemical heterogeneity, contact line topology, roughness and the wetted fraction of the surface [6, 27]. The “roll-off” superhydrophobic paper possesses a high metastable state energy and a low barrier energy which may be due to the increased contact line tension and roughness, and decreased wetted surface fraction. This energy combination causes water drops to “hop” or “skid” on the surface in search of a lower energy state, thereby causing roll-off with a CA hysteresis of 3.4° . In comparison, the “sticky” superhydrophobic paper possesses a relatively low metastable state and a very high barrier energy combination. This may be due to the decreased contact line tension and roughness, increased wetted surface fraction and chemical heterogeneity. As a result, the water drop remains pinned at the initial lower energy location without the ability to move, thereby displaying a high CA hysteresis of $\sim 150^\circ$.

Silicon wafers are flat and should not be affected significantly by treatment with oxygen plasma. Indeed, the CA and hysteresis were found to be the same for both treatment methods, with higher CA and lower hysteresis than for an untreated sample as a result of the fluorocarbon coating. It can therefore be concluded that the natural micrometer topography of the paper surface is responsible for the “sticky” superhydrophobicity. In order to obtain “roll-off” superhydrophobicity, a secondary nanoscale roughness must be added, which was achieved by uncovering the implicit nanostructure of microfibrils via oxygen etching. Both “roll-off” and “sticky” substrates are of considerable interest for applications in which it is important to manipulate the mobility of water droplets; a key

question that will be addressed in the following chapter (Chapter 3) is whether it is possible to tune contact angle hysteresis between these extremes.

2.3.7 Robustness and stability of the superhydrophobic paper substrates

As discussed previously, a critical property of superhydrophobic paper substrates for practical applications is the robustness of the small micron and submicron scale features. Even if a substrate is superhydrophobic immediately after creation, several operational factors can affect stability by decreasing the water CA and/or increasing CA hysteresis: condensation of water vapor in the air pockets at the liquid-solid interface, external pressure applied to the liquid, which compresses the air pockets, and damage to the fragile nanometer-scale features. The first two issues are related to the application in which the substrates will be used. With regards to the issue of mechanical robustness, it is expected that the robustness of the roughness generated by our process, which originates from the internal morphology of cellulosic fibers, should be better than for structures created by traditional polymer grafting or nanoparticle deposition. To confirm this hypothesis, the robustness of the surface was tested with a standardized scotch tape test (ASTM) [60, 61]. The result was damage to the paper substrate, with a layer of fibers adhering to the scotch tape. Apparently, adhesion failure occurred at fiber-fiber interfaces rather than fiber-PFE interfaces, which demonstrates that the PFE film has excellent adhesive bonding to the fibers. This observation is consistent with the fact that cellulose has numerous -OH moieties which can serve as reaction sites for covalent bonding to a crosslinked PFE film. This covalent bonding is stronger than the fiber-fiber hydrogen bonding. After failure of the scotch tape test to confirm the robustness of the topology of

the modified surface, we performed another simple wear test on the roll-off superhydrophobic handsheet by pressing it firmly with a bare finger. Although this is not a standardized test, the procedure closely replicates common handling of paper and paperboard and therefore offers insight into practical use of the modified paper surfaces. After this test, the handsheet showed an average CA and CA hysteresis of $157.1^\circ \pm 4.2^\circ$ and $21.4^\circ \pm 14.5^\circ$, respectively, which indicates that the superhydrophobicity was retained. The slight decrease of CA, increase of CA hysteresis and increased variability of both parameters after the wear test is likely due to contamination of the surface by grease/dust from the finger, although partial destruction of the nanometer-scale structures cannot be excluded. It was also observed that the superhydrophobicity of these paper substrates was retained even after feeding the treated paper substrates through a desktop printer (Xerox Phaser[®]) showing an average CA and CA hysteresis of $165.1 \pm 2^\circ$ and $29.8 \pm 2.9^\circ$ respectively. The robustness of the superhydrophobic paper substrates makes it possible to pattern these substrates with different surface energy ink islands via desktop printing to control the mobility of the drops. Patterning of superhydrophobic papers will be discussed in greater detail in Chapters 5 and 6. In addition, CA and CA hysteresis values of SH-treated handsheets and copy paper were constant after storage for several weeks under ambient conditions ($T \sim 25^\circ \text{C}$, relative humidity $\sim 40\%$). These studies establish the stability of the PFE film under ambient conditions against surface oxidation by atmospheric oxygen.

2.4 Conclusions

A two step (etching and deposition) process was used to fabricate superhydrophobic surfaces. The first step was the domain selective etching of the amorphous portions of the cellulose in an oxygen plasma to expose the crystalline cellulose structures. The second step was to coat the etched surface with a thin film of pentafluoroethane deposited via plasma enhanced chemical vapor deposition. Two types of superhydrophobicity were obtained by varying the oxygen plasma etch time: “roll-off” (contact angle (CA): $166.7^\circ \pm 0.9^\circ$ and CA hysteresis: $3.4^\circ \pm 0.1^\circ$) and “sticky” (CA: $153.4^\circ \pm 4.7^\circ$ and CA hysteresis: $149.8 \pm 5.8^\circ$) superhydrophobicity.

The superhydrophobic paper surfaces are robust, flexible, breathable, biodegradable [21] and may also be recyclable. Such conclusions are partly derived from previous studies in our group that have demonstrated that plasma deposited fluorocarbon films are flexible and breathable [44]. Preliminary studies also suggest that due to the hydrophobic nature of the fluorocarbon layer, the coated fibers can easily be separated through froth floatation during paper recycling. Hence, the approach described to fabricate superhydrophobic paper surfaces may lead to applications of superhydrophobic substrates that were thus far uneconomical. Specifically, these results should find application in textile, packaging, printing, de-inking (paper recycling), biomedical and chemical industries.

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CHAPTER 3

TUNABILITY OF THE ADHESION OF WATER DROPS ON A SUPERHYDROPHOBIC PAPER SURFACE*

3.1 Introduction

In 1805, Young [1] proposed a relationship between the forces acting at an interface between a liquid and a solid. At that time, he did not formulate the equation that now bears his name, but expressed his ideas in words: "...for each combination of a solid and a fluid, there is an appropriate angle of contact between the surfaces of the fluid, exposed to the air, and to the solid..." With several assumptions, Gibbs later reformulated this hypothesis into the equation that has been widely accepted as Young's equation [2]:

$$\gamma_{SV} = \gamma_{SL} + \gamma_{LV} \cos \theta \quad (3.1)$$

where θ is the static contact angle made by the liquid on the solid surface and γ_{SV} , γ_{SL} , and γ_{LV} are the interfacial energies between solid-vapor, solid-liquid and liquid-vapor phases respectively. One of the important assumptions in Young's equation is that the energy of the system reaches a global minimum surrounded by infinitesimally close non-

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equilibrium states in the energy landscape [2]. However, this scenario is not realistic for real surfaces because of the presence of physical and chemical heterogeneities, which give rise to the presence of local minima surrounding the global energy minimum. As a result, most real surfaces exhibit a variety of contact angles for a particular solid-liquid-vapor combination, rather than a unique “...appropriate angle of contact...” as proposed by Young. Although Wenzel [3, 4] and Cassie [5] modified Young’s equation for physical and chemical heterogeneities respectively, the inherent assumption of the system being in the global energetic minimum still exists. Indeed, Johnson and Dettre [2, 6-8] described the limitations that may be encountered when applying Young’s hypothesis, as well as other mathematical formulations (*e.g.*, Wenzel’s and Cassie’s equations), to real surfaces.

More than one hundred years after Young’s essay, the first observation of multiple contact angle (CA) values for the same solid-liquid-vapor system was reported [9]. Despite the observation of various CA values, well-defined maximum and minimum contact angles could be identified and the term “contact angle hysteresis” was introduced to designate the difference between these two contact angles [9]. The maximum and minimum contact angles were later termed advancing and receding contact angles, respectively. In order to properly characterize the wetting properties of all substrates, it is critical to take contact angle hysteresis into account.

The presence of hysteresis becomes critical with regards to superhydrophobic surfaces (as defined by a static or advancing $CA > 150^\circ$). Superhydrophobic surfaces are often correlated with water repellency, but the presence of hysteresis on a superhydrophobic

surface might disguise such behavior as it causes water drops to stick to the surface. Water droplets easily roll off superhydrophobic surfaces if the CA hysteresis is below $\sim 10^\circ$, depending on the size of the drops. For larger CA hysteresis, water droplets tend to stick even to superhydrophobic substrates with advancing CA greater than $\sim 150^\circ$. As a result, the functional classification of superhydrophobic substrates depends on CA hysteresis. Approximately 10 years ago it was noted that for the characterization of (super)hydrophobic substrates it is necessary to determine both the advancing and receding CAs (*i.e.*, hysteresis) [10, 11]. Over the past few decades various names have been used to describe surfaces with static or advancing CA larger than 150° (Table 3.1).

Table 3.1 Terminology used to describe superhydrophobic surfaces

Static or advancing CA	CA hysteresis	Terms used to describe superhydrophobicity
$>150^\circ$	$<10^\circ$	Absolutely hydrophobic [12], Water-repellant [5], Ultrahydrophobic [10, 13] and Superhydrophobic [14-16]
$>150^\circ$	$>10^\circ$	Super water-repellent [17], Highly hydrophobic [18], Water repulsive [19, 20] and Superhydrophobic[21-24]
$>150^\circ$	NA reported	Ultra water-repellant [25-27], Ultrahydrophobic [28, 29], and Superhydrophobic [15, 30]

A recent review on superhydrophobicity indicates that only surfaces with both an advancing CA greater than 150° and CA hysteresis less than 10° should be termed superhydrophobic [31] ; however, this strict definition excludes surfaces with CA hysteresis $> 10^\circ$, which exhibit interesting behavior. In the previous chapter (Chapter 2) superhydrophobic substrates with vastly different CA hysteresis values were reported and, to avoid confusion, simple terms were proposed to describe two limiting cases of

superhydrophobicity: roll-off superhydrophobicity ($CA > 150^\circ$ and CA hysteresis $< 10^\circ$) and sticky superhydrophobicity ($CA > 150^\circ$ and CA hysteresis $> 10^\circ$).

“Roll-off” superhydrophobic surfaces have been extensively investigated in reference to self-cleaning substrates, often referred to as surfaces mimicking the lotus effect. By comparison, much less information has been reported on the fabrication of “sticky” superhydrophobic surfaces that possess a unique combination of high adhesion force and small liquid-solid contact area. The mobility of water droplets on such sticky superhydrophobic surfaces can, in principle, be manipulated by tuning the adhesion force, which may enable novel applications in the chemical and biomedical arenas. Examples are the development of “tweezers” for liquid drops, controlled mobility of liquid drops for lab-on- chip (LOC) applications, selective permeability in membranes, as well as proteomics and genomics microarrays.

To our knowledge, the first sticky superhydrophobic surface was fabricated by Johnson and Dettre [7], and displayed an advancing $CA \sim 160^\circ$ with CA hysteresis $\sim 100^\circ$; this discovery has gone largely unnoticed, because of the lack of knowledge about possible applications at that time. More recently, a number of publications have appeared that report superhydrophobic surfaces with high hysteresis values [12, 13, 17, 21]. In addition, a number of publications have reported superhydrophobic surfaces that are able to immobilize water droplets up to various tilt angles [18, 21, 32-36], without actually reporting CA hysteresis values. Finally, several publications have reported the stickiness of a superhydrophobic surface in terms of adhesion forces [22-24, 37]. Because the

adhesion force of a droplet to a surface depends on factors like surface inclination angle and drop size [33, 38-41], CA hysteresis is a more appropriate parameter to enable comparison of surface characteristics.

So far, there have been no reports regarding the fabrication of superhydrophobic surfaces that can be tuned from sticky to roll-off behavior by simply controlling the receding contact angles from superhydrophilic ($<10^\circ$) to superhydrophobic ($>150^\circ$). The publications on sticky superhydrophobic surfaces referenced in the preceding paragraph suggest that droplets can be removed completely from a sticky substrate if the applied force is sufficiently large. However, if the receding contact angle is superhydrophilic ($<10^\circ$), the droplet forms a bridge and breaks during its separation from the sticky superhydrophobic surface. Reports to date have indicated only a small window for hysteresis tunability, with the maximum hysteresis less than 100° [7, 12, 13, 21, 32].

This chapter reports a methodology for tuning CA hysteresis (*i.e.*, stickiness) of superhydrophobic paper surfaces between $149.8 \pm 5.8^\circ$ and $3.5 \pm 1.1^\circ$. This was achieved by controlling the formation of physical heterogeneities (roughness) of the paper fibers during selective etching of the amorphous cellulose domains in an oxygen plasma. Low surface energy is obtained by depositing a thin film of fluorocarbon onto the etched paper surfaces using pentafluoroethane (PFE) monomer from a plasma environment (as described in Section 2.2.2).

3.2 Experimental

3.2.1 Handsheet formation

Four types of substrates were prepared with different fiber combinations and drying methods as shown in Table 3.2. A more detailed description of this procedure can be found in Section 2.2.1 (Chapter 2).

Table 3.2 Method of drying and fiber type of handsheets

Handsheet	Method of drying	Fiber type
HS-OD	Overnight drying	50% hardwood - 50% softwood
H	Rapid drying	100% hardwood
S	Rapid drying	100% softwood
HS	Rapid drying	50% hardwood - 50% softwood

3.2.2 Plasma etching/deposition

The details of the parallel plate rf (13.56 MHz) plasma reactor configuration and operational procedures can be found in section 2.2.2 (Chapter 2). The experimental conditions for oxygen etching to create appropriate roughness of the cellulose fibers and subsequent deposition of a fluorocarbon film from pentafluoroethane (PFE) monomer are listed in Table 3.3.

Table 3.3 Plasma reactor parameters for the etching of and deposition on handsheets

Parameters	Etching	Deposition
Gas	Oxygen	Pentafluoroethane (PFE) and Argon
Flow rate	75 sccm	20 sccm (PFE) and 75 sccm (Argon)
Temperature	110° C	110° C
Pressure	0.55 Torr	1 Torr
Power	10 W	120 W

3.2.3 SEM

SEM analysis was performed on the treated paper substrates after depositing a thin layer of gold (~15 nm). A more detailed description of this method and apparatuses used can be found in section 2.2.3 (Chapter 2).

3.2.4 Water contact angle measurements

A more detailed description of this method can be found in section 2.2.7 (Chapter 2).

3.3 Results and discussion

3.3.1 Tunability of the stickiness

In the previous chapter (Chapter 2), the fabrication of “sticky” and “roll-off” superhydrophobic paper surfaces with contact angle hysteresis of $\sim 4^\circ$ and $\sim 150^\circ$, respectively, was reported. After establishing these extreme values of hysteresis for the sticky and roll-off superhydrophobic paper substrates, this work focuses on systematically varying the hysteresis. It was expected that the stickiness (hysteresis) could be manipulated by controlling the physical heterogeneity of the paper surfaces. This was achieved by selectively etching the amorphous portions of the cellulose fibers by selective oxygen etching (explained in Chapter 2). To investigate this process, handsheets (HS-OD) were etched in an oxygen plasma for different durations (0 to 60 minutes). Fluorocarbon film deposition was then performed onto the handsheets (HS-OD) for 2 and 15 minutes to form thin PFE films.

Figure 3.1 shows the advancing CAs, receding CAs and CA hysteresis for handsheets (HS-OD) processed for different etching times. It is evident from Figure 3.1 that while

the advancing CA remains in the superhydrophobic regime for all samples ($>150^\circ$), the receding CA gradually increases with etching time from a superhydrophilic value of $8.4 \pm 6.8^\circ$ to superhydrophobic value of $155.4 \pm 1.6^\circ$. This smooth transition of the receding CA results in tunable CA hysteresis values from $147.2 \pm 6.8^\circ$ (sticky) to $7.6 \pm 1.6^\circ$ (roll-off). Controlling the oxygen etching times thus enables the generation of paper surfaces with various degrees of stickiness.

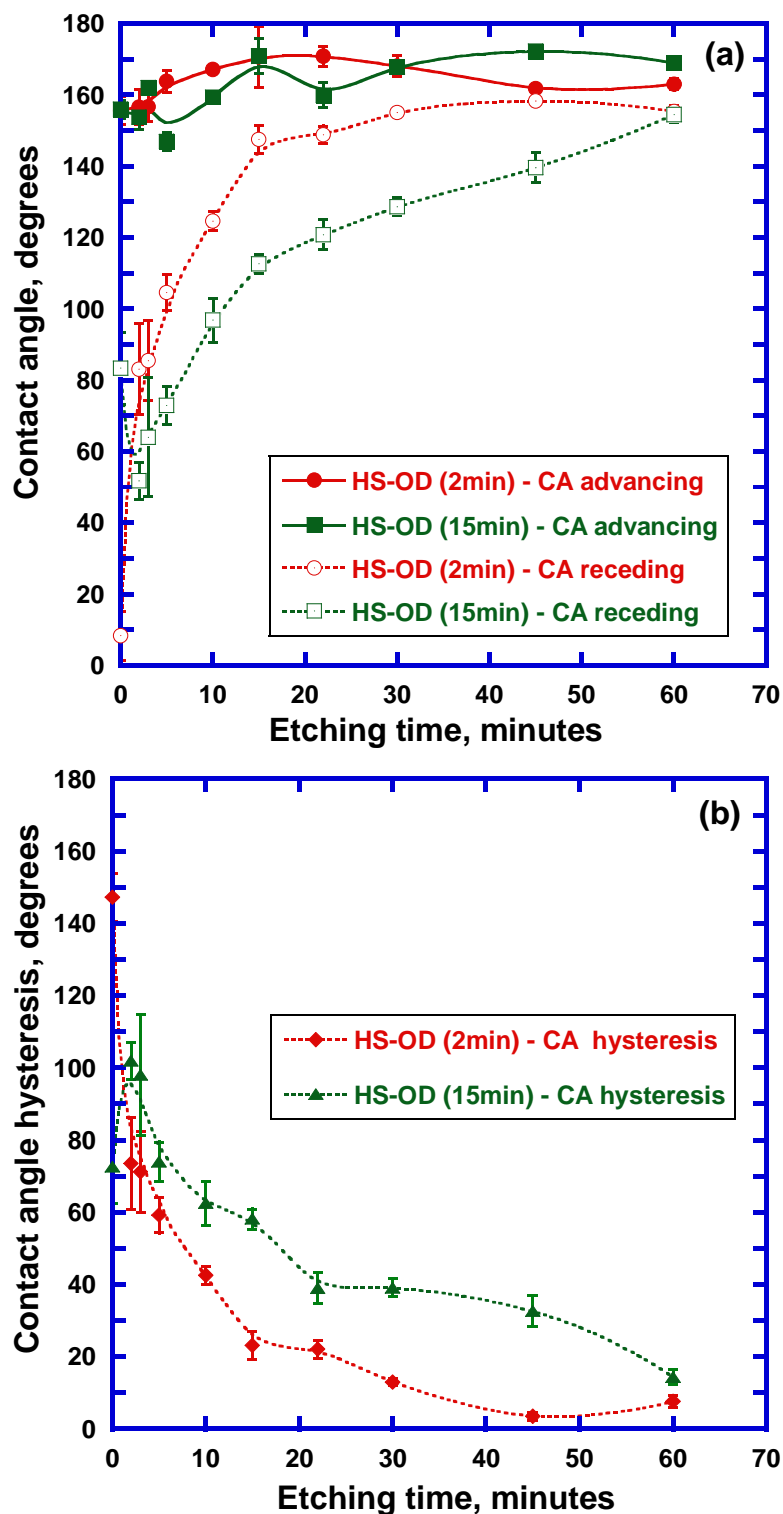


Figure 3.1 a) Plot of advancing CA and receding CA of handsheets (HS-OD) with respect to oxygen plasma etching time for 2 min and 15 min PFE depositions, b) Plot of CA hysteresis of handsheets (HS-OD) with respect to oxygen plasma etching time for 2 min and 15 min PFE depositions

SEM images of the handsheets (HS-OD) for selected etch times and 2 minute PFE deposition are shown in Figure 3.2a (high magnification) and 3.2b (low magnification). On the nano-scale (Figure 3.2a), we see that as oxygen etching proceeds, the fiber surfaces are roughened due to the formation of small nano-scale features. After 10 minutes etching, followed by fluorocarbon deposition, the formation of mushroom-like features is apparent. The feature spacing increases with prolonged etching (30 minutes) and when 60 minutes etching is followed by deposition, the mushroom-like features disappear; at this point the nano-scale roughness consists of solid ridges. On the micro-scale (Figure 3.2b), the number density of fibers on the surface decreases with etching time, thereby enhancing the micro-scale roughness. The oxygen plasma apparently etches away soft fiber dust and primary layers of fibers. It is evident from the figures that oxygen etching roughens the handsheet surface on two separate length scales; the extent of both roughness scales qualitatively increases with etching time, at least up to 30 minutes etching.

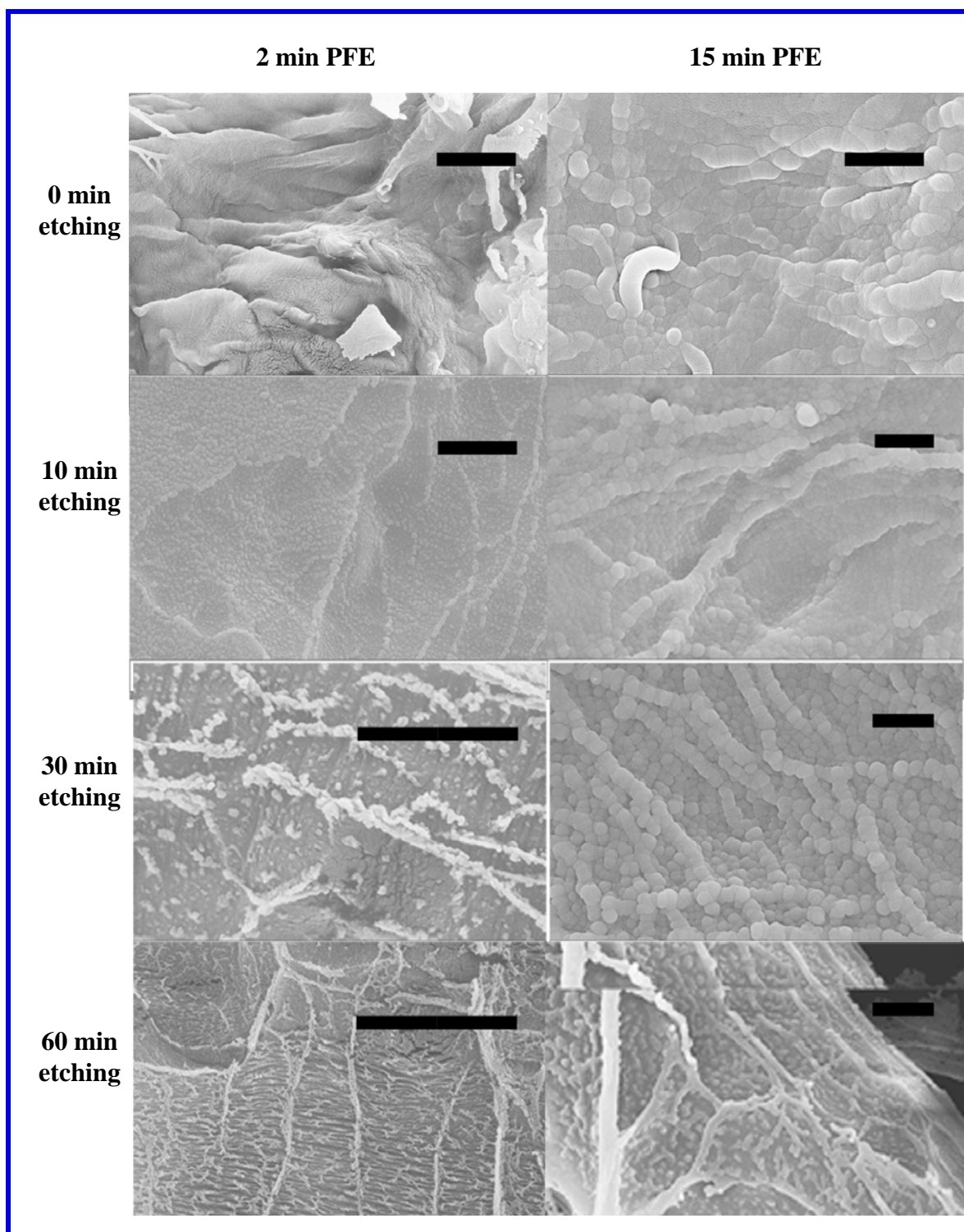


Figure 3.2a High magnification images of the HS-OD (with 2 min PFE deposition) and HS-OD (with 15 min PFE deposition) for 0, 10, 30 and 60 minutes etching times. The scale bars correspond to 2 μm

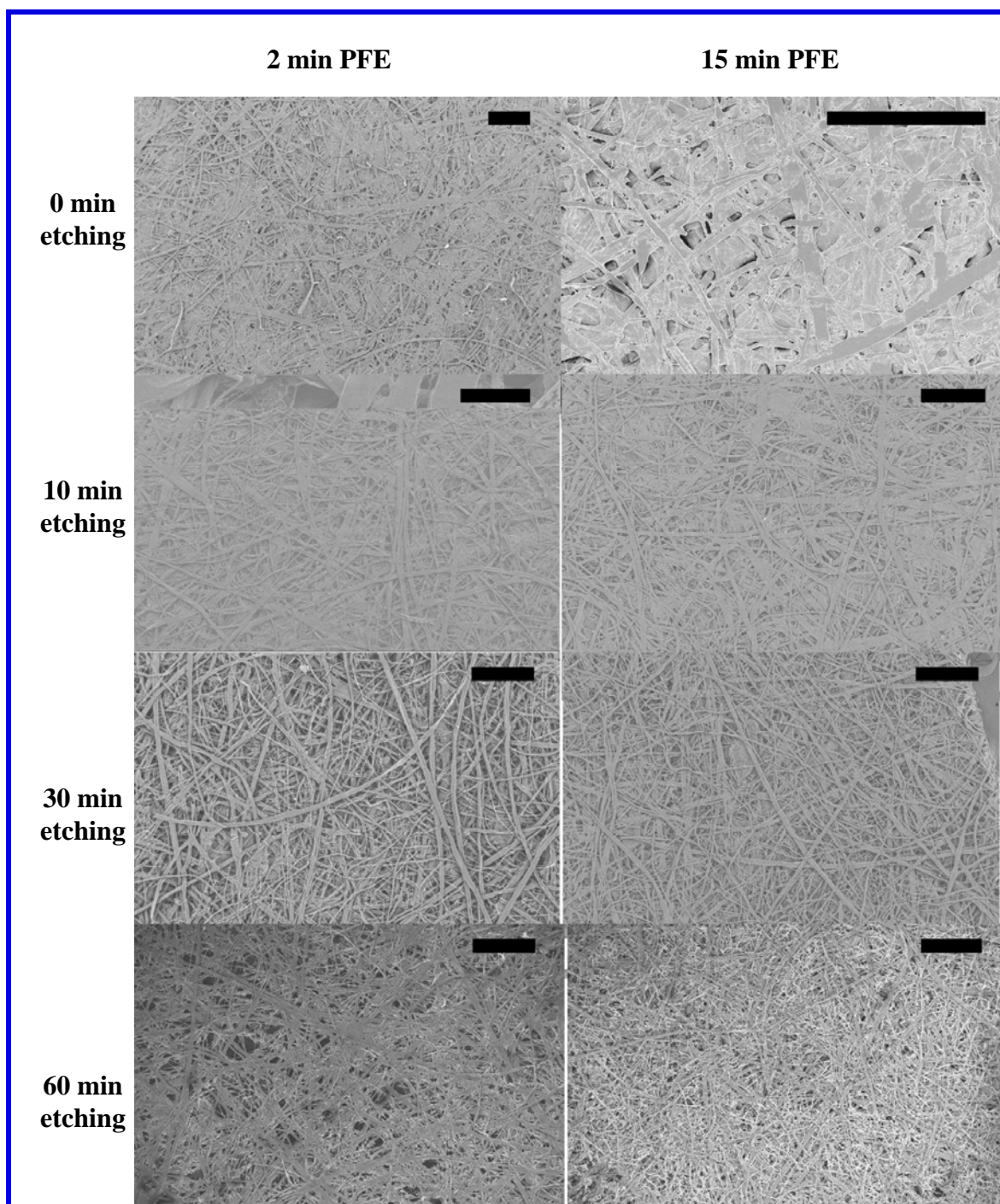


Figure 3.2b Low magnification images of the HS-OD (with 2 min PFE deposition) and HS-OD (with 15 min PFE deposition) for 0, 10, 30 and 60 minutes etching times. The scale bars correspond to 400 μm

3.3.2 Mechanism for tunability of the stickiness

The interactions of a liquid drop with a physically heterogeneous (rough) solid surface can be explained by two classic equations: the Wenzel [3] and Cassie [5] equations.

When a liquid fully penetrates into the rough grooves of the surface, the apparent contact angle made by the liquid drop with the surface is described by Wenzel [3] (Equation 1.1).

On the other hand, when the liquid does not completely penetrate into the solid grooves, leaving air voids at the apparent solid-liquid interface, the apparent contact angle of the liquid drop with the surface was described by Cassie and Baxter [5] (Equation 1.2).

Johnson and Dettre [2, 7] later combined these two equations (1.1 and 1.2) to model the transition from a Wenzel type wetting to a Cassie type wetting as a function of roughness.

They concluded that for a hydrophobic surface, the CA hysteresis increases with roughness if the wetting is governed by the Wenzel state. On the other hand, the CA hysteresis decreases with roughness if the wetting is governed by the Cassie state.

Paper is a porous substrate and consists of many tortuous pores formed by the complex network of cellulose fibers. When a thin PFE film is deposited (~100 nm), only the individual fibers are coated and the porosity and roughness of the paper substrate remain essentially unchanged [42]. Therefore, even without the presence of roughness created by oxygen etching, an inherent micro-scale roughness of the paper surface results from the network of fibers and the porosity of the paper; this causes the interaction of a liquid drop with paper to be in the Cassie regime. From Figure 3.1 it is evident that the hysteresis of the handsheets (HS-OD) decreases with respect to the etching time, which suggests that the etched handsheets are in the Cassie regime.

The static CA for a smooth PFE film deposited on a silicon wafer is approximately 105° . After a continuous PFE film covers the fiber surface, the surface chemistry of the fiber becomes identical to that of a PFE film (data not shown). The PFE deposition time (2 minutes) for these experiments was chosen in order to obtain a continuous coating on the fibers with a thickness (~ 100 nm) that is sufficient to retard water absorption into paper [42]. As a result, the fiber surface after a 2 minute PFE deposition should have a Young's CA approximately equal to 105° .

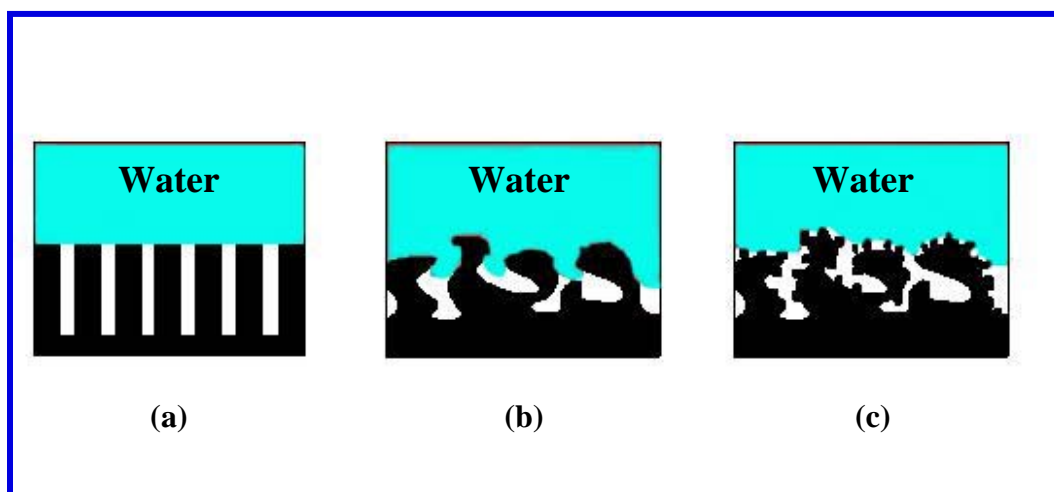


Figure 3.3 Schematics of interactions of water with surfaces at an ideal Cassie state (a), sticky superhydrophobic state (b), and roll-off superhydrophobic state (c)

If the water-substrate interaction occurs in the ideal Cassie regime (Figure 3.3a) with chemical heterogeneity provided by the PFE film and air, the minimum possible receding CA obtained would be $\sim 105^\circ$ [6, 8]. However, in Figure 3.1 we see that for 0 minutes etching, the receding CA values are well below 105° , as low as $\sim 10^\circ$. Such low receding CA values for these samples are only possible if pinning of the three phase contact line occurs by physical heterogeneities on the handsheet fibers. Since the wetting of handsheets seems to be in the Cassie regime (irrespective of the etching time), the high

hysteresis could result only if there is a possibility of a nano-scale Wenzel state [3, 4] on the top of each fiber of the handsheet, characterized by an enhanced liquid-substrate contact area (figure 3.3b). This type of interaction would increase the energy barrier between metastable states. For a surface with a very low solid-liquid contact area, a higher energy barrier greatly decreases the receding CA while only slightly increasing the advancing CA [8], as we observe for the handsheets. These types of interactions occur in the unetched handsheet with 2 min PFE deposition, resulting in a high, superhydrophobic advancing CA and a low, superhydrophilic receding CA. This type of substrate can be categorized as Cassie state on the micro-scale and Wenzel state on the nano-scale. The unique combination of Cassie and Wenzel states provides a large energy barrier for the movement of the three phase contact line, resulting in a very low receding CA, in spite of a superhydrophobic advancing CA. Similar types of interactions for different roughness scales are also referred to as “air pocket trapping” [43], “Cassie impregnation” [37, 44], “petal effect”[37], “transitional superhydrophobic state between Wenzel’s and Cassie’s states”[45] and “Cassie–Wenzel wetting transition” [44, 46-50].

After etching, the formation of the nano-scale features on the fiber surface (see *e.g.* 30 minutes in Figure 3.2) provides a different liquid-solid interface. The additional roughness created by the nano-scale features results in Cassie state wetting at the nano-scale (Figure 3.3c). These nano-scale features greatly reduce the wetting area, as well as the energy barrier between metastable states. As a result, the vibration energy of the drop is now high enough to overcome the energy barriers, giving rise to roll-off behavior (low CA hysteresis). This type of interaction can be classified as Cassie regime on the nano-

and micro- scale. This type of interaction is commonly referred to as the “lotus state” [45], in reference to its natural occurrence on the leaves of lotus plants.

The schematics in Figure 3.3b and Figure 3.3c can be directly compared with the SEM images (Figure 3.2a and 3.2b) for 0 minutes and 60 minutes etched handsheets (HS-OD). Figure 3.1 shows that by varying the etching time, the hysteresis can be tuned from a sticky superhydrophobic to a roll-off superhydrophobic value. It is believed that this tunability is obtained by a smooth transition from the Wenzel to Cassie state at the nano-scale (fiber surface) due to the evolution of the nano-scale features. This conclusion is supported by the SEM images in Figure 3.2a and 3.2b (2 min PFE). The transition of a drop’s interaction from Cassie state to Wenzel state has been previously obtained by inducing pressure [51] and vibration[44, 46, 48-50]. Here, we obtain this transition without using an external stimulus by controlling the topography of the fibers or, in other words, by evolution of the nano-scale features. In conclusion, the unique combination of transition in wetting that occurs at the nano-scale, while maintaining a Cassie state at the micro-scale results in variable receding CA while maintaining a superhydrophobic advancing CA.

3.3.3 Significance of nano -scale roughness on the tunability of stickiness

Fiber diameters in the handsheet range from 10 to 40 μm . Thus, a PFE film thickness of less than $\sim 1 \mu\text{m}$ would not be expected to affect the micro-scale roughness. Two PFE deposition times were used to coat the etched handsheets (2 and 15 minutes) to study the effect of nano-scale roughness on tunability. As explained previously, the 2 minute PFE

deposition (~100 nm) was chosen because it just exceeds the minimum thickness to retard water absorption [42]. The 15 minute PFE deposition (~1 μm) was chosen because we believe this relatively thick film would not disturb the micro-scale features.

Comparisons of the advancing CA, receding CA and CA hysteresis for the 2 min and 15 min PFE deposited handsheets (HS-OD) are shown in Figure 3.1a and 3.1b. The 15 min PFE deposition on an unetched paper results in a CA hysteresis of $72.6 \pm 10^\circ$ which is lower than the value obtained for a 2 min PFE deposition. This observation is consistent with the SEM images in Figure 3.2a. With increased PFE deposition time (15 minutes), the presence of globular features is visible on the unetched handsheets (0 min). It is believed that these nano-scale features decrease the hysteresis by enhancing roughness. On the other hand, the CA hysteresis goes through a maximum value after 2 minutes etching and then decreases, following the same trend as the handsheets with 2 min PFE deposition. As explained previously, hysteresis increases with roughness if the interaction between the water drop and surface is in the Wenzel state. Hence, it can be concluded that 2 minutes etching increases the roughness scales (from 0 minutes etching) appropriate for Wenzel state wetting. After 3 minutes etching, the fiber topography is appropriate for wetting in the Cassie regime, resulting in a decrease in the CA hysteresis with respect to etching (roughness). It is evident from Figures 3.1a and 3.1b that the hysteresis values for the etched fibers with 15 min of PFE deposition are consistently higher than for 2 min PFE deposition. This increased hysteresis may be attributed to the larger nano-scale features observed on handsheets after 15 min PFE deposition (Figure 3.2a) which provide a larger solid-liquid contact area.

The morphological changes at the nano-scale due to longer deposition times (thicker PFE film) are a unique feature of the plasma deposition process. The plasma deposition process takes place at reduced pressure (~ 1 Torr). At this pressure, the reactant species flux onto the surface, which establishes the deposition rate, is determined by surface topography. That is, the film thickness at a particular surface location depends upon the acceptance angle for species impingement and thus on the local geometry [52]. Ideally, at the top of a 90° step edge, reactant flux impinges from a 270° cone and hence the edge receives a higher flux than the flat surface (180° cone) or the corner at the bottom of the step edge (90° cone). Because deposition rates scale with these fluxes, This phenomenon plays a crucial role in the modification of the oxygen etched fiber morphology due to the variation of PFE deposition thickness and deposition time.

Figure 3.4, shows a plot of CA hysteresis and advancing CA for the different handsheets at various etching times (0 to 60 minutes) with either 2 min PFE deposition (HS-OD, H, S and HS) or with 15 min PFE deposition (HS-OD only). This figure reiterates that almost all substrates (except two) are superhydrophobic by the classical definition ($CA > 150^\circ$), but that they span a wide range of stickiness (CA hysteresis).

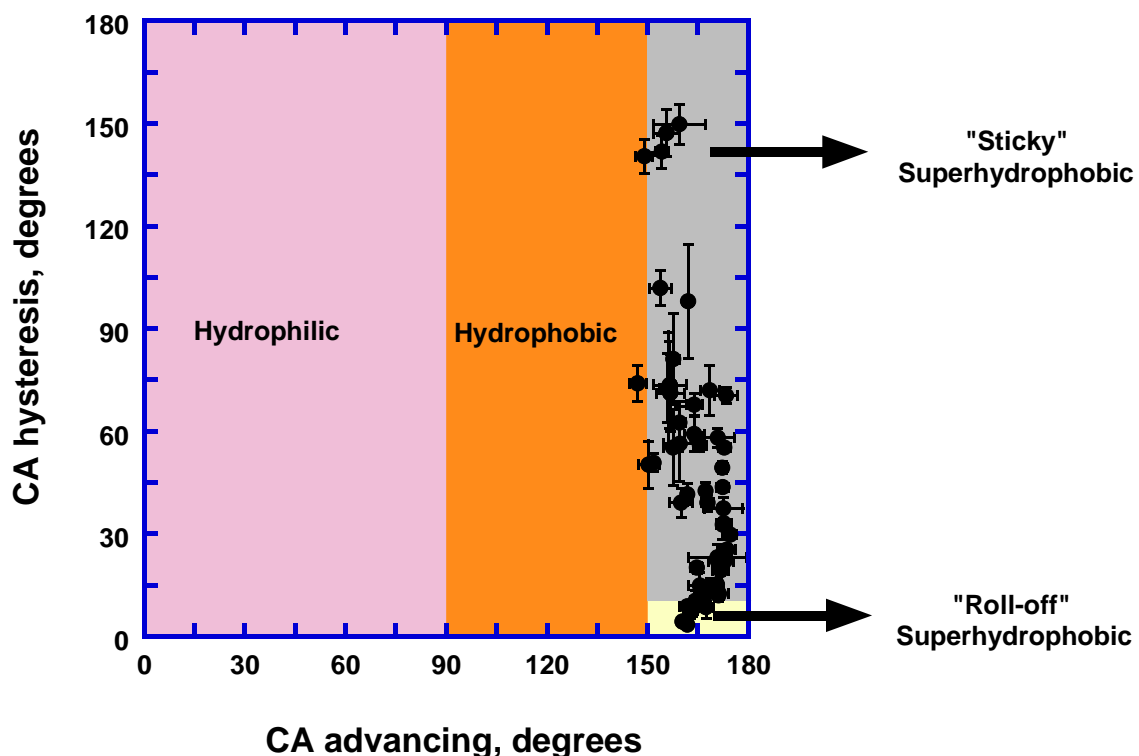


Figure 3.4 CA hysteresis versus CA for all samples investigated for various etching times: HS-OD, H, S and HS with 2 min PFE deposition and also HS-OD with 15 min PFE deposition

3.4 Conclusions

The fabrication of superhydrophobic paper surfaces with a wide range of contact angle hysteresis values has been achieved, which results in variable stickiness of water drops on these substrates. The tunability of hysteresis was obtained by creating physical heterogeneity (roughness) on the cellulose fibers via selective oxygen plasma etching. Tunable stickiness of water drops on paper-based substrates can be important in controlling the mobility of drops on the surface. Sticky superhydrophobic surfaces can also be used to make “tweezers” for liquid drops that allow the transfer of a large

volume of liquid with low solid-liquid contact area. These concepts may open new windows of opportunity in chemical and biomedical research areas such as microfluidics, lab-on-Paper (LOP), MEMS, proteomics and genomics. In addition, generation of tunable surface properties on biodegradable, inexpensive, recyclable paper substrates may widen the scope of these possible applications.

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CHAPTER 4

INSIGHTS INTO THE DESIGN OF SUPERHYDROPHOBIC PAPER SURFACES*

4.1 Introduction

Paper/cellulose is an inexpensive commercial polymer which, depending upon the specific application, is manufactured in different forms as determined by the choice of fiber source and paper making method. Hence, it is reasonable to expect that different types of paper will display significant differences in micro- and nano-scale roughness, which are key factors in determining their superhydrophobic behavior. In this chapter, the impact of variations in fiber type and paper making process on the superhydrophobic properties of paper surfaces is investigated. These experiments provide insight into the appropriate plasma processing conditions for the design and fabrication of superhydrophobic paper surfaces for particular applications. In addition, the effect of water drop size on contact angle and contact line geometry as observed in superhydrophobicity measurements on heterogeneous porous substrates such as paper is discussed.

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4.2 Experimental

4.2.1 Paper substrates

Five types of paper substrates were used for superhydrophobicity studies as described in Table 4.1. Handsheets (H, S, HS) were fabricated following the TAPPI method T205 sp-02 with southern hardwood kraft (Alabama River Pulp co.) and/or southern softwood kraft (North Carolina International Paper). A more detailed description of the preparation of these handsheets is given in Sections 2.2.1 (Chapter 2) and 3.2.1 (Chapter 3).

Commercial copy paper substrates, "Premium white copy paper", were obtained from Office Depot. Commercial paper towels, SCOTT[®] High Capacity Hard Roll Towels (product code 01000) manufactured by Kimberly-Clark Professional were also used.

Table 4.1 Paper substrates used for superhydrophobicity studies

Substrates	Description
H	Handsheet (100% hardwood fibers)
S	Handsheet (100% softwood fibers)
HS	Handsheet (50% hardwood - 50% softwood)
CP	Copy paper
PT	Paper towel

4.2.2 Plasma etching/deposition

A 6 inch parallel plate rf (13.56 MHz) plasma reactor was used for plasma etching and deposition sequences; substrates were heated to 110° C. Details of the reactor configuration and operational procedures for the treatment of paper substrates can be found in section 2.2.2 (Chapter 2). Experimental conditions for oxygen etching and fluorocarbon (pentafluoroethane (PFE)) film deposition are listed in Table 3.3 (Chapter 3).

4.2.3 SEM

SEM micrographs were obtained using a LEO scanning electron microscope after sputtercoating (EMS 350) the paper substrates with a thin gold film (~15 nm). A more detailed description of the method and apparatus can be found in section 2.2.3 (Chapter 2).

4.2.4 Water contact angle measurements

For standard contact angle measurements, a 4 μL water drop was used. Advancing and receding contact angles were measured by moving the substrate left to right with respect to the drop, following the methodology discussed by Gaudin *et al.* [1]. For substrates with receding CAs less than 10° , the standard “volume decrement” method was used. A more detailed description of the methods used to measure contact angles can be found in section 2.2.7 (Chapter 2).

4.2.5 Microscopic imaging of the contact line

Water drops with appropriate volumes were dispensed onto sticky superhydrophobic paper (HS) surfaces that had been attached to microscope slides. The existence of a superhydrophobic contact angle made it impossible to obtain a clear image of the three-phase contact line with the microscope due to lensing effects of the drop. The water drops were therefore allowed to evaporate until a hydrophilic contact angle was observed and the contact line was then imaged with a Leica microscope (DM4500 B; Leica Microsystems, Inc., Bannockburn, IL) using a 10X objective.

4.3 Results and discussion

4.3.1 Effects of papermaking parameters on achievement of superhydrophobicity

The superhydrophobicity imparted to paper substrates results from the combination of a low surface energy film and two-scale roughness (nano-scale and micro-scale). The nano-scale roughness originates from the exposure of crystalline domains on fiber surfaces by selective plasma etching/removal of the surrounding amorphous domains as discussed extensively in Chapters 1 and 2. On the other hand, the micro-scale roughness is determined by the topography of the paper fibers, in particular the fiber size and mesh size of the cellulose web. In this study, the effect of paper making parameters on the evolution of surface roughness length scales and the resulting superhydrophobicity of the paper substrates is investigated. The two main categories by which paper fabrication parameters can be classified are: 1) fiber source (chemical composition and physical characteristics) and 2) paper making technology used during processing, including pulping and paper machine configuration. The current work indicates how these parameters affect the conversion of a hydrophilic paper/cellulose surface into “sticky” and “roll-off” superhydrophobic surfaces using plasma treatments.

4.3.1.1 Effects of fiber type

Cellulose paper is typically produced from hardwood fibers, softwood fibers, or a combination of both. This classification of cellulose fibers is based on the trees from which they are obtained: hardwood fibers come from angiosperm trees (*e.g.*, American yew, Common juniper, Douglas fir), and softwood fibers originate from gymnosperm trees (*e.g.*, wild plum, peach, pear) [2, 3]. Both fiber types have approximately the same

chemical composition: cellulose (40-50%), hemicellulose (25-35%) and lignin (20-35%), but significantly different physical dimensions [2]. Softwood fibers are usually about twice as large as hardwood fibers as shown in Table 4.2. Considering these facts, it is reasonable to expect that: 1) different cellulose fibers types could show differences with regards to the evolution of nano-scale roughness during etching (exposure of crystalline domains) and 2) the different fiber sizes will impact the micro-scale roughness of the paper surfaces. In order to investigate the role of fiber types in more detail, handsheets were fabricated from three different combinations of hardwood and softwood: 100% hardwood (H), 100% softwood (S) and 50% hardwood-50% softwood (HS). Other than the origin of the fibers, all procedures for handsheet fabrication were constant (see Experimental section).

Table 4.2 Typical dimensions of hardwood and softwood fibers [2]

Fiber type	Fiber length, mm	Fiber width, μm
Hardwood	1.0 - 1.5	16 - 22
Softwood	3.0 - 3.7	27 - 38

Figures 4.1a-c show high and low magnification SEM images of untreated handsheets for 100% softwood, 100% hardwood and a 50%-50% hardwood/softwood mixture, respectively. The larger size of softwood fibers in comparison with hardwood fibers is confirmed by the SEM images of Figure 4.1a-c. In addition, it was confirmed by SEM images (not shown) that a thin film of PFE ($\sim 100\text{nm}$) deposited on unetched handsheets does not significantly alter the roughness (either micro- or nano-scale) of the handsheet surfaces (H, S, HS). This PFE deposition without oxygen etching yielded “sticky” superhydrophobic properties for all three handsheets with the following advancing and

receding CAs: H (CA_{adv}/CA_{rec}) - $154.3^{\circ} \pm 1.9^{\circ} / 12.5^{\circ} \pm 5.0^{\circ}$, S (CA_{adv}/CA_{rec}) - $149.0^{\circ} \pm 2.5^{\circ} / 8.5^{\circ} \pm 5.0^{\circ}$ and HS (CA_{adv}/CA_{rec}) - $159.4^{\circ} \pm 7.7^{\circ} / 9.7^{\circ} \pm 5.8^{\circ}$. The fact that the advancing and receding CAs are similar for the three handsheets shows that differences in micro- and nano-scale roughness due to variations in fiber type do not significantly affect the “sticky” superhydrophobic behavior.

In subsequent experiments, handsheets were etched in an oxygen plasma for different durations before depositing the PFE film. Figure 4.2 displays plots of advancing and receding CA for the different handsheets as a function of oxygen etching time. The graph shows the transition from “sticky” to “roll-off” superhydrophobicity (CA hysteresis $< 10^{\circ}$) after ~ 30 minutes of etching for all substrates. The curves in Figure 4.2 overlap, showing that the rate of change of advancing and receding CAs, which is closely connected to the evolution of nano-scale features, was indistinguishable for the three handsheets within experimental error. Indeed, there were no noticeable differences between the sizes of nano-scale features formed on etched hardwood and softwood fibers (SEM images not shown). These results provide further evidence that there is no significant difference between the nano-scale features formed on fibers of different types. In conclusion, different fiber types do not affect the superhydrophobic behavior provided that the paper making procedures are constant.

4.3.1.2 Effects of paper making:

The pulping process and the paper machine configuration vary from mill to mill in order to optimize paper properties for specific applications [2]. The process involves the

following steps: after wood chips are pulped and bleached, the paper web is formed in the paper machine, after which it undergoes a variety of mechanical treatments (pressing, drying and calendering) before finally being collected on a large roll [4, 5]. All steps of the papermaking process ultimately affect the roughness of the final paper surface. In this section, the impact of the paper making process conditions on imparting superhydrophobicity to paper surfaces is explored. Two different paper types (in addition to laboratory-made handsheets) were used that were fabricated for unique applications: 1) a commercial copy paper (CP), which is moderately hydrophobic to yield good printability and 2) a paper towel (PT), which is extremely hydrophilic to yield good absorptivity.

The SEM images in Figures 4.1d and 4.1e show high and low magnification SEM images of untreated CP and PT, respectively. Of these two samples, copy paper is most similar to the handsheets; the main difference is the presence of (inorganic) filler particles on the fiber surface (shown in Figure 4.1d). The paper towels have a noticeably more porous surface with very loosely cross-linked fibers; these substrates are designed for superior absorption properties. From the SEM images it is evident that these substrates have very different surface roughness values prior to plasma treatment.

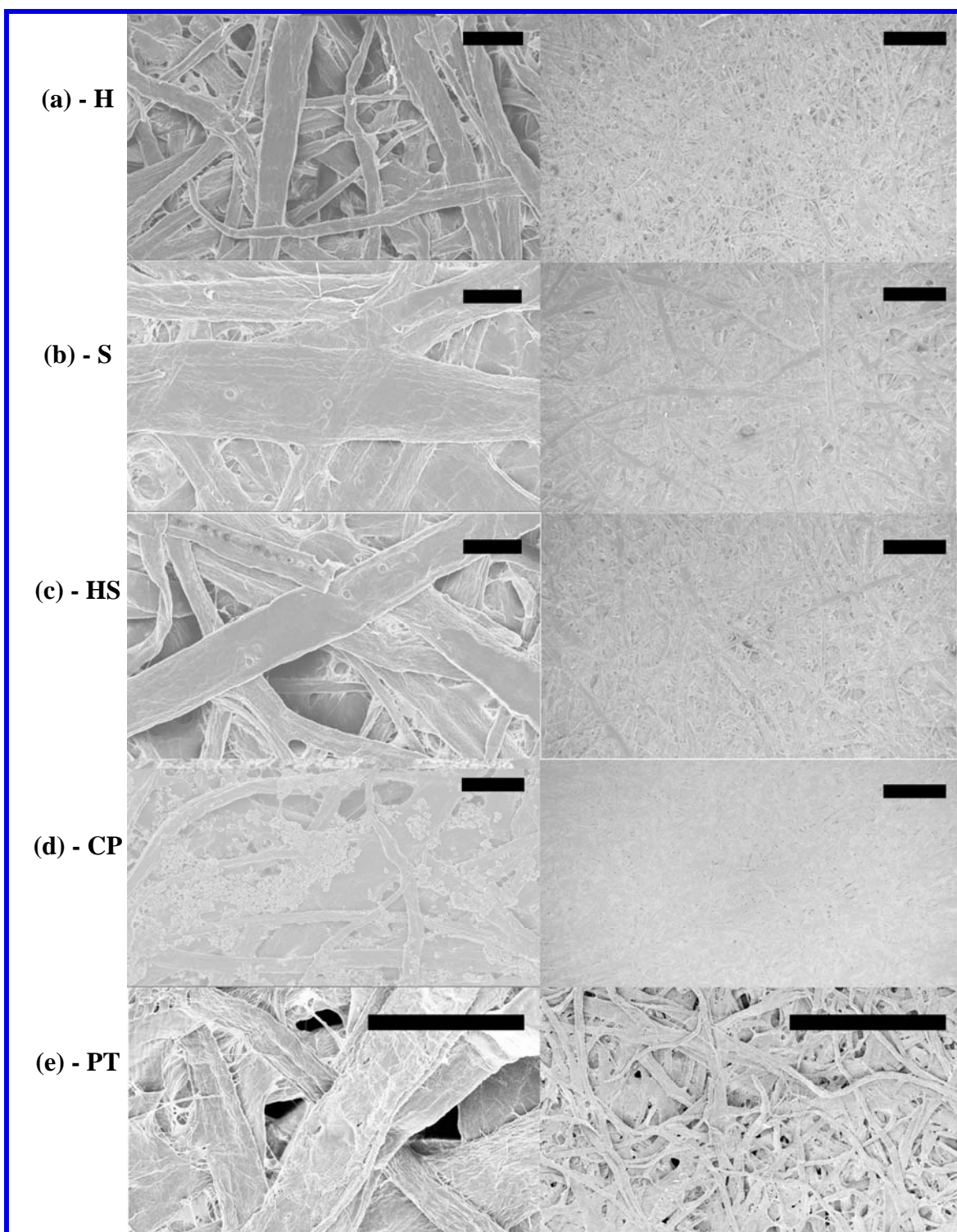


Figure 4.1 Low and high magnification SEM images of laboratory handsheets made with (a) 100% hardwood (H), (b) 100% softwood (S), (c) 50%-50% hardwood and softwood (HS), and two commercial paper sample, (d) copy paper (CP) and (e) paper towel (PT). Scale bars correspond to 40 μm (high magnification) and 4000 μm (low magnification)

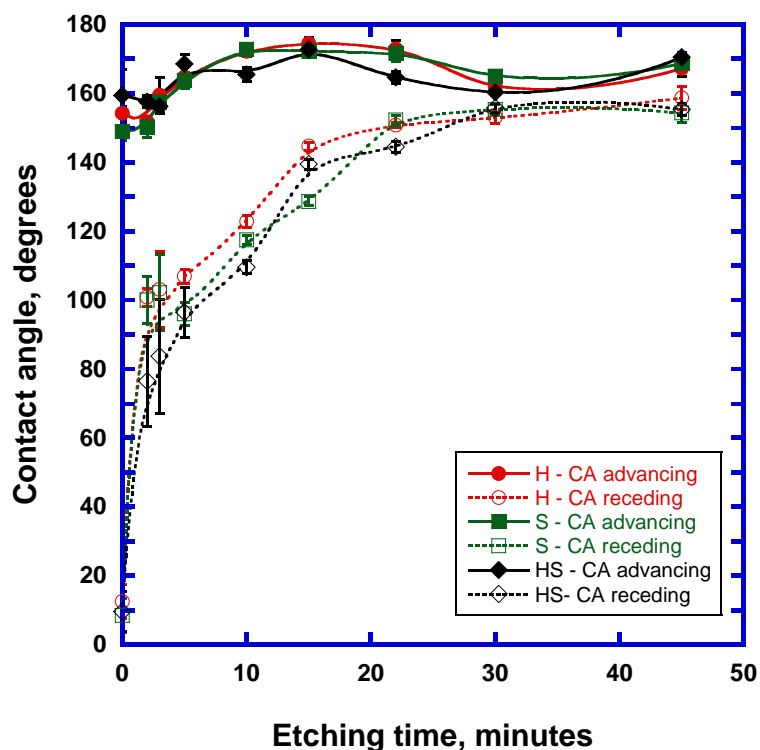


Figure 4.2 Plot of advancing CA and receding CA of handsheets (H, S, HS) with respect to oxygen plasma etching time for 2 min PFE deposition (~100 nm)

The untreated copy paper displayed an advancing CA $\sim 79.2^\circ \pm 3.4^\circ$, which confirms its moderately hydrophobicity. For the untreated paper towel, the water drop was absorbed into the paper within one second, so that CA values could not be measured. After deposition of a thin PFE film (without oxygen etching), the CP and PT substrates yielded different superhydrophobic behavior than the HS substrate, as shown in Figure 4.3. The difference in receding CA can be attributed to differences in the micro- and nano-scale roughness that result from the distinct processing conditions in the paper mills (evident from the SEM images in Figure 4.1). The advancing and receding CA values for CP, which is most similar to the HS handsheet with regards to fiber composition, are analogous to the values obtained for HS. However, the PT showed a very different receding CA as compared to HS and CP. The increased values of the receding CA

(decreased CA hysteresis) of the PT can be attributed to the increased micro-scale roughness associated with the increased porosity of these substrates. In addition, the PT showed different superhydrophobic behavior on the two sides of the substrate (labeled PT-top and PT-bottom in Figure 4.3). Although the SEM images did not reveal a significant difference between the two sides, we believe that the distinct CA values are due to the different roughness scales generated on the felt and wire side of the paper during the manufacturing process, usually referred to as “two sidedness of paper” [2, 6]. The copy paper did not show difference in superhydrophobic behavior between the top and bottom side, which is expected since the applications of copy paper require that it has the same physical and chemical properties on both sides.

The paper substrates (CP, PT-top and PT-bottom) were subsequently etched in an oxygen plasma environment for different durations prior to PFE deposition. The advancing and receding CA of these substrates with respect to oxygen etching times are shown in Figure 4.4. It is evident from Figure 4.4 that “roll-off” superhydrophobic behavior could be obtained for all samples tested, in spite of significant differences in paper making methods. Indeed, the nano-scale roughness established by oxygen etching, which is responsible for the “roll-off” superhydrophobic behavior, was found to be similar for all papers (SEM images not shown).

In conclusion, the difference in CA hysteresis between various substrates (Figure 4.3) results in a difference in the adhesion of water drops on these substrates. This demonstrates that by controlling the paper making processes, adhesion of water drops on a superhydrophobic paper surface can be tuned. Also, after the paper substrates are

etched, the formation of nano-scale roughness dominates the superhydrophobic behavior, thereby leading to similar wettability for all tested paper substrates. Although these experiments do not represent a comprehensive study, they do provide a general picture of the effects of paper making parameters on superhydrophobicity as established by our plasma treatment process.

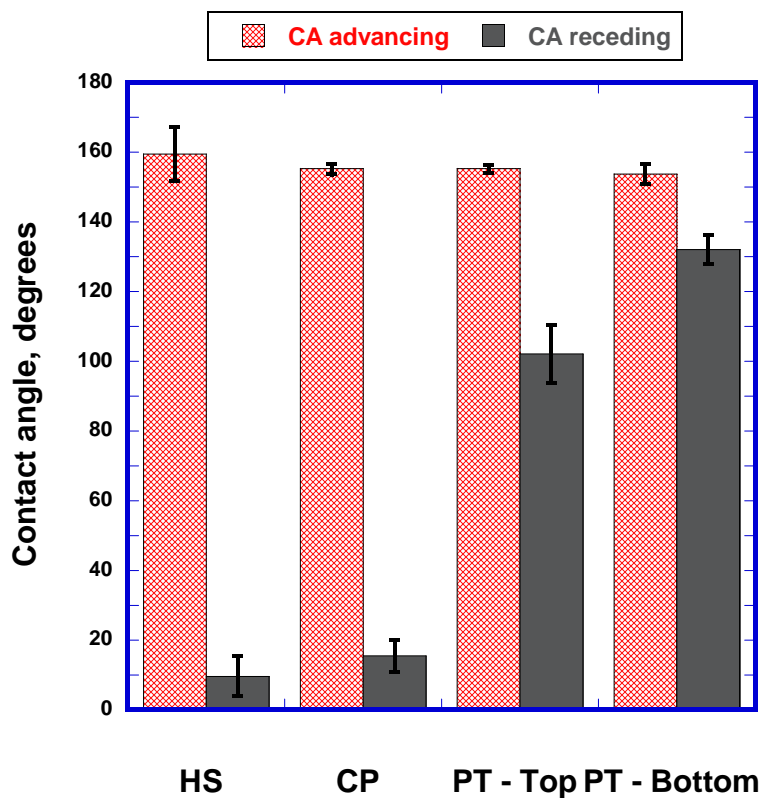


Figure 4.3 Plot of advancing CA and receding CA of handsheet (HS), coppypaper (CP), paper towel-top side (PT-top) and paper towel-bottom side (PT-bottom), after 2 min PFE deposition (~100 nm) and without oxygen etching

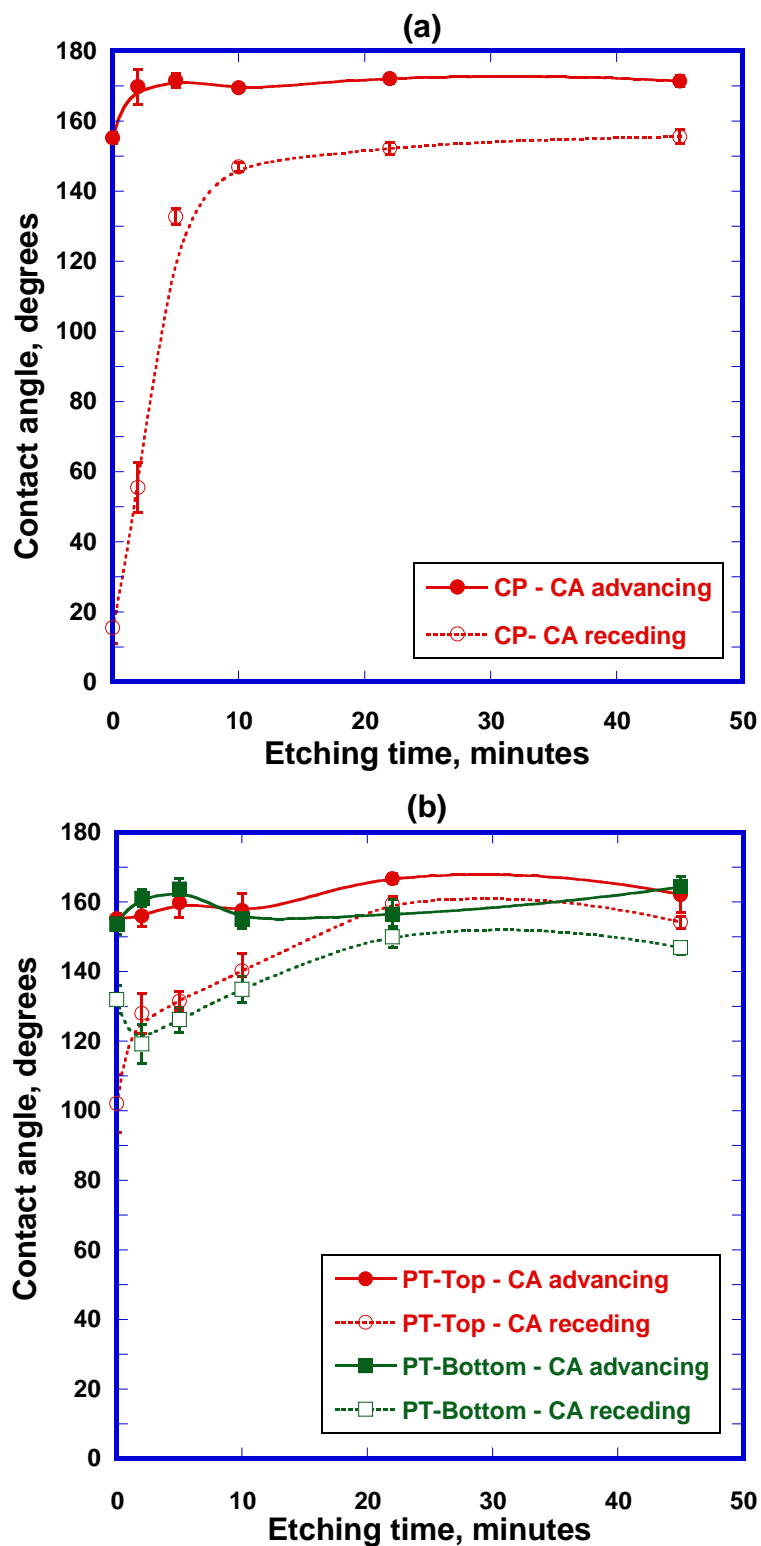


Figure 4.4 Plot of advancing CA and receding CA of copypaper (CP) (a) and paper towel (PT-top and PT-bottom) (b) with respect to oxygen plasma etching time for 2 min PFE deposition

4.3.2 Design of superhydrophobic paper surfaces via control of fiber type and plasma processing conditions

Longer softwood fibers are usually responsible for paper strength, while shorter hardwood fibers are predominantly responsible for paper shininess because of reduced roughness [7]. Experimental results indicate that oxygen etching ultimately reduces the shininess of the paper by creating nano-scale roughness, so that the presence of hardwood fibers no longer provides enhanced optical properties in the etched handsheets. Therefore, fabrication of superhydrophobic paper based on softwood fibers is the most desirable approach because of the expected increase in physical strength.

For longer PFE deposition times, roll-off superhydrophobic behavior can only be achieved after 60 minutes etching (data not shown), while for 2 min PFE deposition times roll-off is observed after much shorter etching times (~30 minutes) (Chapter 2). These results are due to smoothing of the topography of the roughened surface that occurs during the deposition of a thicker PFE film. Since prolonged oxygen etching damages the fiber surfaces, it has a significant negative impact on the strength of the paper, which is undesirable. Therefore, it is most desirable to obtain roll-off superhydrophobicity at reduced etch times. It is expected that an optimum PFE thickness exists: thick enough to retard the absorption of water, yet thin enough to prevent smoothing of the morphology created by oxygen etching. Results suggest that a ~100 nm film obtained from a 2 min PFE deposition is a near-optimum thickness to achieve a roll-off superhydrophobic paper surface with good physical properties. Of course, optimizing the papermaking process to tune the micro-scale roughness may offer an additional degree of freedom for the design of superhydrophobic paper surfaces.

4.3.3 Discussion of the effect of drop size on contact angle and edge geometry

Measurement of contact angles on rough surfaces is more complex when the drop sizes are comparable to the roughness length scale of the substrate. Advances in drop dispense technologies have made it possible to vary the dispensed volume of a water drop in a controlled manner from a few picoliters to a few microliters. This section describes the significance of the water drop size when measuring CAs on paper surfaces.

Figure 4.5 shows the contact lines established by water drops of four different volumes on a HS handsheet with CA characteristics shown in Figure 4.3. Solid lines were drawn along the three-phase contact lines to highlight the contact line geometry. Clearly, the contact line is more distorted by the topography of the fiber network for smaller drops (Figures 4.5a and 4.5b) than for larger drops (4.5c and 4.5d).

In order to determine the effect of contact line distortion on CA measurement, the drop volume was varied from 0.1 μL to 16 μL , and the advancing CA values were measured. Figure 4.6 shows the advancing CA as a function of drop volume for HS substrates etched for three different etching durations. Figure 4.7 shows the images of water drops corresponding to a 30 min etched HS (“roll-off” superhydrophobic) and 0 min etched HS (“sticky” superhydrophobic). From the data in the graph it can be concluded that the advancing CA increases significantly up to a volume of $\sim 2 \mu\text{L}$. This suggests a lower limit of drop volume that should be used to measure CA on superhydrophobic paper surfaces. On the other hand, the upper limit is determined by the angular resolution of the CA goniometer. As the drop volume increases, the drop flattens due to gravity, which makes it difficult to locate the three phase contact point from a side view of the drop. This limitation can greatly affect the accuracy of CA values, which is consistent with the

observed slight decrease in advancing CA values for the 16 μL drop (Figure 4.6). Hence, in order to mitigate ambiguity in CA values when measuring contact angles on porous, heterogeneous substrates such as paper, it is important to select drop sizes that 1) are larger than the surface roughness length scales and 2) provide sufficient image resolution for the goniometer to identify the three-phase contact line.

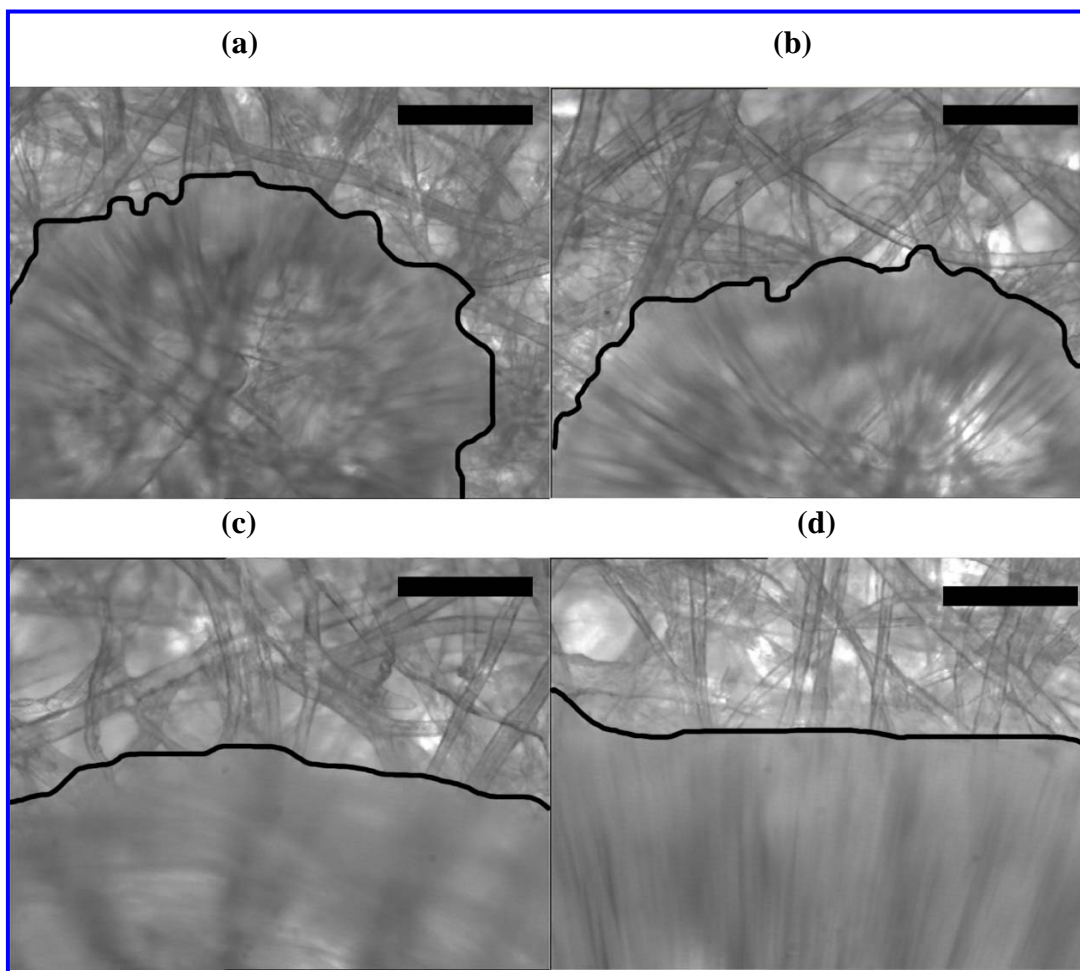


Figure 4.5 Contact lines formed by 0.1 μL (a), 0.2 μL (b), 4 μL (c) and 8 μL (d) water drops on a 2 min PFE deposited (without etching) HS substrate. Scale bars correspond to 160 μm

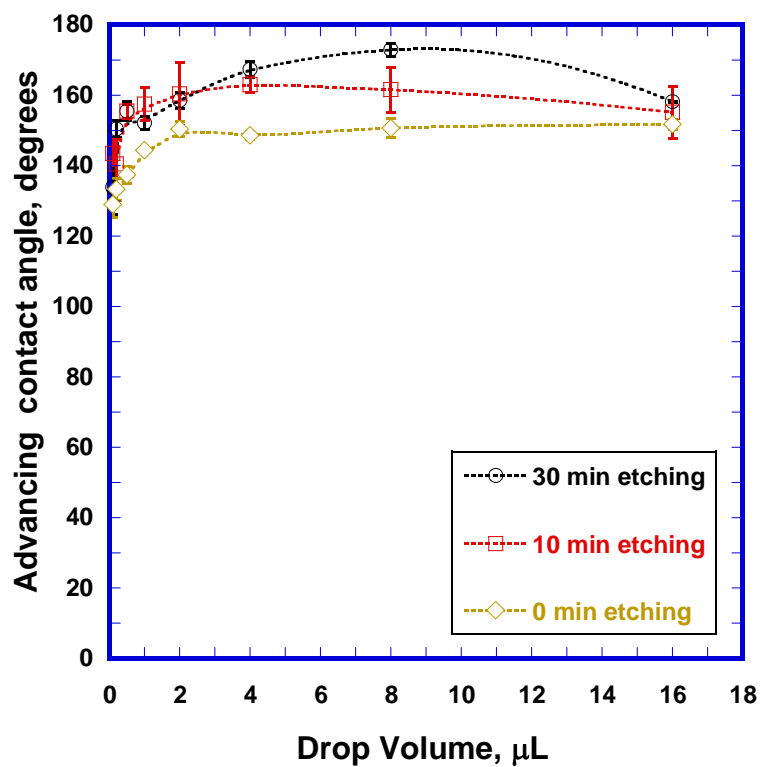


Figure 4.6 Plot of advancing CA with respect to drop volume for oxygen etched (0 min, 10 min and 30 min) and PFE deposited (2 min) handsheet (HS) surfaces

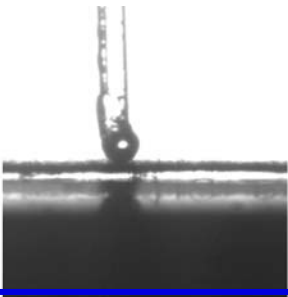

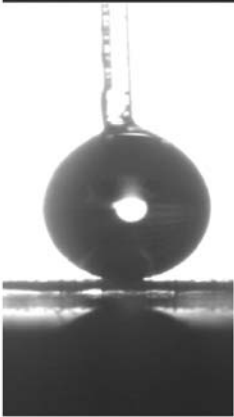
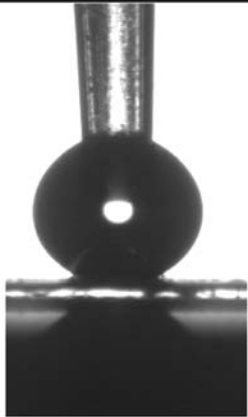
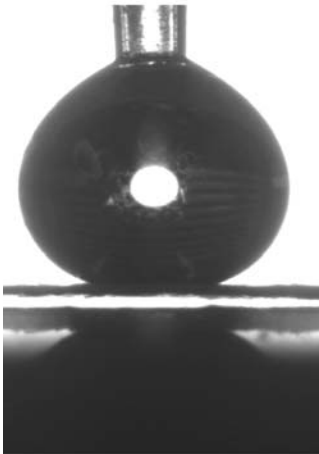
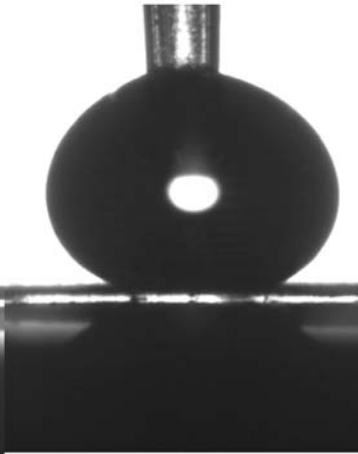
	“Roll-off” Superhydrophobic HS	“sticky” Superhydrophobic HS
0.1 μL		
4 μL		
8 μL		

Figure 4.7 Photographs of advancing CA for different drop volumes for “roll-off” superhydrophobic (0 min oxygen etched and 2 min PFE deposited) and “sticky” superhydrophobic (30 min oxygen etched and 2 min PFE deposited) handsheet (HS) surfaces

4.4 Conclusions

The effects of fiber types and paper making parameters on the creation of superhydrophobic paper surfaces were studied. While the fiber type does not significantly affect superhydrophobicity, the paper making process plays a significant role in determining the microscale roughness, and thus the superhydrophobic behavior of paper substrates. Paper made from softwood fibers is likely to be more suitable for superhydrophobic applications because of improved physical properties with this fiber type, in particular paper strength. A PFE film of $\sim 100\text{nm}$ PFE represents a near-optimum thickness to obtain superhydrophobicity. The importance of water drop volume in the measurement of CAs on superhydrophobic surfaces fabricated on heterogeneous and porous substrates such as paper has been discussed. It was concluded that drop sizes that are larger than the surface roughness length scales and provide sufficient image resolution for the goniometer to identify the three-phase contact line are critical in obtaining unambiguous results while measuring CA values for porous substrates.

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CHAPTER 5

PATTERNING OF SUPERHYDROPHOBIC PAPER TO CONTROL THE MOBILITY OF MICRO-LITER DROPS FOR TWO- DIMENSIONAL LAB-ON-PAPER APPLICATIONS*

5.1 Introduction

This chapter describes patterning methods that provide local control over droplet adhesion on superhydrophobic paper and an approach to develop novel paper-based LOC microfluidic devices that enable manipulation (storage, transport, mixing and sampling) of drops of test fluids on the substrate, without absorption of these fluids into the porous paper. These two-dimensional devices can be used for qualitative analytical fluid testing, as well as storage of large arrays of drops for transportation and further quantitative analysis.

In the early stages of their development, LOC microfluidic devices were fabricated with technologies originally developed for the microelectronics industry, in particular photolithography and etching, and thus were fabricated from silicon wafers or glass substrates [1]. Subsequently, researchers began investigating polymers as substrates (especially PDMS) in combination with soft lithography techniques because of the

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advantages of these substrates over silicon- or glass-based devices: transparency, flexibility, biochemical compatibility and permeability [1-4]. However, even PDMS-based devices require the use of clean room facilities for the fabrication and incorporation of complex components such as valves, pumps and mixers [1, 5-7]. Fluid actuation in these devices relies mostly on electrokinetic or pneumatic actuation, which requires an external power source (high voltage power supply, batteries, or compressed gas/vacuum sources [8, 9]). Overall, in spite of breakthrough advances in LOC concepts, most of the devices remain unsuitable for low-tech applications like biomedical diagnostics in developing countries due to the lack of simplicity and affordability.

Paper-based LOC devices (also referred to as lab-on-paper (LOP) [10]) have emerged as a promising alternative technology. For fluid actuation on these devices one can rely on capillary forces inside the porous paper and thus avoid external power sources. In a recent report on the top ten biotechnologies for improving health in developing countries “modified molecular technologies for affordable, simple diagnosis of infectious diseases” was ranked as the number one priority [8, 11]. Another report on the grand challenges for global health ranked the development of technologies to “measure disease and health status accurately and economically in poor countries” first among the top 14 priorities [8, 12]. Due to their affordability and potentially simple fabrication technology, LOP devices may offer improved global availability of medical technology.

In its simplest form, the concept of LOP dates back to the 1950s, when paper-based strips [13-18] were first used for biomedical diagnostics. However, applications of these LOPs

were limited by the fact that they could not perform multiplex analysis: *i.e.*, it was impossible to perform multiple biochemical analyses on a single sample with the same strip. This limitation inspired the fabrication of multiple channels with barriers within a paper substrate, analogous to a microfluidic device, to enable multiplex analysis. Creation of hydrophilic channels with hydrophobic barrier layers for biochemical assay devices was originally proposed in 1995 and 2003 [19, 20]. More recently, this concept has been adapted by using modern photolithography techniques to create hydrophobic photoresist barriers [21-23]. This work has since been expanded to three-dimensional LOP devices by layering sheets of patterned paper with perforated barrier tape to guide the exchange of liquids between paper layers [24]. A disadvantage of these LOPs was the limited flexibility due to the use of rigid photoresists (SU-8 or PMMA), which has been addressed by printing PDMS as a barrier polymer using a desktop plotter, thus creating flexible LOP devices [25]. However, the low surface tension of uncrosslinked PDMS limits the spatial resolution of the patterns, resulting in broad and irregularly shaped barrier wall structures [25, 26]. A new two-step method for patterning straight barrier walls was proposed: hydrophobize the entire paper substrate with Alkyl Ketene Dimer (AKD) and then create hydrophilic channels via a plasma patterning process [26]. Although both PDMS- and AKD-based LOPs are flexible, the channels are relatively wide (1-2 mm) because of the patterning limitations [25, 26]. Controlled fabrication of channels with widths of several hundred micrometers has been achieved by printing hydrophilic patterns via inkjet printing [27]. The use of widely available technology to design LOP devices, for example a standard desktop printer, clearly offers substantially enhanced versatility, since it enables end-users to “program” LOP devices according to

specific needs. A recent report has noted that programmable LOCs would be the next critical innovation in this technology [28]. Most current LOP technologies limit the ability of non-expert users to program their own devices because of the complex chemicals, methods, and/or equipment needed for device fabrication. Furthermore, all the LOP concepts discussed above depend on absorption of test fluids into the hydrophilic areas of porous paper and use capillary forces for fluid actuation. As a result, the products of reactions occurring inside a LOP cannot easily be extracted for further biochemical analysis. This is particularly important because the analysis in LOPs is currently semi-quantitative at best; the accuracy and sensitivity cannot compete with traditional analytical equipment [10].

One option to overcome this issue is to prevent absorption of the liquids into the paper matrix. By restricting droplets to the surface of the substrate, the samples are accessible for post-processing and quantitative analysis in a centralized testing center [8], while simple qualitative biochemical characterizations can still be performed at the point-of-care (POC). In order to achieve this, droplets must be manipulated on a two dimensional substrate that enables basic unit operations: storage, guided transport, mixing and sampling. 2D microfluidic lab-on-chip devices have been previously obtained via electrowetting [29, 30] and optoelectrowetting (OEW) [31-33], but these approaches require external power sources for operation and complicated fabrication methods. Ideally, a 2D LOP should be inexpensive, enable design flexibility and operate without an external power source.

An approach to develop a 2D LOP device capable of storage, transfer, mixing and sampling of liquid drops by decorating superhydrophobic paper substrates with high surface energy ink patterns (lines and dots) is discussed in this chapter. Surface energy and gravitational forces are used to manipulate and transfer drops, thus eliminating the need for an external power source. The key feature of this device is that patterning changes the local contact angle (CA) hysteresis, resulting in sticky ink spots on non-sticky superhydrophobic paper; the substrates are therefore referred to as Hysteresis Enabled Lab-on-Paper (HELP) substrates. This study demonstrates that patterns to manipulate microliter drops can be designed using standard word processing software and a commercially available desktop printing process that deposits waxy inks. The simplicity of the soft- and hard-ware ensures that end-users can readily develop their own patterns to achieve desired functionality of the LOP devices [28]. Finally, the HELP substrates can serve as an inexpensive storage medium for test fluids, reagents and/ or reaction products in the form of arrays of drops, which can then be transported to the centralized testing centers for detailed quantitative analysis after initial semi-quantitative on-chip analysis.

This chapter demonstrates the fundamental principles behind our LOP concept by measuring and modeling the adhesion of water drops on patterned substrates, and applies this basic knowledge to the design of LOP building blocks with advanced functionality, such as droplet storage, transfer, merger, mixing and sampling.

5.2 Experimental

5.2.1 Superhydrophobic paper

Handsheets were used as model paper substrates and were fabricated following TAPPI-standardized protocol T205 sp-02, using southern hardwood kraft (Alabama River Pulp Co.) and southern softwood kraft (North Carolina International Paper). Handsheets were placed inside a 13.56 MHz parallel plate plasma reactor to undergo a two step process (oxygen etching for 60 min to generate roughness and fluorocarbon (Pentafluoroethane monomer) film deposition for 1 min to establish surface hydrophobicity) that results in “roll-off” superhydrophobicity as described previously (Chapter 2). A more detailed discussion regarding handsheet preparation and plasma processing can be found Section 2.2.1 and 2.2.2 (Chapter 2) respectively.

5.2.2 Patterning

The patterns were designed using standard word processing software (Microsoft® Word 2007). Two types of simple patterns were used: dots and lines. The size of the dots and lines were varied using the font size in “pt”- units, as provided by Microsoft® Word 2007. The “roll-off” superhydrophobic handsheets were pasted on sheets of regular copy paper using Scotch® tape and fed through a Xerox® Phaser 8500n printer to print the patterns designed in the word processing software onto the superhydrophobic handsheet with black phaser ink. Brightfield microscopy images of the printed substrates (Leica microscope DM4500 B) were used to determine the conversion factor of the patterns from pt units to μm ; images were analyzed using Image J software. For dots, the

conversion factor for the dot diameter was found to be 118.5 $\mu\text{m}/\text{pt}$; the conversion factor for line width was 404.9 $\mu\text{m}/\text{pt}$.

5.2.3 Water contact angle measurements

Water contact angle measurements were obtained with a Rame-Hart contact angle goniometer (model 100, Netcong, NJ). Advancing and receding contact angles were measured by placing a drop of known volume on the substrate and dragging the paper substrate left to right with respect to the drop; a more detailed description of this method can be found in section 2.2.7 (Chapter 2). Values of the advancing and receding CAs of non-patterned SH paper (after passage through the printer) are $\theta_{ASH}=165.1\pm2^\circ$ and $\theta_{RSH}=135.3\pm2.9^\circ$; for a substrate with full coverage of the ink film, $\theta_A=113.8\pm2.7^\circ$ and $\theta_{RI}=84.7\pm2^\circ$.

5.2.4 Sliding drop experiments

The substrates were mounted on a flat surface attached to a rotating optical stage. The plate was tilted gradually until the drop rolled-off. The angle (in degrees) at which the drop started to slide was defined as the critical sliding angle.

5.2.5 Drop transfer experiments

The drop was placed on a horizontally placed paper substrate which had the “from” pattern (dot). Another substrate having the “to” pattern (line or dot) was then inverted and manually aligned to the drop, to obtain a configuration in which the drop touched the ink

patterns on both substrates, which were placed parallel to each other. Then the “to” substrate was carefully lifted in a direction perpendicular to the “from” substrate and the resulting drop dynamics was recorded.

5.3 Results and Discussion

5.3.1 Adhesion on patterned paper

5.3.1.1 Sliding drops on sticky islands

In the middle of the 20th century, four research groups independently reported that for a drop sliding on a homogeneous surface, the ratio of the force exerted on the drop (F) to the width of the drop perpendicular to the direction of sliding (W_{drop}) is constant [34-37]:

$$\frac{F}{W_{drop}} = \frac{\rho V g \sin \alpha}{W_{drop}} = K_1 \quad (5.1)$$

where ρ is the density of the liquid drop, V is the volume of the drop, g is the acceleration due to gravity, and α is the critical sliding angle. The constant K_1 was then related to the work functions associated with wetting ($\gamma_{LV} (1 + \cos \theta_A)$) and dewetting ($\gamma_{LV} (1 + \cos \theta_R)$) of the substrate by the drop [35, 36]:

$$\frac{\rho V g \sin \alpha}{W_{drop}} = \gamma_{LV} (\cos \theta_R - \cos \theta_A) \quad (5.2)$$

where γ_{LV} is the surface tension of the liquid, and θ_A and θ_R are the advancing and receding contact angles of the drop on the surface. This semi-empirical equation, hypothesized by Bikerman [34] in 1950 and derived by Kawasaki [36] in 1960, is popularly known as the Furmidge equation [35], in reference to the researcher who reiterated it in 1962. This equation is based on a force balance calculation on the receding and advancing edges of a 2D drop sliding on an inclined surface. For a 3D drop, the contact angle varies continuously along the three-phase contact line, which complicates mathematical analysis. Currently, some disagreement exists in literature as to whether the local contact angles at the advancing and receding edge of the sliding drop are equal to the experimentally measured maximum (θ_A) and minimum (θ_R) contact angles [38]. In spite of this ambiguity, it has been demonstrated that the Furmidge equation is a good empirical approximation for 3D sliding drops [39] and the Furmidge equation is used by most researchers [38]. After rearranging Eqn. 2 we obtain:

$$\frac{V \sin \alpha}{W_{drop}} = \frac{\gamma_{LV} (\cos \theta_R - \cos \theta_A)}{\rho g} \quad (5.3)$$

For a specific liquid-surface combination, the right hand side of Equation 5.3 is constant. On homogeneous substrates, the volume (V) and width (W_{drop}) of a drop are typically connected via simple geometrical relations, so that each drop size corresponds to a unique sliding angle α . This work aims at disrupting this one-to-one correspondence with the objective to independently control the critical sliding angle (α) and the drop volume (V). Eqn. 3 suggests one possible route to achieve this: by making the drop width (W_{drop})

independent of the drop volume (V). The following paragraphs will demonstrate how this can be accomplished.

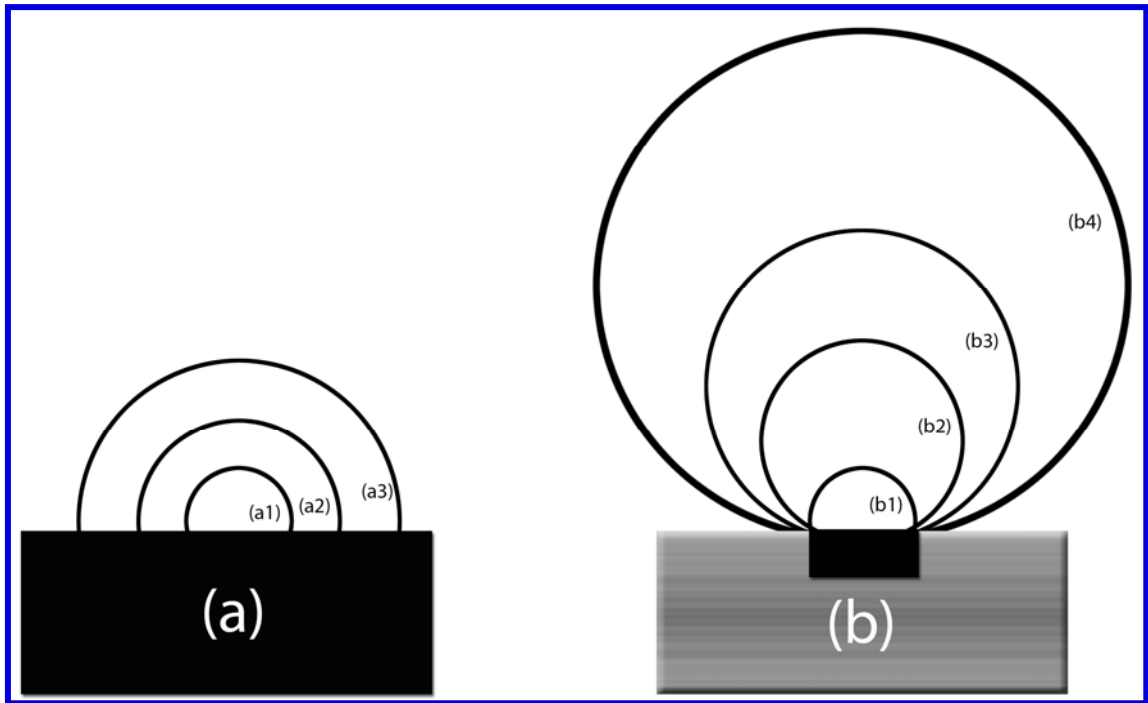


Figure 5.1 Schematics of side view profiles for various drop volumes (a) on a homogenous hydrophobic ($CA \sim 90^\circ$) surface and (b) on a superhydrophobic ($CA > 150^\circ$) surface with hydrophobic ($CA \sim 90^\circ$) pattern (b)

For simplicity, we will assume that the liquid in the following description is water. Figure 5.1a shows a schematic of the side view of drops with various volumes dispensed on a homogeneous hydrophobic surface (advancing contact angle $\sim 90^\circ$). When the drop volume increases, the width of the drop also increases in order to maintain a constant contact angle on the substrate. Next, consider the patterned substrate shown in Figure 5.1b, where a hydrophobic island (advancing contact angle $\sim 90^\circ$, same as in 1a)) is surrounded by a superhydrophobic surface (advancing contact angle $> 150^\circ$) which is extremely water repellant. In this case, when more liquid is added to the drop, it does not expand its contact line periphery onto the superhydrophobic substrate until the advancing

contact angle of the surrounding superhydrophobic substrate is reached ($> 150^\circ$). As a result, the drop width initially remains constant, while the contact angle changes: the width of the contact area between drop and substrate (W_{drop}) is equal to the size of the sticky island and independent of the volume (V). Only for sufficiently large drops (b4 in Figure 5.1b), when the advancing contact angle of the surrounding superhydrophobic substrate is reached, will the base of the drop expand beyond the sticky island. In conclusion, with such patterned substrates, the critical sliding angle (α) at constant drop volume (V) can be manipulated by changing the dimensions of the sticky island (Equation 5.3).

The patterned substrates were experimentally obtained as follows: “roll-off” superhydrophobicity was first achieved on the paper substrates using plasma etching and deposition. The hydrophobic island on the superhydrophobic surface was then obtained by printing “•” (from the symbol menu in Microsoft® Word 2007; designated “dots” in the remainder of this manuscript) using a commercially available phaser printer (Xerox 8500n) and standard black ink. The difference in advancing contact angle between superhydrophobic paper ($165.1 \pm 2^\circ$) and homogeneous full-coverage films of black phaser ink ($113.8 \pm 2.7^\circ$) is sufficient to create the scenario depicted in Fig. 1b. The superhydrophobic paper substrate was robust enough that the bending and pressing of the printing process did not affect its advancing CA. Figure 5.2 shows the critical sliding angle (angle at which the drop started to slide) versus drop volume for different dot sizes along with the sliding angles of two homogeneous control substrates (blank superhydrophobic paper after passage through the printer (SH) and a full coverage ink

film printed on the superhydrophobic paper surface (I)). Data points on the line that marks the critical slide angle of 90° represent the largest drop that did not slide from vertical substrates. Drop behavior on the control substrates (SH and I) was in good agreement with predictions from a modified Furmidge equation, as will be discussed later in this section. As expected, for each substrate, the critical slide angle decreases monotonically with increasing drop volume (within experimental error). When comparing the results for patterned substrates with the homogeneous control substrates, it is evident from Figure 5.2 that at a constant drop volume (V), the critical slide angle (α) increases with increasing dot width. Another interesting observation is that patterned surfaces with large dot sizes (e.g., 2.37 and 4.22 mm) require larger slide angles (*i.e.*, surfaces are more sticky) than a homogenous ink substrate.

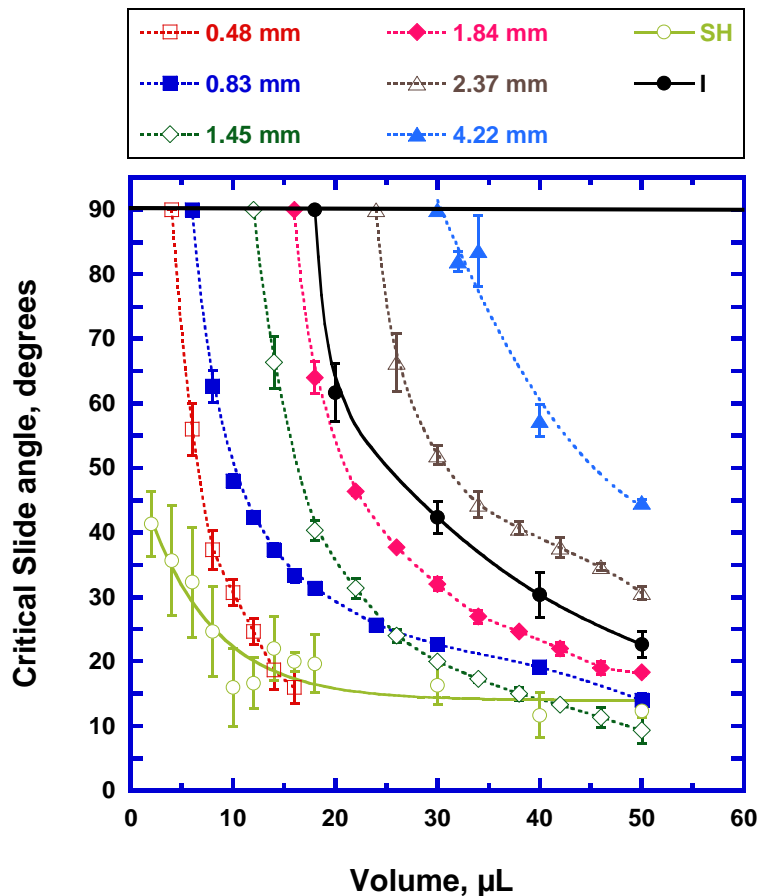


Figure 5.2 Critical slide angle versus drop volumes on patterned substrates (for various dot sizes) and control substrates (SH and I); curves are to guide the eye

Since a continuous ink film essentially is a dot with infinite width, one might have expected that α for the ink film would be higher than for all printed dots. However, this apparent anomaly can be explained using the Furmidge equation. For a drop to slide on a surface, it must deform so that the advancing and receding edge of the drop both reach the experimentally measured advancing and receding CA, respectively, for that substrate. If the drop slides from a printed dot, the advancing CA is set by the superhydrophobic paper, while the receding CA is that of the ink film. In contrast, for homogeneous substrates both advancing and receding contact angles are for the same surface material. For experiments with the smaller dots (< 2.37 mm) the initial CA of the drop after

dispensing it onto the horizontal substrate was essentially equal to the advancing CA of the superhydrophobic paper surface (similar to configurations b3 or b4 in Figure 5.1b), because the drops are large relative to the dot. For the bigger dots, however, the initial CA was closer to the advancing angle on the ink film (configuration b1 in Figure 5.1b). Hence, on larger dots a drop must deform to a greater extent before its advancing CA reaches $\sim 165.1 \pm 2^\circ$ and the drop starts to slide; thus resulting in higher critical sliding angles for the bigger dots. Figure 5.2 demonstrates clearly that the critical sliding angle for drops of any size can be tuned by adjusting the size of the dots printed on the superhydrophobic surface.

The data in Figure 5.2 can be interpreted quantitatively by inspecting the Furmidge equation (Eqn. 5.2) more closely. This equation essentially represents a force balance,

$$F_E [= \rho V g \sin \alpha] = F_P [= W_{drop} \gamma_{LV} (\cos \theta_R - \cos \theta_A)] \quad (5.4)$$

where F_E is the experimentally measured gravitational force that is necessary to slide a drop on a surface and F_P is the force that can be predicted theoretically from the values of W_{drop} , θ_A and θ_R , which can be determined via independent experiments.

Based on Figure 5.1b, the width of the drop (W_{drop}) should be equal to the width of the dot (W_{dot}) for a wide range of drop sizes. It was observed experimentally that for large drops, gravity deformed the drops sufficiently to extend the contact line of the drop beyond the dot width (W_{dot}) as shown in Figure 5.1b. This is denoted as the “outside”

configuration (b4) and the corresponding predicted force as F_{PO} . Once the contact line of the drop extends beyond the ink periphery, the surface energy of the ink film no longer affects the size of the contact area, W_{drop} , or the contact angle θ_A [40-43]. Thus for the F_{PO} configuration, these parameters are determined solely by the properties of the superhydrophobic paper and W_{drop} can therefore be obtained independently by measuring the drop width on non-patterned superhydrophobic paper substrates as a function of drop volume. The results from these experiments (data not shown) were used to calculate W_{drop} for any drops for which $W_{drop} > W_{dot}$. When the substrate in this “outside” drop configuration is tilted, there are two contributions to the adhesion force F_{PO} : one from the part of drop in direct contact with the ink dot (advancing CA of paper substrate (SH) and receding CA of ink (I)) and another from the part of the drop only in contact with the superhydrophobic paper (both advancing and receding CA of SH). Assuming that these force contributions are additive, the predicted force F_{PO} can then be modeled as:

$$F_{PO} = [W_{dot}\gamma_{LV}(\cos\theta_{RI} - \cos\theta_{ASH})] + [(W_{drop} - W_{dot})\gamma_{LV}(\cos\theta_{RSH} - \cos\theta_{ASH})] \quad (5.5)$$

where θ_{ASH} is the advancing CA of SH paper, θ_{RI} and θ_{RSH} are the receding CAs of the ink film and SH paper, respectively.

If, on the other hand, the drop is confined to the perimeter of the dot (configurations b1-3 in Figure 5.1b), the force F_{PP} needed to slide the drop depends only on a single length scale, the dot size (W_{dot}). For this configuration, the second term on the right hand side of Eqn. 5 disappears ($W_{dot} = W_{drop}$), resulting in:

$$F_{PP} = [W_{dot} \gamma_{LV} (\cos \theta_{RI} - \cos \theta_{ASH})] \quad (5.6)$$

The scenario for which $W_{drop} < W_{dot}$ was not encountered with the dot sizes and drop volumes in this study. For homogeneous control substrates (I and SH) there is no need to modify the original Furmidge equation (Eqn. 5.3), provided that the appropriate CA values are used to calculate the adhesion force (F_I for ink film and F_{SH} for SH paper).

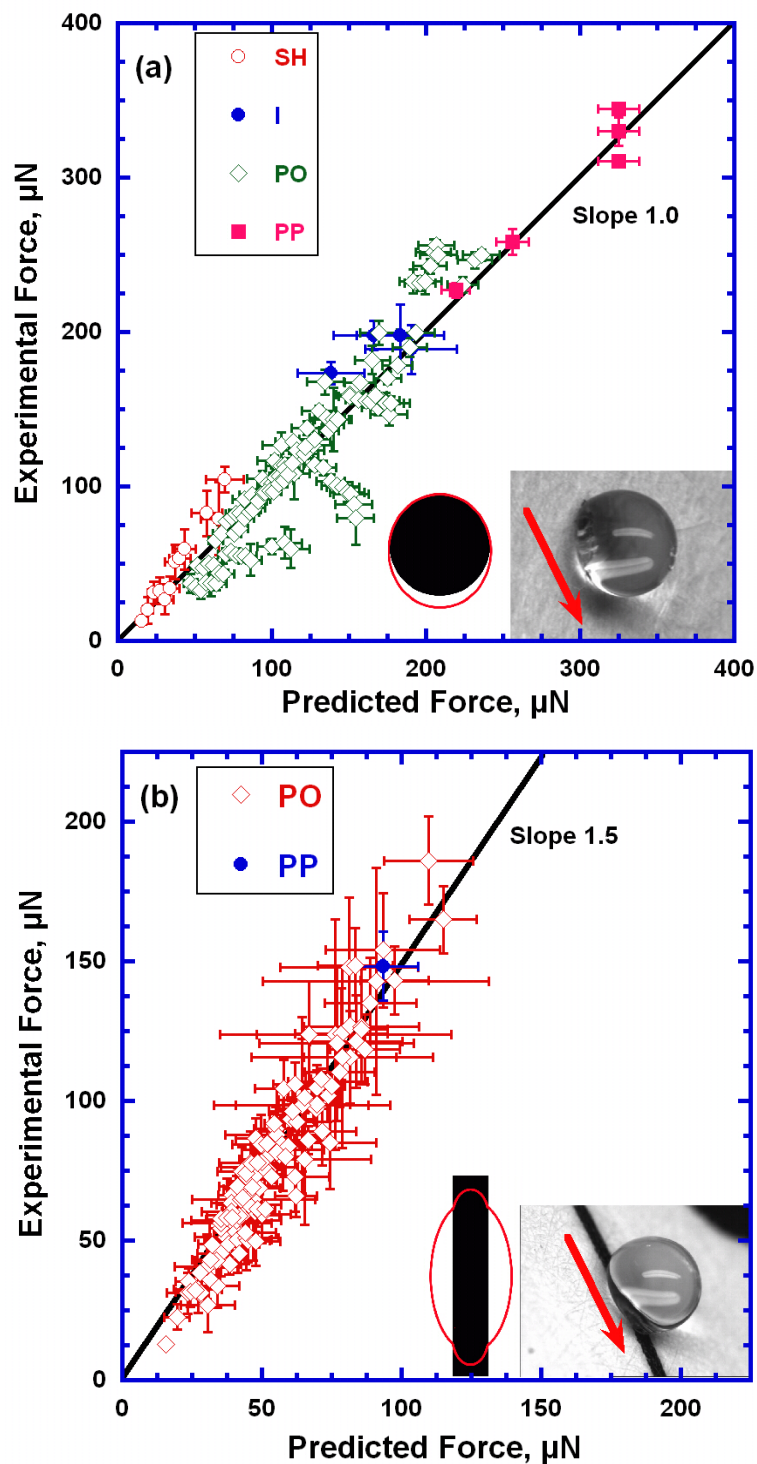


Figure 5.3 Experimental vs predicted drag-adhesion force for dots (a) and lines (b) for the following substrates: Superhydrophobic paper after printing a blank pattern (SH), ink film on a SH paper (I), configuration in which the contact line of the drop is outside the ink pattern's periphery (PO) and configuration in which the contact line of the drop is on the ink pattern's periphery (PP). (Insets: Schematic of contact line profile compared to

the pattern geometry and photograph of a 4 μL drop just before sliding on a 0.83 mm dot (a) and a 0.3 mm line (b))

Figure 5.3a shows a plot of experimentally determined sliding force F_E versus predicted values $F_{PO}/F_{PP}/F_I/F_{SH}$ for a range of drop volumes and dot diameters (0.36 to 4.22 mm, or 3 to 36 pt) including the data presented in Figure 5.2. It is evident that the data correspond quantitatively to the predictions from the modified Furmidge equation, which is based on simple geometrical arguments and has no adjustable parameters. The sliding angle measurements were performed manually, so that slight vibrations induced during the measurements could not be avoided. Also, the SH paper and ink film are both heterogeneous with respect to topography. The deviations from the model in Figure 5.3a for some substrates are believed to be a result of these inevitable experimental errors.

The experiments were subsequently extended to a significantly more complicated ink pattern: lines. For these tests, lines were generated using Microsoft® Word 2007 and printed on superhydrophobic paper. The sliding behavior of drops along the printed lines for different line widths (0.10 to 2.83 mm or 0.25 to 7 pt) was investigated. It must be noted that the motion of drops on line patterns is anisotropic (parallel vs. perpendicular to the line). The initial experiments focused entirely on drop sliding parallel to the lines. In this case, the Furmidge model for the adhesion force for drops that extend outside the line (F_{PO}) can be expressed as

$$F_{PO} = [W_{Line}\gamma_{LV}(\cos\theta_{RI} - \cos\theta_{AI})] + [(W_{drop} - W_{line})\gamma_{LV}(\cos\theta_{RSH} - \cos\theta_{ASH})] \quad (5.7)$$

where the only difference between Eqns 5.5 and 5.7 is that the advancing CA in the first term of equation (5.7) is now the advancing CA of the ink; as the drop slides along the line, the part of the drop that resides on the line always remains in contacts with the ink film. Similar to dots, if the drop is contained within the line, $W_{drop} = W_{Line}$ and the second term of the equation vanishes so that

$$F_{PP} = [W_{Line} \gamma_{LV} (\cos \theta_{RI} - \cos \theta_{AI})] \quad (5.8)$$

Figure 5.3b plots experimental sliding force versus predicted force (F_{PO} and F_{PP}) for the line patterns. Although a quantitative correlation between model and experiments is obvious, one-to-one correspondence was not observed. The experimental forces always exceeded the model prediction and linear least-square regression yielded a simple correction factor of 1.5, as indicated by the line in Figure 5.3b. Although we have no quantitative explanation for this correction factor, the inset of Figure 5.3b clearly shows its qualitative origin: the complicated geometry of the contact line of a drop on a line. The fact that the correction factor is larger than unity can be interpreted as an enhancement of the length of the contact line, which can be attributed to the curvature of the contact line induced by the printed line (see inset in Figure 5.3b).

5.3.1.2 Transfer of drops between substrates

The experiments and models in the previous section give excellent insight into the ‘drag adhesion’ of drops sliding on substrates patterned with ink dots and lines. This section focuses on the force of adhesion that is observed when drops are pulled-off perpendicular

to the patterned substrates. This kind of adhesion is referred to as ‘extensional adhesion’. The objective of these experiments was to capitalize on differences in extensional adhesion between different patterned substrates to permit transfer of drops between substrates.

In 1896, Dupré rearranged Young’s classical contact angle equation to describe the work of adhesion for a drop to detach from a surface [44, 45]:

$$W_{adh} = \gamma_{LV} (1 + \cos \theta) \quad (5.9)$$

where θ is the equilibrium CA of the liquid drop on the surface. At that time, the scientific community did not define maximum (advancing) and minimum (receding) CAs [46]. When a drop detaches from a surface in a direction perpendicular to the plane of the surface, the contact line of the drop experiences a receding CA value rather than the equilibrium CA [47]. As a result, the above equation must be adapted to

$$W_{adh} = \gamma_{LV} (1 + \cos \theta_R) \quad (5.10)$$

This work of adhesion can be converted to the force of adhesion by multiplying the right hand side of the equation by a characteristic length scale, L_{char} :

$$F = L_{char} \gamma_{LV} (1 + \cos \theta_R) \quad (5.11)$$

It is reasonable to assume that this length scale is proportional to the characteristic size of the ink pattern (dot diameter or line width). Hence, Equation 5.11 becomes:

$$F = (\beta W) \gamma_{LV} (1 + \cos \theta_{RI}) \quad (5.12)$$

Where β is the proportionally constant and W is the width of the ink pattern. Therefore, it is expected that a drop positioned on a small ink island will experience a smaller force of adhesion than the same drop sitting on a relatively large ink island. This fact inspired us to determine whether this difference in adhesion force can be used to overcome gravity and transfer a drop from a substrate with a smaller ink island to a substrate with a larger ink island. The resulting force balance between gravity on the drop and adhesive forces of the two substrates is:

$$\rho V g = \beta_1 W_1 \gamma_{LV} (1 + \cos \theta_{RI}) - \beta_2 W_2 \gamma_{LV} (1 + \cos \theta_{RI}) \quad (5.13)$$

where the subscripts 1 and 2 denotes the large and small ink islands, respectively. By rearranging the above equation, one can predict the maximum volume of the drop that can be transferred between two substrates:

$$V_{predicted} = \frac{\beta_1 W_1 \gamma_{LV} (1 + \cos \theta_{RI}) - \beta_2 W_2 \gamma_{LV} (1 + \cos \theta_{RI})}{\rho g} \quad (5.14)$$

The model prediction was tested by determining the *maximum* drop volume that could be transferred (lifted) for a wide variety of dot-dot size combinations. Figure 5.4a shows a plot of the experimental versus predicted pickup volumes (from Equation 5.14) for the various dot-dot configurations used in this study. A linear least-squares regression (data not shown) was performed to find the proportionality constants β_1 and β_2 (in Equation 5.14), and the values were found to be $\beta_1 = \beta_2 = 1.35$.

A closer look at the force equations for extensional-adhesion (Equation 5.12) and drag-adhesion (Equation 5.2) reveals that the force needed to overcome extensional-adhesion will always be greater than for drag-adhesion for a specific liquid-substrate combination. This can be explained by comparing Equations 5.2 and 5.12:

$$F_{\text{extensional-adhesion}} \sim (1 + \cos \theta_R) > F_{\text{drag-adhesion}} \sim (\cos \theta_R - \cos \theta_A) \quad (5.15)$$

To achieve equality between the adhesive forces, the advancing CA (θ_A) would have to be 180° which is not possible with our superhydrophobic substrates. Hence, the extensional adhesion is always greater than the drag adhesion for our substrates. When the substrates are not parallel to each other during the drop transfer, a combination of drag and extensional adhesion may be experienced and the force balance becomes more complex. The drop transfer experiments were performed manually without using tools to optimize alignment. In spite of this, the good correlation between experimental and predicted values supports our simple model hypothesis.

The model was further extended to the more complex line configurations by determining the transfer of drops from a dot to a line. This configuration (dot to line) was selected because of its applicability to the proposed LOP device that is discussed below. In this scenario, the width of the lines and dots were used as the characteristic length scales W_1 and W_2 , respectively. The β_2 value of 1.35 was used, as determined from dot-dot transfer experiments. The β_1 value of 1.75 was obtained by linear least-squares regression as discussed previously. Figure 5.4b shows the experimental versus the predicted pickup volume for the various dot-line configurations used in this study.

In these studies, we have taken advantage of differences in the surface energy between ink and superhydrophobic paper to control the adhesion forces exerted on the drop. Two modes of adhesion exist for the drops, which were designated as ‘extensional-adhesion’ force and ‘drag-adhesion’ force, respectively, and these two adhesion forces can be tuned by varying the length scales of the ink patterns.

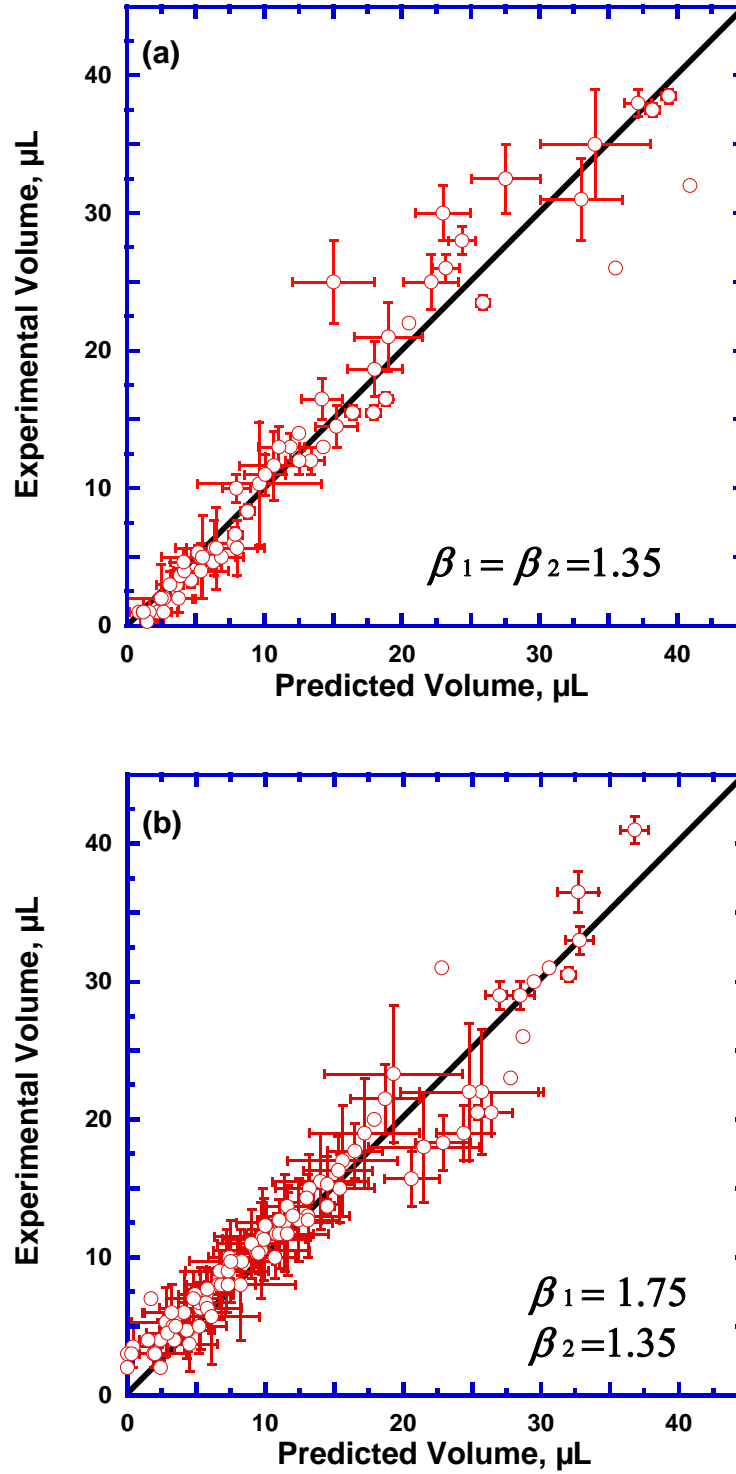


Figure 5.4 Experimental versus maximum drop pick-up volume for transfer from dot to dot (a) and dot to line (b); β_1 and β_2 are the fit parameters in Eqn. 14

5.3.2 Functional unit operations with patterned substrates

In the previous section, it was shown that patterned superhydrophobic paper substrates can be used to control the mobility of liquid drops on these substrates, both parallel and perpendicular to the substrate. The underlying mechanism is that the ink patterns locally increase the contact angle hysteresis on a low hysteresis superhydrophobic surface.

Therefore these patterned substrates are referred to as Hysteresis Enabled Lab-on-Paper (HELP) substrates. In this section, we will discuss how the fundamental wetting properties of these substrates can be used to engineer unit operations that can then be combined in lab-on-paper devices. The following paragraphs discuss the implementation of four basic functionalities that are critical for any device based on droplet manipulation: storage, transfer, mixing and sampling. The complete list of possible functionalities is certainly longer, especially when more complex ink patterns and modifications in ink chemistry are considered.

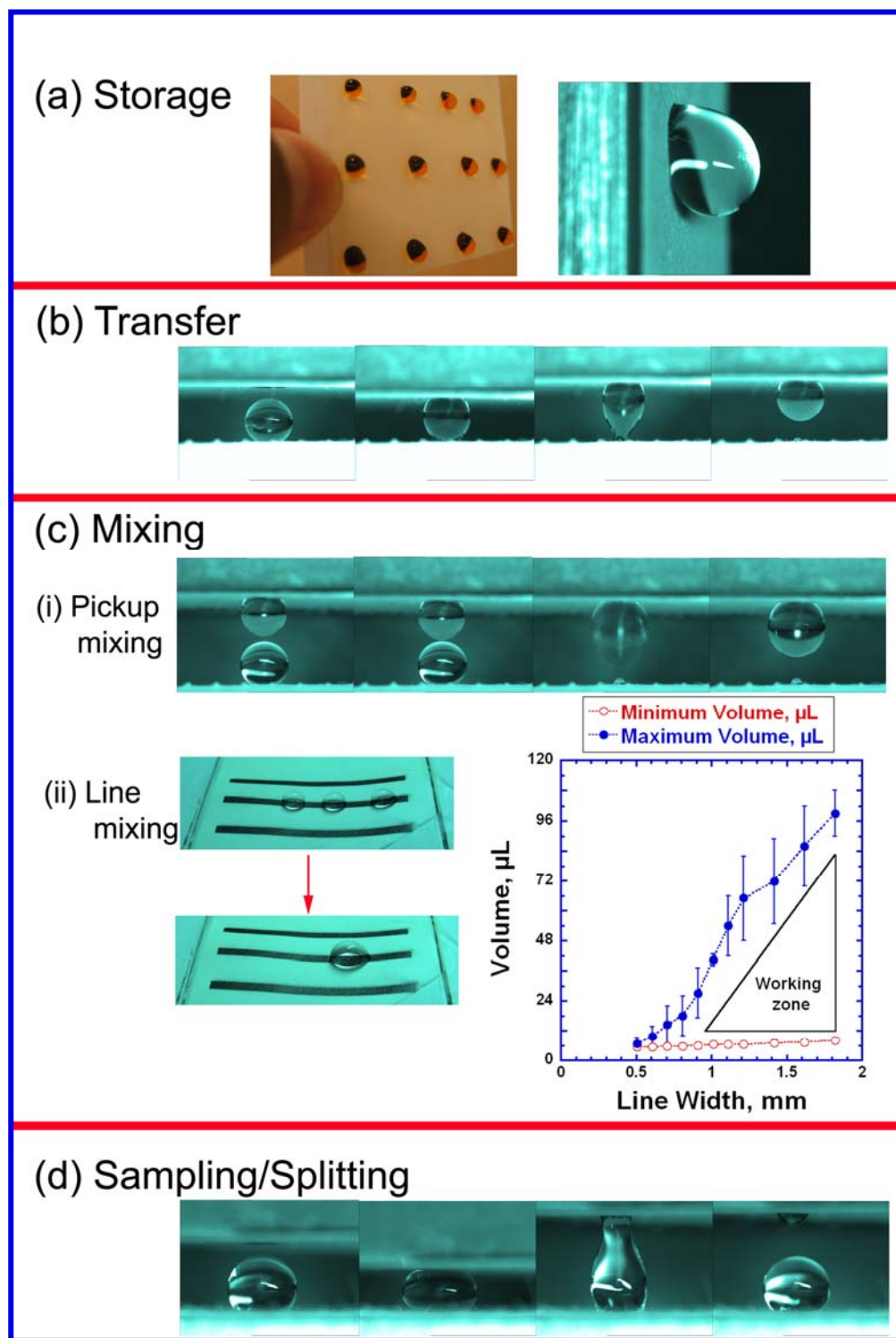


Figure 5.5 a) Photographs of an array of drops (food coloring was added to enhance contrast) and a high magnification image of a single drop stored on a vertical substrate, b) series of snapshots of a drop being transferred between two substrates, c) photographs of merging and mixing: i) via “pickup mixing” (two drops), ii) “line mixing” (three drops) and plot that shows the working zone of drop volumes suitable for line mixing, d) photographs of drop splitting between two substrates

5.3.2.1 Storage

The patterned paper substrates have a peculiar combination of two extreme wetting properties: minimal contact area between liquid and substrate due to high advancing contact angles and good adhesion due to hysteresis. It is believed that these properties make HELP substrates potential candidates to serve as storage media for arrays of microliter drops of test fluids and reagents. This type of storage is generally achieved with more expensive well plates that confine liquids in 3D wells with larger interfacial contact areas. It is believed that our patterned substrates can provide an inexpensive replacement for current technologies for the storage of array of drops for high throughput screening. The photograph on the left in Figure 5.5a shows an array of water drops (colored using food coloring) stored on a vertical substrate. The three rows had dot sizes of 1.84, 2.37 and 3.32 mm and drops volumes of 15, 20 and 25 μL , respectively. The photograph on the right in Figure 5.5a shows a high magnification image of a 12 μL water drop stored on a vertically placed substrate with a 1.7 mm dot. In spite of the low interfacial contact area, the drop withstands a tilt of 90 degrees.

5.3.2.2 Transfer

In Section 1.2, the transfer of microliter drops between two patterned substrates with different pattern dimensions was investigated. It was also demonstrated that the maximum drop volume that can be transferred between substrates can be predicted using the modified Young- Dupré equation. Figure 5.5b shows a sequence of frames from a movie that captured the transfer of a 4 μL water drop from a 0.4 mm dot to a 1.45 mm dot. This functionality enables selective transfer of drops from an array by carefully

tuning the size and location of dots on a pick-up substrate; the superhydrophobicity of the base paper will guarantee that no transfer will occur in non-patterned areas of the pick-up substrate.

5.3.2.3 Mixing

The patterned paper substrates can also be used to merge and mix liquid drops. Two strategies for drop mixing were explored, which are referred to as pickup mixing and line mixing. Figure 5.5c (i) shows how two 4 μL water drops (attached to the “from” (0.4 mm) and “to” (1.45 mm) substrates, respectively) are merged into a single drop using pickup mixing. After the two drops are roughly aligned and the substrates are brought together, the drops touch and merge. The final position of the merged drop depends upon the competing adhesive forces of the upper “from” and lower “to” dots. As demonstrated in Section 1.2, the size of “from” and “to” substrates can be tailored to enable pickup mixing for a variety of drop volumes.

The second mixing strategy, line mixing (Figure 5.5c (ii)), enables mixing of two or more drops on a line by taking advantage of the fact that the mobility of the drop on the line depends on the drop’s configuration on the line. There are three basic configurations that are important for line mixing. 1) If the drop is positioned on the line without touching the end points of the line, the force needed to induce sliding along the line is given by Equation 5.7 or 5.8, as discussed previously and shown in Figure 5.3b. 2) When a drop slides and reaches the end of the line, its advancing edge contacts the SH paper and hence θ_{AI} in Equation 5.7 and 5.8 must be replaced by θ_{ASH} , which results in a significantly

increased value of F_{PO} and F_{PP} . Therefore, the sliding angle for a drop at the end of a line is always greater than at other positions. 3) The mobility of a drop perpendicular to the line is much lower than along the line for the same reasons, which restricts drop movement to the line. These considerations concerning drop mobility on a line can be used to design a mixing strategy for drops positioned on the line, by simply tilting the line back and forth drops can be moved towards the line edge, where they become pinned, so that trailing drops can merge. Subsequent rocking of the substrate then moves the merged drop back and forth along the line, which induces internal mixing of the drop. Figure 5.5c (ii) shows pictures of the merging and mixing of three 20 μL drops into a single 60 μL drop via this line mixing strategy.

There are some limitations to this type of mixing. When drops start sliding, they initially accelerate along the line; if their momentum becomes too large, the adhesive force at the end of the line may be insufficient to ensure adherence to the line edge. Similarly, vibrations can provide energy for the drop to break away from the line. Both effects will limit the maximum drop size for line mixing, but it is beyond the scope of this work to model these phenomena in detail. Instead, the overall effect was explored by evaluating the reproducibility of drop size limits for line mixing. Three individuals with different levels of experience and skill sets performed line mixing experiments and independently determined the minimum and maximum drop volumes that could be mixed for different line widths. The plot in Figure 5.5c (ii) shows the results of these tests; the ‘working zone’ drawn between the two curves denotes the drop volume range for mixing as a function of line width. It is important to point out that these experiments were performed

with lines that were 3cm long. According to our experience with line mixing, longer lines expand the ‘working zone’ because of improved operator control over the drop position.

5.3.2.4 Sampling/ splitting

Similar to the transfer of a drop between two substrates (shown in Figure 5.5b), the patterned substrates can also be used to sample small volumes of liquid from a single drop. Figure 5.5d shows the sampling of a small volume of liquid to a 1.45 mm dot from an 8 μ L drop resting on another 1.45 mm dot. By using closely matched “from” and “to” dot sizes, patterned SH substrates can be used to collect small sample volumes of liquid from a single drop. This functionality is useful in LOP applications in which it is desirable to obtain multiple samples from individual drops for multiplex analysis.

5.3.3 Integrated lab-on-paper (LOP) concepts based on HELP substrates

By using the functionalities described in the previous section, it is possible to confine microliter drops to specific locations on a storage substrate, selectively transfer (pick up) drops between substrates, combine/mix drops and sample/split the products into multiple drops. These unit operations can be combined to create a simple lab-on-paper (LOP) device. Figure 5.6 shows a schematic of such a LOP device that can be fabricated using the HELP substrates. As starting point, an array of drops is positioned on a substrate using a pattern of dots. In step 1, selected drops are transferred to another substrate which has larger dots printed at specific locations. In step 2, these selected drops are picked up by a third substrate with a line pattern. Finally, in step 3, the three drops on the line are merged into a single drop and mixed via ‘line mixing’. Thus, selected reactants can be

picked up from an array of reactants and mixed to obtain the final product. We should emphasize that this specific LOP configuration is just one example of the possible configurations that can be fabricated. The versatility of the printing technique provides the opportunity to design and create new LOP configurations at the end-user level.

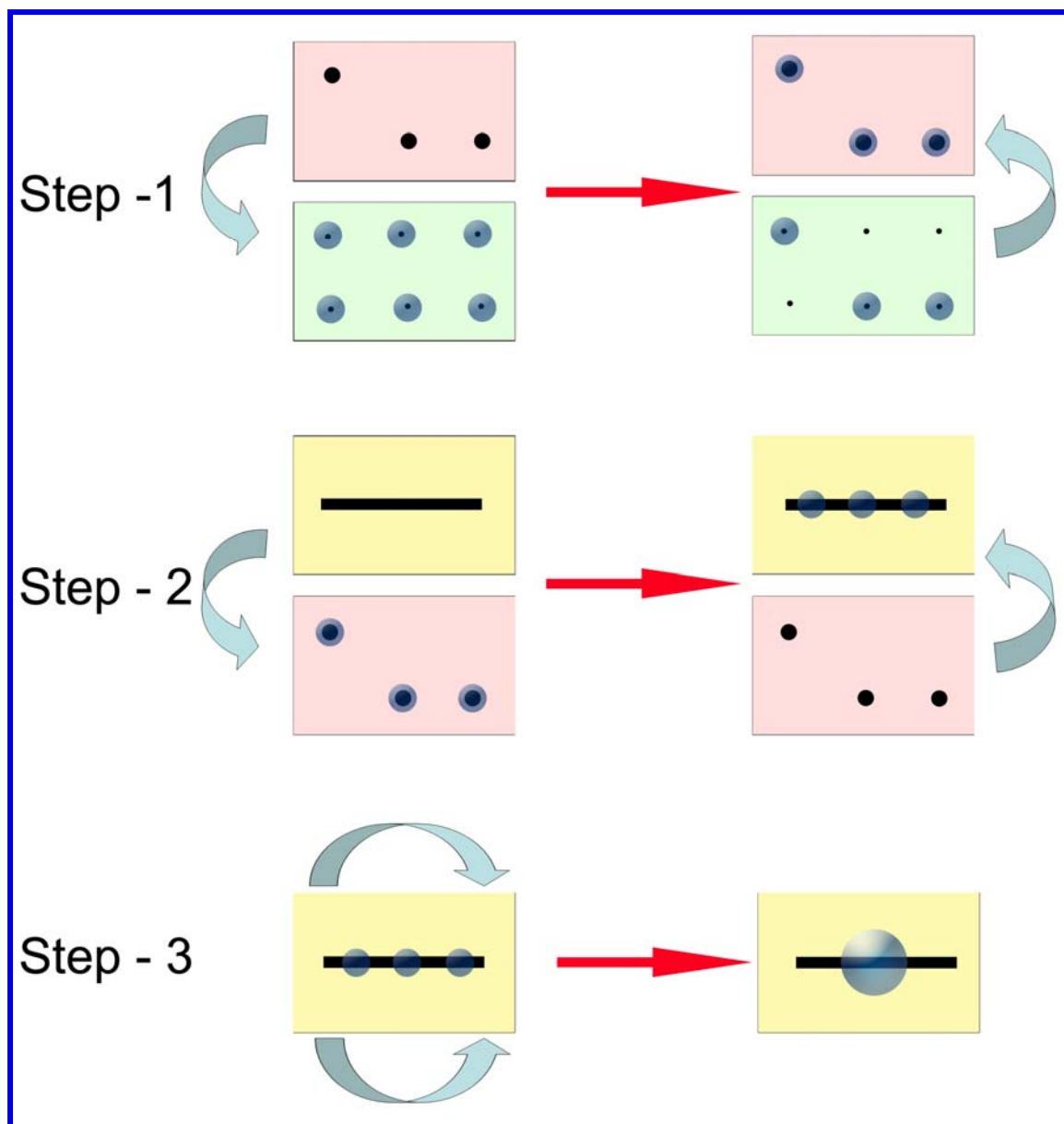


Figure 5.6 Schematic of a simple LOP that can be fabricated using the HELP substrates

5.4 Conclusions

A commercially available phaser printing technology was used to pattern superhydrophobic paper substrates with high-hysteresis ink patterns (dots and lines). By tuning the shape and size of the ink patterns, the drag-adhesion and extensional adhesion of liquid drops to the substrate can be controlled. Experimental results for the adhesive forces of water drops on these patterned substrates are in good agreement with classical models for drag-adhesion (Furmidge equation) and extensional-adhesion (modified Dupré equation) over a wide range of pattern sizes and drop volumes. Fundamental knowledge of the dependence of adhesive forces on pattern parameters and the resulting control over drop mobility were then used to design substrates for four basic functionalities that are relevant for lab-on-paper (LOP) applications: drop storage, drop transfer, drop mixing and sampling.

Previous LOP devices depended on the capillary forces inside the paper to enable the transfer and mixing of test fluids. Hence, this approach rules out the possibility to extract multiple samples of the reaction products for various analyses. Using HELP substrates, the required unit operations were obtained by merely manipulating the liquid drops on top of the substrate by tuning adhesive forces. Another unique advantage of the HELP substrates over existing LOP devices is the ability to store liquid drops after initial, qualitative on-chip analysis for further testing with specialized methods at centralized testing centers, which would be relevant for bioanalytical applications in resource-limited settings [1]. The substrates can also be used as a disposable storage medium in

conjunction with high-throughput screening and potentially replace well plate technology.

Finally, the simplicity of the patterning techniques using commercially available desk top printing technology and standard word processing software provides extreme flexibility in substrate design. End-users can easily program their own substrates according to specific needs, using the design rules presented in this paper, which represents the first step towards creating successful LOP devices. Further success of this technology can be achieved by developing custom printing ink formulations to enhance compatibility with various test fluids and facilitate deposition of reagent species.

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CHAPTER 6

DIRECTIONAL MOBILITY AND ADHESION OF WATER DROPS ON PATTERNED SUPERHYDROPHOBIC SURFACES*

6.1 Introduction

The intriguing self-cleaning property of the lotus leaf was identified and studied extensively in the 1970s and has been termed “lotus effect”. On a lotus leaf, water drops easily slide off in all directions. But there are also several natural surfaces where directionality is observed in droplet roll-off behavior. For instance, rice leaves exhibit anisotropic dewetting: the adhesion to drop sliding depends on whether the drop moves perpendicular or parallel to the orientation of micropapillae [1]. Similar anisotropic drop sliding behavior has been observed in bamboo, screw pine and some grasses; for these plants, water can roll off only along the length of the leaf [2]. On duck and goose feathers, water drop roll-off can only occur along the vein of the feather outward [3]. Drop mobility on the wings of butterflies is even more intriguing; in addition to uni-directional drop sliding, the wing surface can also reverse the direction of sliding based on its posture (upward or downward) [4]. Finally, on the petals of the red rose (*rosea Rehd*), water drops with volumes less than 10 μL do not roll-off in any direction despite near-spherical droplet shapes with a high contact angle ($\text{CA} \sim 152.4^\circ$) [5].

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Many properties of natural materials have served as motivating factors for research, which has become evident through the emergence of the field of biomimetics. Advances in nanofabrication technology over the past few decades have enabled the fabrication of artificial surfaces that mimic the behaviors listed above on a variety of organic and inorganic substrates. For instance, the ‘lotus effect’ (advancing CA>150° and CA hysteresis <10°), which is also classified as ‘roll-off’ superhydrophobicity, has been obtained on a variety of surfaces [6-11]. The fabrication of ‘sticky’ superhydrophobic surfaces (advancing CA>150° and CA hysteresis > 10°) that mimic the behavior of a ‘rose petal’ has also been reported [5, 12-14]. Chapter 2 described the fabrication techniques for obtaining roll-off and sticky superhydrophobic behaviors on paper surfaces via plasma processing. Chapter 3 and 4 described the ability to tune adhesion forces on superhydrophobic paper surfaces between two extremes (roll-off and sticky) by controlling the topography of the cellulose fibers by means of selective plasma-enhanced etching, fiber types and paper making techniques.

Recently, there has been increased interest in fabricating superhydrophobic surfaces with directional preference in the drop sliding behavior [15]. Although there have been numerous studies that report anisotropic drop sliding [16-18] and wetting [2, 19-24] on artificial hydrophobic surfaces, directional drop sliding behavior on superhydrophobic surfaces has not been extensively studied. Superhydrophobic surfaces with parallel, grooved microtextures were fabricated to create anisotropic drop sliding behavior: aligned carbon nanotubes (ACNT) were grown with ordered pillar patterns similar to the topology of a rice leaf and hydrophobized; these structures mimic the anisotropic sliding

behavior observed on rice leaves [1]. Other studies created an S-shaped surface pattern on a roll-off superhydrophobic glass surface using a cutting blade [25]. With the proper tilt angle, a water drop (4 μL) was able to slide along the S-shaped curve. In both these cases, directional adhesion was obtained by changing surface topography. For anisotropic sliding behavior on ACNT films, the difference between orthogonal critical sliding angles was small due to the weak adhesion forces [1]. Similarly, for patterned glass slides, droplet manipulation was not successful for volumes larger than 20 μL and for tilt angles larger than 45° because of the limited adhesion force exerted by the difference in topography [25]. To the authors' knowledge, superhydrophobic surfaces that offer directional drop mobility with a large difference between sliding angles for a wide range of drop volumes, have not been reported

In Chapter 5, we discussed the ability to control drag and extensional adhesion of water drops on superhydrophobic paper surfaces by patterning the surfaces using a commercial desktop printer. In that work, the difference in surface energy between the superhydrophobic paper substrate and the less hydrophobic ink film was exploited to manipulate the adhesion and mobility of water drops. The objective of this chapter is to create surfaces with anisotropic drop mobility by using ink patterns on superhydrophobic paper. The guiding hypothesis is that directionality can be obtained by imposing geometrical constraints on the movement of the three-phase contact line between drop, paper and air through the shape of the ink patterns. In addition, we investigated the tunability of sliding angle by creating chemical heterogeneity in the form of 'checkered' ink patterns. Controlled anisotropic drop mobility and adhesion on superhydrophobic

surfaces has potential application in the formation of fundamental components for 2D microfluidic devices, such as flow paths, gates/diodes, junctions and drop size filters.

6.2 Experimental

6.2.1 Superhydrophobic paper

Paper handsheets were prepared with a combination of 50% hardwood and 50% softwood fibers according to TAPPI standardized method T205 sp-02. The handsheets were selectively plasma etched with oxygen gas (0.55 torr) and a thin (100 nm) fluorocarbon film was subsequently deposited onto the etched surface by plasma-assisted deposition (1 torr) using pentafluoroethane (PFE) as precursor. A more detailed discussion of the methods to fabricate superhydrophobic paper can be found in Section 2.2.1 and 2.2.2 (Chapter 2).

6.2.2 Patterning

All printed ink patterns other than the ‘checkered’ ones were designed with Adobe Photoshop (version 7.0; Adobe Systems Inc., San Jose, CA). The ‘checkered’ patterns were designed through MATLAB (version 7.6.0; The MathWorks Inc., Natick MA). For the ‘checkered’ patterns in Figure 6.5a, 6.5b, 6.5c and 6.5d a predefined square area was divided into square matrices with 750 X 750, 250 X 250, 1400 X 1400 and 1200 X 1200 cells, respectively. All pixels constituting the cells of the matrix were designated ‘white’ ([R B G] = [255 255 255]) to ensure a white background. For Figures 6.5a and 6.5b, starting from the left top cell (0, 0), the pixels in every second cell (in both x and y

directions) were designated ‘black’ ($[R\ B\ G] = [0\ 0\ 0]$). For Figures 6.5c and 6.5d, starting from the left top cell (0, 0), the pixels in every fourth cell (in both x and y directions) were designated ‘black’ ($[R\ B\ G] = [0\ 0\ 0]$).

The SH paper was pasted on a piece of copy paper using Scotch® Tape and fed through a Xerox Phaser 8500n (Xerox Corporation, Norwalk CT) printer to generate the patterns.

6.2.3 Image Analysis

‘Checkered’ patterns with variable areal ink fractions were imaged with a Leica microscope (DM4500 B) using a 2.5X objective as described in Section 4.2.5 (Chapter 4). The images were then processed using Image-J (Java-based image processing program developed at the National Institutes of Health). For all imaged ‘checkered’ substrates, a threshold was manually set for the pixel brightness to differentiate ‘ink’ pixels from the SH paper background. Then the ‘area’ and ‘area fraction’ of the ink were obtained using the “Analyze particles” feature in Image-J. The ‘area’ value obtained from the Image-J was divided by the number of features present on the image and the square root of this ratio was taken to calculate the size of the ink features. The distance between features was calculated by measuring the distance (x value, in pixels) at the center of two adjacent features and converting the difference to units of μm using the microscope conversion factor. This procedure was repeated for a variety of pixel pairs to obtain statistically averaged values.

6.2.4 Sliding drop measurements

The patterned SH paper substrates were mounted on a flat surface attached to a rotating optical stage. For the directional substrates, a drop with a specific volume was placed at the center of the pattern. For the ‘checkered’ pattern substrates, a drop with a specific volume was placed randomly. After placing the drop on the horizontal substrate, the plate was tilted gradually until roll-off was observed; the critical sliding angle was obtained from the angular scale of the optical stage.

6.3 Results and discussion

6.3.1 Control of directionality

6.3.1.1 Design for directional sliding

It has been well-established that drop sliding on a surface is primarily governed by movement of the three-phase liquid/solid/vapor contact line [16, 26]. This contact line can be divided into segments with advancing and receding contact lines, as observed at the advancing and receding edges of a sliding drop. The sliding of drops on inclined surfaces was extensively studied and modeled in the mid-twentieth century [27-30]. In 1962, Furmidge [28] surveyed the literature at that time and formulated a semi-empirical force balance (Equation 5.3). After rearranging this equation, we get:

$$\sin \alpha = \frac{\gamma_{LV} W_{drop} (\cos \theta_R - \cos \theta_A)}{\rho V g} \quad (6.1)$$

where α is the critical sliding angle, γ_{LV} is the surface tension of the liquid, W_{drop} is the width of the drop perpendicular to the direction of sliding, θ_A and θ_R are the advancing and receding contact angles of the drop on the surface, ρ is the density of the liquid drop, V is the volume of the drop and g is the acceleration due to gravity. From Equation (6.1) it is clear that the contact angles established at the advancing (θ_A) and receding (θ_R) contact lines are critical in determining the drop sliding behavior.

Consider a semi-circular ink pattern with diameter D printed on superhydrophobic (SH) paper, as depicted in Figure 6.1. The surface energy of the ink film is higher than that of the SH paper, leading to advancing and receding CAs on the ink film that are always lower than those on the SH paper. If a water drop with a foot print diameter equal to the diameter of the semi-circle is placed on the paper substrate, the semi-circle will be aligned with the contact line of the drop. When the substrate is tilted towards the concave or convex direction of the pattern, the advancing contact line of the drop always experiences the advancing CA on the SH paper. However, the receding contact line of the drop experiences different interactions with the ink pattern depending upon the sliding direction. For sliding in the concave direction, the ink curvature embraces the receding contact line of the drop; as a result, the receding contact line experiences the receding CA on the ink until drop detachment from the pattern occurs. During sliding in the convex direction, the receding contact line is more complex: the receding contact line and the semi-circular ink pattern appear as two intersecting semi-circles. When the drop finally detaches from the pattern, the contact line and the pattern touch each other at a single point, similar to contact between two hard-spheres. In other words, when a drop slides in

the convex direction, a major fraction of the receding contact line is always in contact with the SH paper, while only a small fraction touches the ink film. It is therefore reasonable to assume that the effective receding CA is close to the receding CA of the SH paper. So, while the advancing CA for both sliding directions is the same, the receding CA for the concave direction (similar to ink film) is always higher than the receding CA for the convex direction (similar to SH paper).

The qualitative analysis described above suggests the existence of different CAs for sliding in the concave and convex directions and provides new insight in relation to the Furmidge equation (Equation 6.1). For a particular drop volume-pattern diameter combination, all terms on the right hand side of Equation (6.1) are the same for the two sliding directions except for the receding CA. Because the receding CA for the concave direction is greater than that in the convex direction, as discussed above, it is expected that the critical sliding angle (α) will also be higher for the concave direction. Figure 6.1 shows a plot of the critical sliding angles for the concave and convex directions for various drop volumes on a semi-circular pattern with a diameter of 2.55 mm. It is evident that the critical sliding angle in the concave direction is consistently higher for all drop volumes investigated.

To further investigate this directional sliding behavior, the critical sliding angles for a variety of semi-circle diameters was studied. Figure 6.2a shows a plot of the critical sliding angle vs. semi-circle diameter for the concave direction. For all drop volumes, the critical sliding angle increases with semi-circle diameter. For combinations of large semi-

circles with small drops (diameter 3.33 mm for $V < 30 \mu\text{L}$ and 4.14 mm for $V < 40 \mu\text{L}$), drops slide along the curved semi-circle towards either end of the curve rather than breaking away from the pattern near the cusp of the semi-circle. Data points for such asymmetric roll-off events are not plotted in Figure 6.2a. Figure 6.2b shows a plot of the critical sliding angle vs. the semi-circle diameter for sliding in the convex direction. Unlike the concave direction (Figure 6.2a), the convex sliding direction (Figure 6.2b) displays no clear trend in critical sliding angle with respect to pattern size.

As explained previously, when the droplet slides in the concave direction sliding (Figure 6.2a), the curvature of the ink pattern matches the receding contact line of the drop. As a result, the ink pattern size has a direct impact on the fraction of receding contact line touching the ink film: the larger the pattern size, the larger the fraction of receding contact line that touches the ink film. A larger ink-contact line interface enhances the adhesion of the drop to the ink pattern and hence increases the critical sliding angle. Therefore, the critical sliding angle increases with semi-circle diameter (Figure 6.2a). On the other hand, in convex direction sliding (Figure 6.2b), the receding contact line primarily interacts with the SH substrate, so that the ink pattern size has less effect on drop adhesion and critical sliding angle. In this case, the critical sliding angle is essentially independent of semi-circle diameter (Figure 6.2b).

Mathematical modeling of the critical sliding angle on semi-circle ink patterns is complex because the CA varies continuously along the contact line. Despite the lack of a quantitative model, the qualitative comparisons described above provide a reasonable

understanding of water drop sliding behavior on the ink patterns. Clearly, critical sliding angles can be altered by establishing positive or negative curvature of ink patterns to achieve directional mobility of water drops on superhydrophobic paper surfaces.

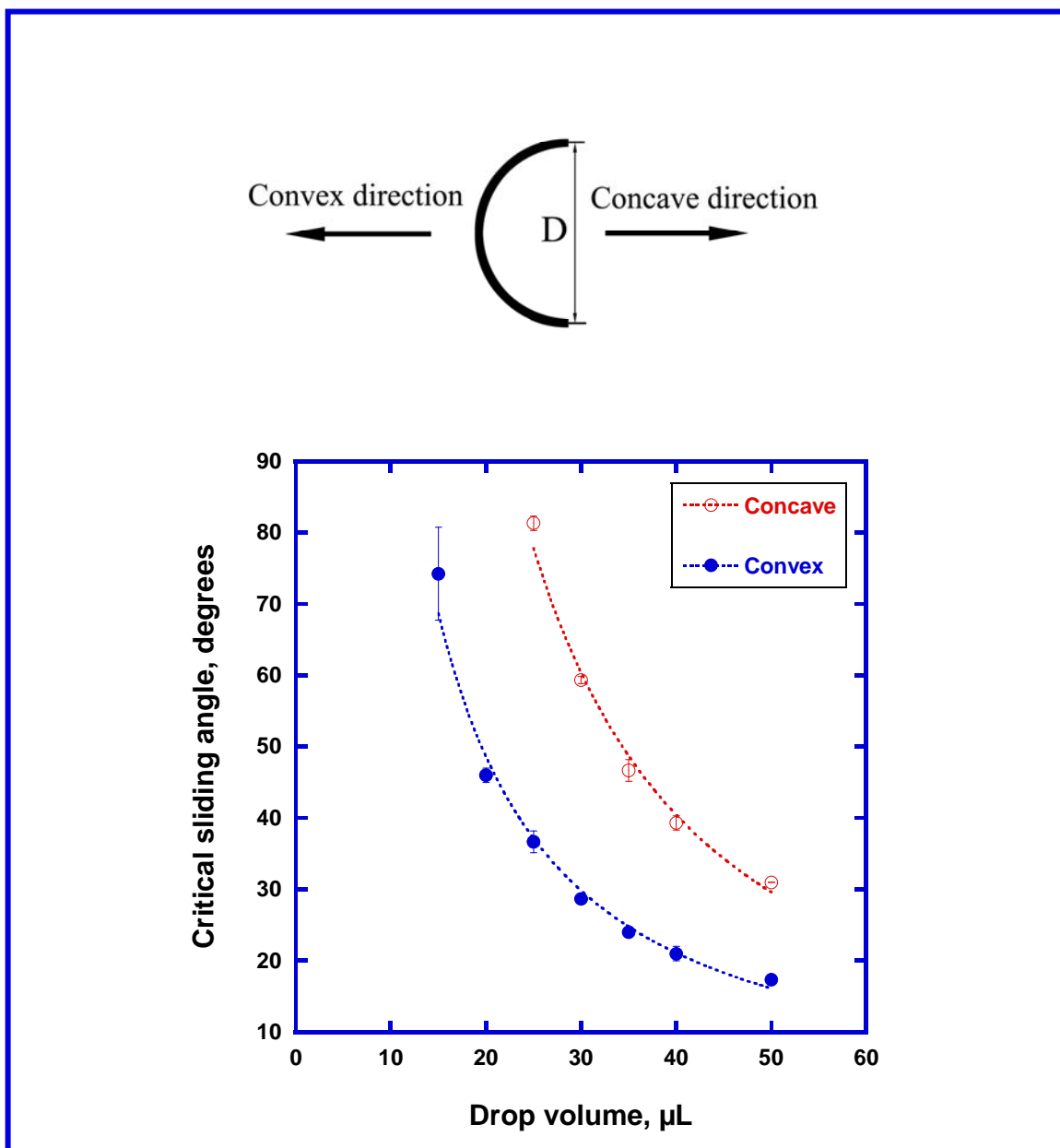


Figure 6.1 Critical sliding angles in different sliding directions vs. drop volume for a semi-circle ink pattern printed on superhydrophobic paper substrate ($D = 2.55\text{mm}$); the slide directions are defined at the top of the figure

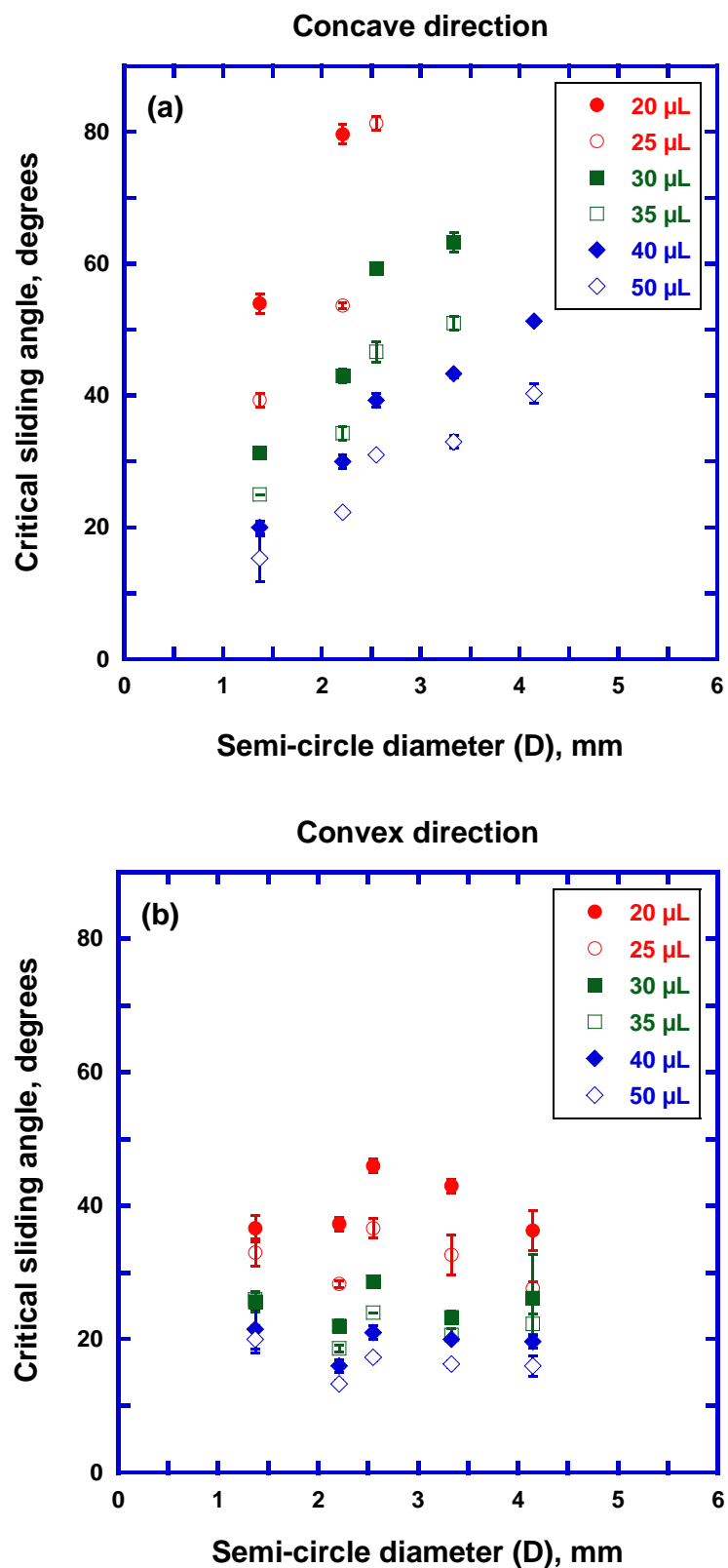


Figure 6.2 Critical sliding angle vs. diameter of semi-circle ink patterns printed on superhydrophobic paper substrates for various drop volumes in concave (a) and convex (b) sliding directions (see Figure 6.1 for definition of sliding direction)

6.3.1.2. Different types of directional sliding

Directional sliding of water drops on superhydrophobic surfaces may offer novel approaches for the design of 2D microfluidic devices. For example, anisotropy can be exploited to create 2D flow paths, valves/diodes, and junctions for drops on these surfaces. In the following, several examples of patterns that apply the concept described in the previous section to achieve specific droplet manipulation are discussed in detail.

Figure 6.3a shows the simplest anisotropic design, a straight line. When a drop is placed on this line pattern, the width of the drop (W_{drop}) in the ‘parallel’ sliding direction is always smaller than in the ‘perpendicular’ sliding direction. This is because of the elongation of the drop along the line as a result of the higher surface energy of the ink line. The receding CAs for both sliding directions are similar to that of the receding CA on the ink. While the advancing CA in the ‘perpendicular’ direction is similar to that of the advancing CA on SH paper, the advancing CA in the ‘parallel’ direction is similar to that of the advancing CA on ink. Such behavior is expected, because as the drop slides along the line (parallel direction), the part of the drop that resides on the line always remains in contact with the ink film. Thus ‘perpendicular’ direction sliding always has a larger width (W_{drop}) and advancing CA (θ_A) relative to the ‘parallel’ direction. Combining this qualitative analysis with the Furmidge equation (Equation 6.1), it is evident that drop sliding parallel to the line is easier than sliding perpendicular to the line. Figure 6.4a clearly shows the anisotropy of drop sliding in directions parallel and perpendicular to a line. A water drop placed on a line moves easily along the line, but it needs a higher tilt angle to detach from the line in the perpendicular direction. This two-directional sliding

functionality of the line pattern can be used to design 2D flow paths for 2D microfluidic devices. For example, lines can be used to ensure directional mobility of drops from one position to another by merely tilting the substrate.

While a simple straight line pattern exhibits two-directional sliding, a semi-circle can be inserted between two lines to establish anisotropy for sliding in one-direction. Figure 6.3b shows such a design which is composed of a semi-circle and two line elements. As described in the previous section, a semi-circle exhibits preferential drop sliding toward the convex side of the semi-circle; lines are included in the design to create a defined flow path for drops moving toward and away from the semi-circle element. The intent in this design is to create a droplet gate or diode, which allows transport of the drop along the line only in one direction. Figure 6.4b presents critical sliding angles in the concave and convex sliding directions for a variety of drop volumes. Comparison of these data with the results for an isolated semi-circle in Figure 6.1 shows that the presence of lines near the semi-circle does not affect the directionality of drop sliding on the semi-circle. At the appropriate tilt angle, while a drop sliding along the line in the concave direction will have its movement inhibited by the semi-circle pattern; a drop sliding in the convex direction will pass the semi-circle unhindered. Of course, different drop volumes result in different sliding angle values for both directions during passage through the semi-circle gate. The ability to exhibit different tilt pass angles for different volume drops sliding in different directions on this pattern allows the fabrication of 2D gates/diodes and valves for 2D microfluidic devices.

Figure 6.3c is a schematic of a pattern created from three circular sections that touch each other. As discussed previously, for a single semi-circle, drop sliding in the convex direction is easier than in the concave direction. The same basic anisotropy applies to each of the three circle sections in Figure 6.3c, as is evident from the results presented in Figure 6.4c. This specific design can be used to design 2D junctions for flow paths on 2D microfluidic devices. For example, two flow channels can be merged into a single channel using this tri-directional junction pattern.

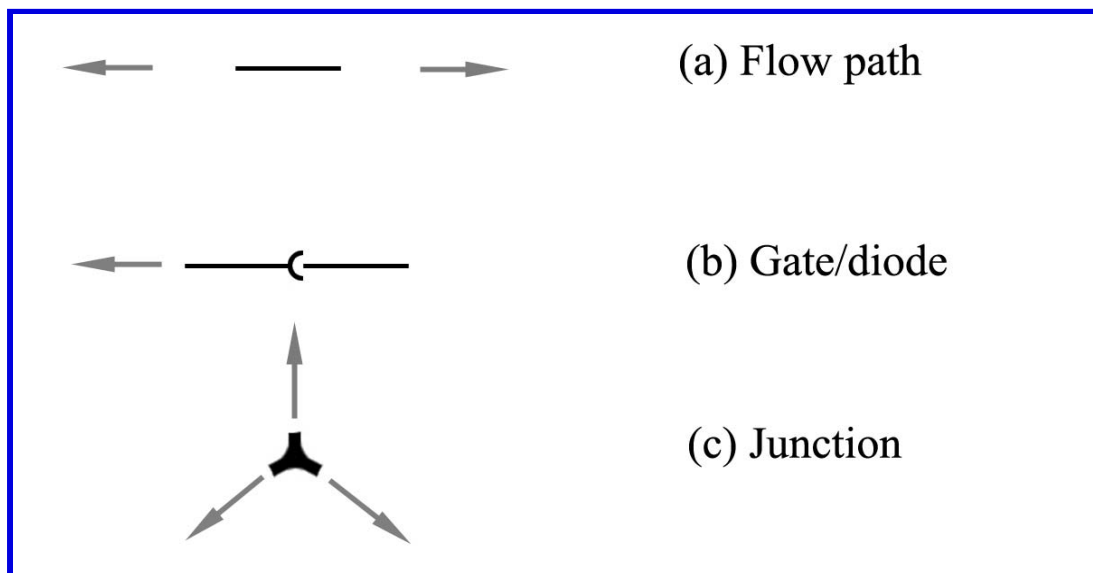
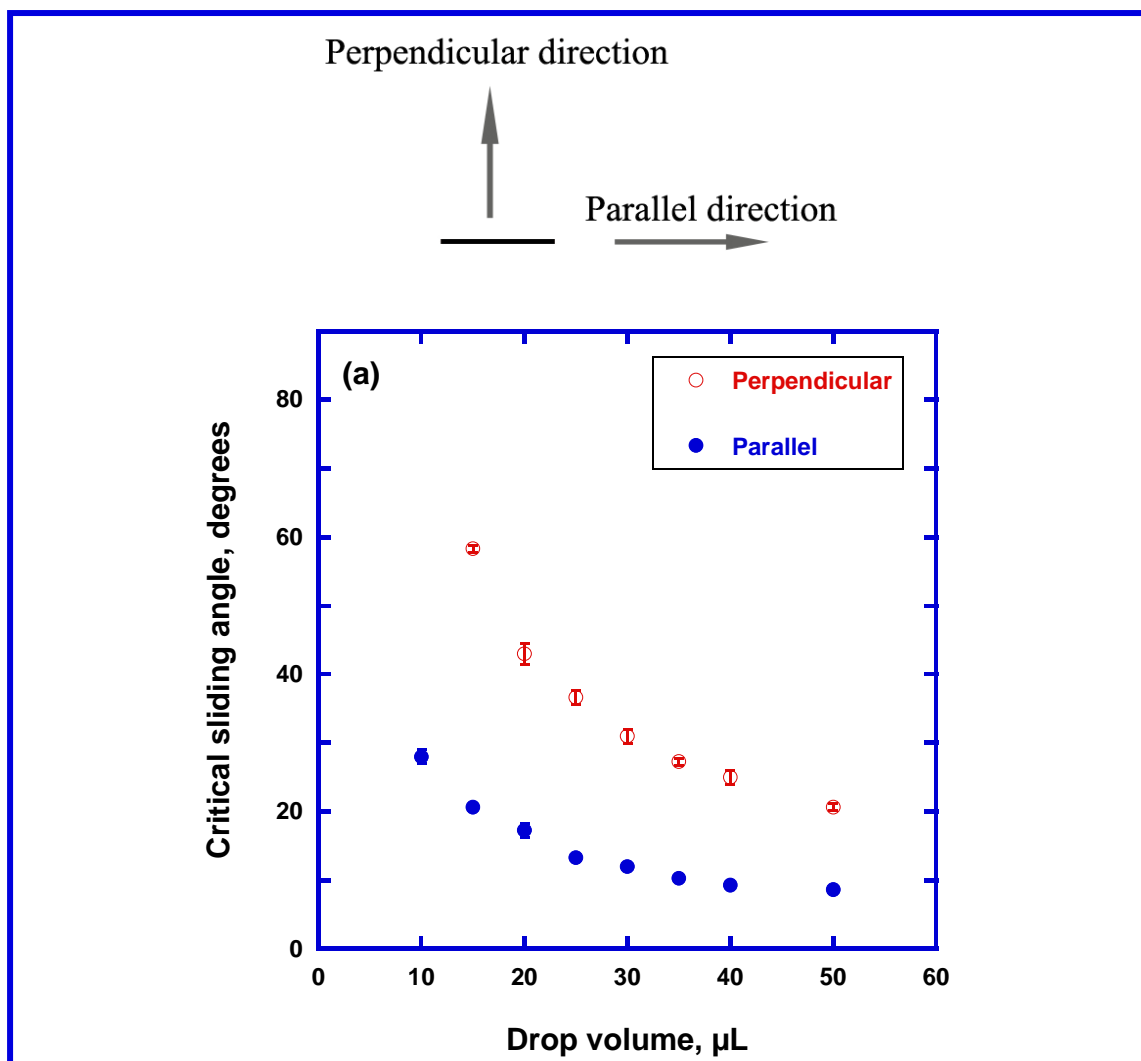
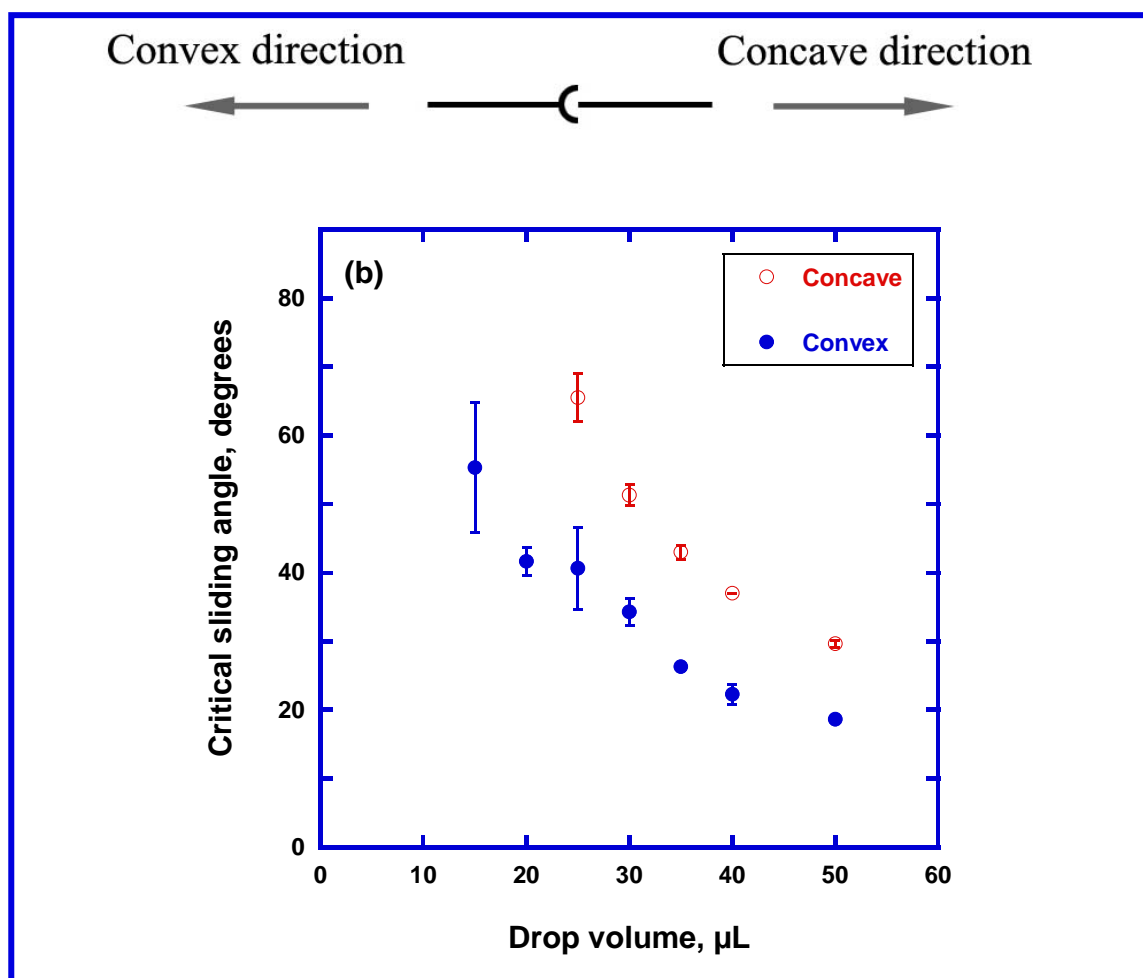


Figure 6.3 Schematics of possible components designed for a 2D microfluidic system generated by SH paper ink patterns





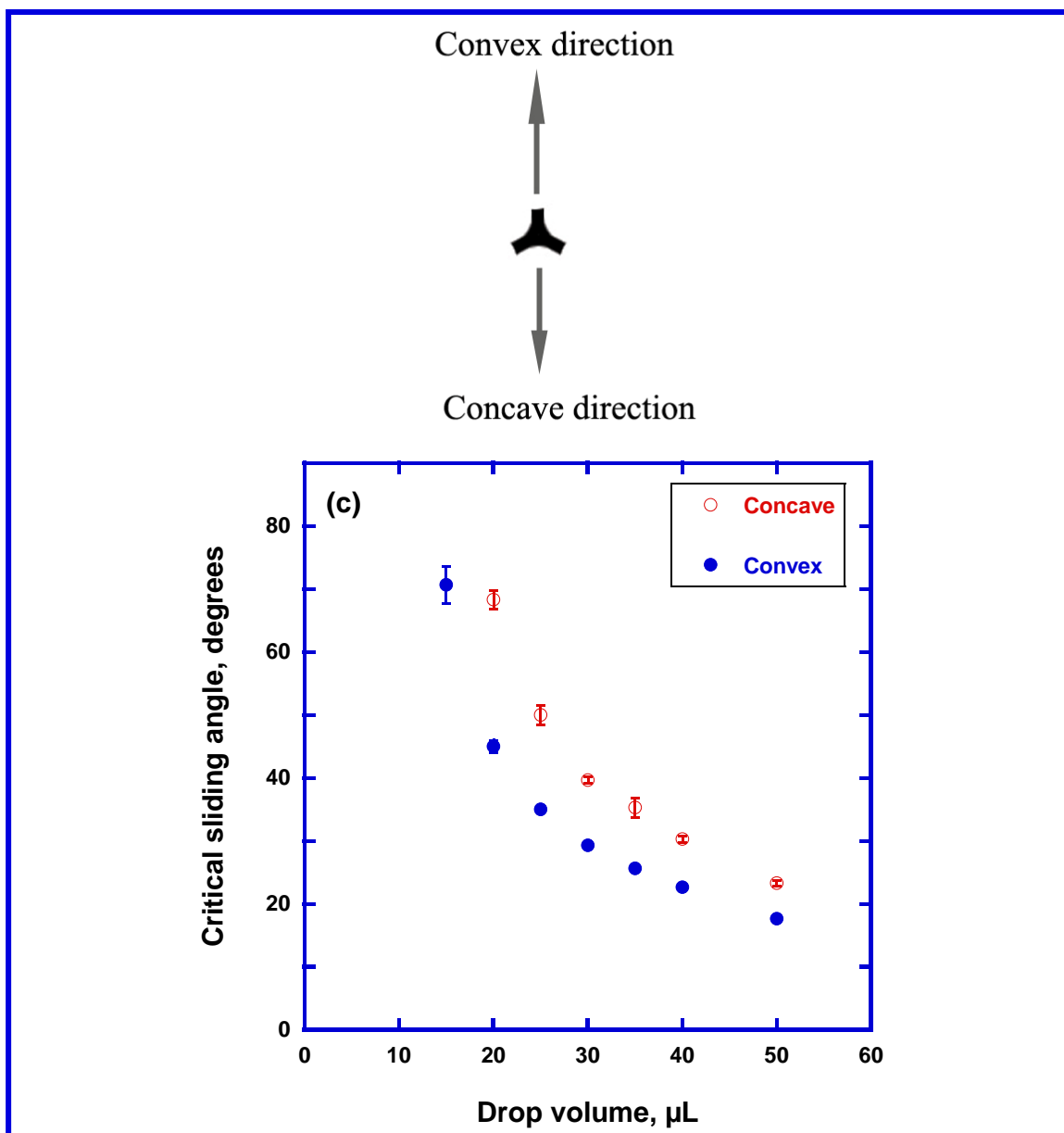


Figure 6.4 Schematics of ink patterns with specific sliding directions (top); plot of critical sliding angle vs. drop volume for flow path (a), gate/diode (b) and junction (c) patterns

6.3.2 Control of adhesion

Interaction of water droplets with ink or superhydrophobic surfaces offers the ability to tune the adhesion of water drops on superhydrophobic paper surfaces using ink patterns.

In 1948, Cassie proposed the following model to predict the apparent contact angles observed on composite surfaces:

$$\cos \theta^* = \sigma_1 \cos \theta_1 + \sigma_2 \cos \theta_2 \quad (6.2)$$

where θ^* is the apparent contact angle, σ_1 and σ_2 are the area fractions of the surface with intrinsic contact angle values θ_1 and θ_2 , respectively. This equation was originally formulated for a static drop. Later, Johnson and Dettre [31] proposed that for a dynamic drop, the contact angle hysteresis (difference between the advancing and receding CAs) is more important than the static CA, and hence the above equation is not valid to explain the scenario of a dynamic drop on a composite surface.

It was also proposed that the contact angle hysteresis can, in principle, be altered by changing the area fractions σ_1 and σ_2 . In addition to the area fractions, the contact angle hysteresis also depends on the vibrational energy of the drop and the height of the energy barriers between metastable states that the drop encounters immediately before roll-off occurs[31]. Johnson and Dettre [31] mathematically showed that a variety of hysteresis values can be observed for the same area fractions provided that the vibrational energy of the drop and the heights of energy barriers between metastable states differ. The two mathematical parameters invoked (defined as vibrational energy and the heights of

energy barriers) were subsequently studied in more detail and related to the topography of the three-phase contact line [9, 32, 33]. The work of Chen *et al.* demonstrated that the topography of the contact line can be altered by changing the sizes/shapes of the different surface elements on the composite surface [9]. Finally, altering the CA hysteresis or, in other words, the advancing and receding CAs, alters the critical sliding angle of water drops on the composite surfaces according to Equation (6.1). Therefore, previous studies in this field suggest that the hysteresis, or the critical sliding angle, can be altered by changing the area fraction and sizes/shapes of the different materials that make up the composite surface.

The fabrication of superhydrophobic surfaces with tunable sliding angles for a wide range of drop volumes is described below. Composite surfaces were prepared by printing ink patterns consisting of square ink dots as described in the experimental section (Figure 6.5). The patterns in Figure 6.5 were designed in such a way that the area fraction and feature size were varied independently. Table 6.1 shows the area fraction of ink, feature size of ink spot and distance between the ink spots for the four patterns shown in Figure 6.5. From Table 6.1, Figure 6.5a and 6.5b have the same area fraction (~25%) of ink printed on the superhydrophobic paper. However, the ink spot size in Figure 6.5b is larger than that in Figure 6.5a. Similarly, Figures 6.5c and 6.5d have the same area fraction (~8%) of ink but have different feature sizes (Table 6.1).

Figure 6.6 shows the sliding angle of water drops as a function of drop volume for the substrates shown in Figures 6.5 a, b, c, and d, along with sliding angles for a continuous ink film, for SH paper after passage through the printer but without printing (SH-control) and for SH paper that was not fed through the printer (SH-virgin). From these results, it is evident that composite surfaces with an ink fraction of $\sim 25\%$ (6a and 6b) display results close to those for the ink film, while the composite surfaces with an ink fraction of $\sim 8\%$ (6c and 6d) display properties more similar to the SH control substrate. These results indicate that by changing the area fraction of the composite surface, critical sliding angle values can be tuned between the values obtained for SH paper (control) and an ink film. Further consideration of the curves in Figure 6.6 with the same area fraction but different feature sizes demonstrates that the sliding angles can be altered by changing the feature sizes. When feature size increases, the contact line becomes smoother, which minimizes the tension of the drop on the contact line and thus lowers the energy barrier height for drop sliding. Indeed, Figure 6.6 shows that patterns with larger features (6.6b and 6.6d) consistently display higher sliding angles than do patterns with smaller features (6.6a and 6.6c) although the area fractions are the same.

The SH-control substrate shows higher critical sliding angle values than SH-virgin paper. It is believed that the enhanced stickiness of the SH-control substrate is caused by minor damage of the plasma-treated surfaces as a result of paper passage through the printer. Heterogeneous damage incurred during passage through the printer is also the likely cause of the relatively large error bars in Figure 6.6 for small drop volumes. Nevertheless, the trends displayed in Figure 6.6 are consistent with differences between

the surfaces and drop sizes investigated. The ability to tune critical sliding angles offers flexibility in the design of 2D microfluidic devices with controlled adhesion. For example, such substrates can be used as drop size filters to slide and transport only drops with a volume above a threshold value.

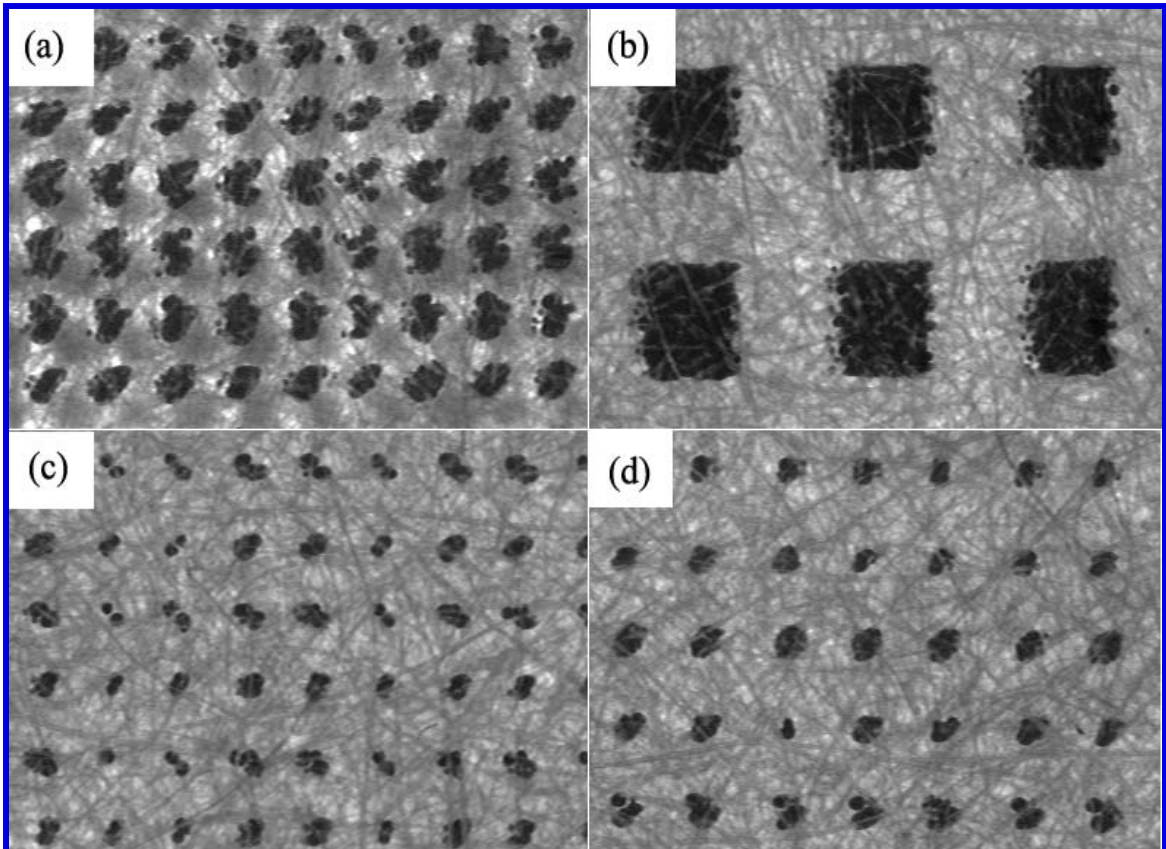


Figure 6.5 Bright field microscope images of ‘checkered’ patterns printed on superhydrophobic paper surfaces

Table 6.1 Area fraction and dimensions of the ink features in ‘checkered’ patterns printed on superhydrophobic paper surfaces

Item	Area fraction, %	Feature Size, μm	Distance between features, μm
a	25.4 ± 0.9	176.3 ± 3.4	346.8 ± 4.6
b	25.7 ± 1.2	532.9 ± 8.1	1041.5 ± 6.0
c	8.1 ± 0.2	107.2 ± 1.7	371.2 ± 1.2
d	7.9 ± 0.4	123.4 ± 2.7	431.7 ± 0.4

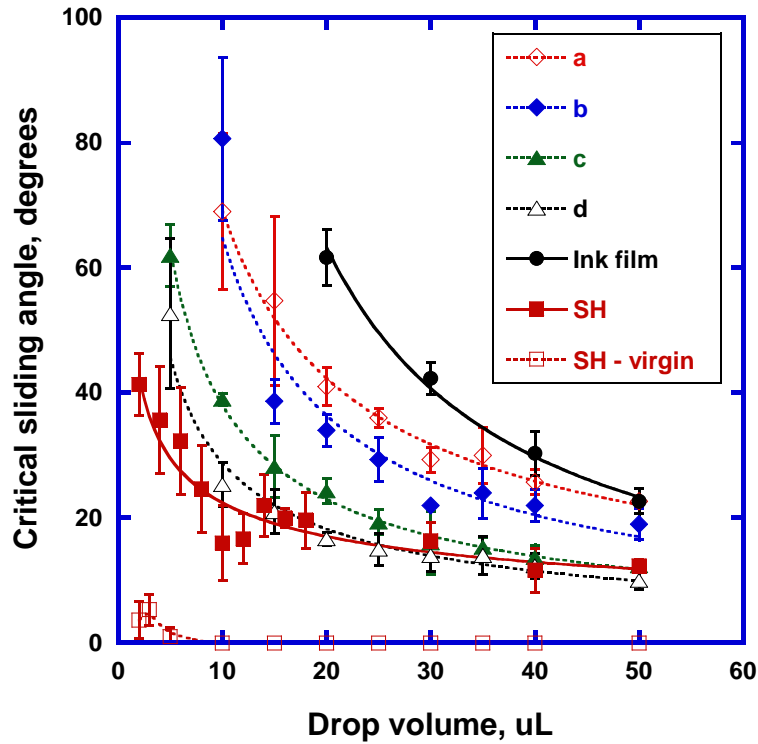


Figure 6.6 Plot of critical sliding angle vs. drop volume for the different ‘checker’ patterns (a, b, c and d) as shown in Figure 6.5, ink film, SH-control and SH-virgin

6.4 Conclusions

The critical sliding angle of a water drop can be manipulated by changing the curvature (positive or negative) of the ink pattern contacting the receding contact line. For instance, a semi-circle ink pattern allows a water drop to slide easily in the convex direction, but not in the concave direction. For a straight line pattern printed on superhydrophobic (SH)

paper, the drop can move along the line more easily than in a direction perpendicular to the line. These concepts were used in this paper to guide the design of fundamental components for 2D microfluidic devices: flow paths, gate/diode, and junctions.

The tunability of critical sliding angle on SH paper obtained by establishing chemical heterogeneity on the paper surfaces was also reported. Area fraction, ink spot sizes, and the distance between ink spots were controlled in order to demonstrate tunability of critical sliding angles on a composite paper surface. These structures allow control of the adhesion of drops on SH paper surfaces to permit the design of 2D drop filters for 2D microfluidic devices.

These studies have demonstrated directional mobility and tunability in the adhesion of water drops on superhydrophobic (SH) surfaces. Control was achieved by imposing geometrical constraints on the contact lines by printing specific ink pattern. The ability to generate these functionalities on paper substrates using commercially available desktop printing techniques offers considerable flexibility in designing paper-based 2D microfluidic devices.

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CHAPTER 7

CONCLUSIONS AND RECOMMENDATIONS

7.1 Conclusions

Most of the artificial superhydrophobic surfaces that have been fabricated to date are not formed on biodegradable, renewable, nor mechanically flexible substrates and the substrate material is often expensive, which limits potential applications. In contrast, cellulose is a biodegradable, renewable, flexible, inexpensive, biopolymer that is abundantly present in nature. Although it satisfies the requirements listed above, cellulose surfaces are not inherently superhydrophobic.

Chapter 2 reported fabrication techniques that were developed to obtain superhydrophobicity on cellulose-based paper. The first processing step invoked domain selective etching of the amorphous portions of the cellulose in an oxygen plasma to ‘uncover’ the crystalline cellulose microfibrils. The etched surface was then coated with a thin fluorocarbon film deposited via plasma enhanced chemical vapor deposition using pentafluoroethane as a precursor. Variation of plasma etch time yielded two types of superhydrophobicity: “roll-off” (contact angle (CA): $166.7^{\circ} \pm 0.9^{\circ}$ and CA hysteresis: $3.4^{\circ} \pm 0.1^{\circ}$) and “sticky” (CA: $153.4^{\circ} \pm 4.7^{\circ}$ and CA hysteresis: $149.8 \pm 5.8^{\circ}$) superhydrophobicity. While “sticky” superhydrophobic substrates combine a high static water contact angle with strong droplet adhesion, “roll-off” superhydrophobic substrates exhibit very low droplet adhesion, resulting in high droplet mobility, reminiscent of the properties of lotus leaves. The key property required for roll-off superhydrophobicity was nano-meter scale roughness, which was achieved by exposing the internal structure

of each fiber via selective etching, instead of depositing foreign material (*e.g.*, polymers, nanoparticles) as had been pursued in previous studies. The combination of this nano-meter scale roughness with the inherent micro-meter scale roughness of cellulose fibers on a paper surface resulted in robust superhydrophobic surfaces. This result was distinct relative to roughened structures created by traditional polymer grafting, nanoparticle deposition or other artificial means.

Chapter 3 described the methods that were established to enable tunability of the contact angle hysteresis on superhydrophobic paper surfaces from $149.8 \pm 5.8^\circ$ to $3.5 \pm 1.1^\circ$. Continuous variations in contact angle hysteresis were achieved by controlling the nano-structure of the fibers through the oxygen etching time. The change in wetting properties was attributed to a transition of the liquid-surface interaction from a Wenzel state to a Cassie state on the nano-scale, while maintaining a Cassie state on the micro-scale.

In Chapter 4, the effects of fiber types and paper making parameters on the fabrication of superhydrophobic paper were studied. No significant difference in the formation of nano-scale features via oxygen etching was observed for the different fiber types, which underlined the generality of the superhydrophobicity treatment. The paper making process controls the structure of the paper web and thus the micro-scale roughness. Paper making parameters therefore strongly affects “sticky” superhydrophobicity, which is dominated by the micro-scale roughness. However, because the etching process creates similar nano-scale roughness on all fibers, all substrates tested in this thesis (from various paper making techniques) became “roll-off” superhydrophobic. The study showed that a

30 minute oxygen etching process in combination with a 2 minute PFE deposition yields the optimum roughness and film thickness (~100nm) for roll-off superhydrophobic paper surfaces with good physical properties. The significance of water drop volume in contact angle measurements was also discussed, which is particularly relevant for the characterization and analysis of superhydrophobic behavior of heterogeneous, porous substrates like paper. It was concluded that while measuring the CA values on heterogeneous, porous substrates like paper, it is important to select drop sizes that 1) are larger than the surface roughness length scales and 2) provide sufficient image resolution for the goniometer to identify the three-phase contact line.

Chapter 5 described a method to pattern superhydrophobic surfaces with high surface energy black ink by using commercially available desktop printing technology. The shape and size of the ink islands control the adhesive forces between the surface and water drops in two directions, parallel ('drag-adhesion') and perpendicular ('extensional-adhesion') to the substrate. Experimental data on the adhesive forces showed good agreement with classical models for 'drag' (Furmidge equation) and 'extensional' adhesion (modified Dupré equation). The tunability of the two adhesive forces was used to design basic patterns that can be used to perform four basic unit operations for the manipulation of liquid drops on the paper substrates: storage, transfer, mixing and sampling. By combining these basic functionalities, it is possible to design simple two-dimensional lab-on-paper (LOP) devices. In this 2D LOP prototype, liquid droplets adhere to the porous substrate rather than absorbing into the paper; as a result, liquid

droplets remain accessible for further quantitative testing and analysis, after performing simple qualitative on-chip testing.

Chapter 6 discussed in detail how directional mobility and controlled adhesion of water drops can be obtained on superhydrophobic paper by creating more complex ink patterns. By changing the curvature/shapes of the ink patterns, anisotropic geometrical constraints can be imposed on the movement of the three-phase liquid/solid/vapor contact line. Through careful pattern design, the critical sliding angle of water drops was manipulated with respect to sliding direction. The adhesion of water drops on superhydrophobic paper was controlled by designing chemical heterogeneity in the form of “checkered” ink patterns. Printing ink patterns with varying area fractions and feature sizes resulted in variable critical sliding angles of water drops. Directional mobility and tunable adhesion of water drops on superhydrophobic paper surfaces can be applied to design novel components for emerging two-dimensional LOP technology, such as 2D flow paths, gates/diodes, junctions and drop size filters. Finally, the use of commercially available desktop printers and word processing software to generate ink patterns enable end-user design of LOP devices for user-specific applications, which adds significant flexibility to this technology.

7.2 Recommendations

This thesis project identified interesting questions and directions that are recommended for future research in this field.

Superhydrophobicity or water repellency on paper surfaces has been extensively discussed in this thesis. Water is a strongly hydrogen bonded system with high cohesive forces, resulting in a surface tension of ~ 73 mN/m. The high cohesive forces inhibit water penetration into the nano-scale features that are coated with a low surface energy coating, and thus give rise to superhydrophobicity of the modified substrate. Non-polar liquids have surface tensions that are much lower than those for water. Because of the weaker cohesive forces, these liquids can more easily penetrate into the nano- and micro-scale structures that are coated with fluorocarbon layers. As a result, it is challenging to design a surface chemistry and roughness topography that yield roll-off behavior for non-polar liquids. Repellency to non-polar liquids, which has been dubbed superoleophobicity, is an emerging field. Recently, it has been reported that the angle at which the liquid enters a geometrical feature, referred to as 're-entrant angle', determines to a great extent the penetration ability of the liquid [1]. Therefore designing appropriate re-entrant angles on the roughened nano-scale features should be a viable strategy to render a surface superoleophobic [1]. Methods to impart controlled re-entrant angles on the nano-scale features on cellulose surfaces should be devised and the resulting superoleophobic behavior determined. Superoleophobic paper surfaces would have significant application in the food packaging industry.

Chapters 5 and 6 described the operating principles and functionalities that have been designed to create LOP microfluidic devices. Nevertheless, we have not presented a bench-scale prototype of a functional LOP microfluidic device. Further research should be performed to design and fabricate a prototype LOP device by combining the fundamental building blocks described in this thesis work. Use of ink chemistries containing indicators and analyte molecules should further enhance the functionality of LOP devices, but this requires development of the appropriate printing technology, which was beyond the scope of the current thesis.

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VITA

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